

Modeling mass independent of anisotropy

A comparison between Milky Way and Andromeda satellites
(and musings on density slope determinations)



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Yale: Marla Geha



Ricardo Munoz



Heigh Ho...

Collaborators

Outline



1. A new mass estimator: accurate without knowledge of anisotropy/beta
2. Applications of new mass determinations for MW dSphs
3. Comparison between MW and M_{31} dSphs
4. The skinny on slope determinations



Cardinal rule about dwarfs:

It's not the size of the boat, but
the motion of the ocean...



Mass modeling of hot systems

Many gas-poor dwarf galaxies have a significant, usually dominant hot component. They are dispersion supported, not rotation supported.

Consider a spherical, dispersion supported system whose stars are collisionless and are in equilibrium. Let us consider the Jeans Equation:

$$r \frac{d(\rho_{\star} \sigma_r^2)}{dr} = \frac{-GM(r)}{r} \rho_{\star}(r) - 2\beta(r) \rho_{\star} \sigma_r^2$$

We want mass

*Unknown:
Anisotropy*

$$\beta \equiv 1 - \frac{\sigma_t^2}{\sigma_r^2}$$

Free function

*Assume known:
3D deprojected
stellar density*

*Radial
dispersion
(depends
on beta)*

Mass modeling of hot systems

Jeans Equation

$$r \frac{d(\rho_{\star} \sigma_r^2)}{dr} = \frac{-GM(r)}{r} \rho_{\star}(r) - 2\beta(r) \rho_{\star} \sigma_r^2$$

Velocity
Anisotropy
(3 parameters)

$$\beta(r) = (\beta_{\infty} - \beta_0) \frac{r^2}{r_{\beta}^2 + r^2} + \beta_0$$

Mass modeling of hot systems

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Velocity Anisotropy
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Mass Density
(6 parameters)

$$\rho(r) = \frac{\rho_s e^{-r/r_{cut}}}{(r/r_s)^c [1 + (r/r_s)^a]^{(b-c)/a}}$$

Mass modeling of hot systems

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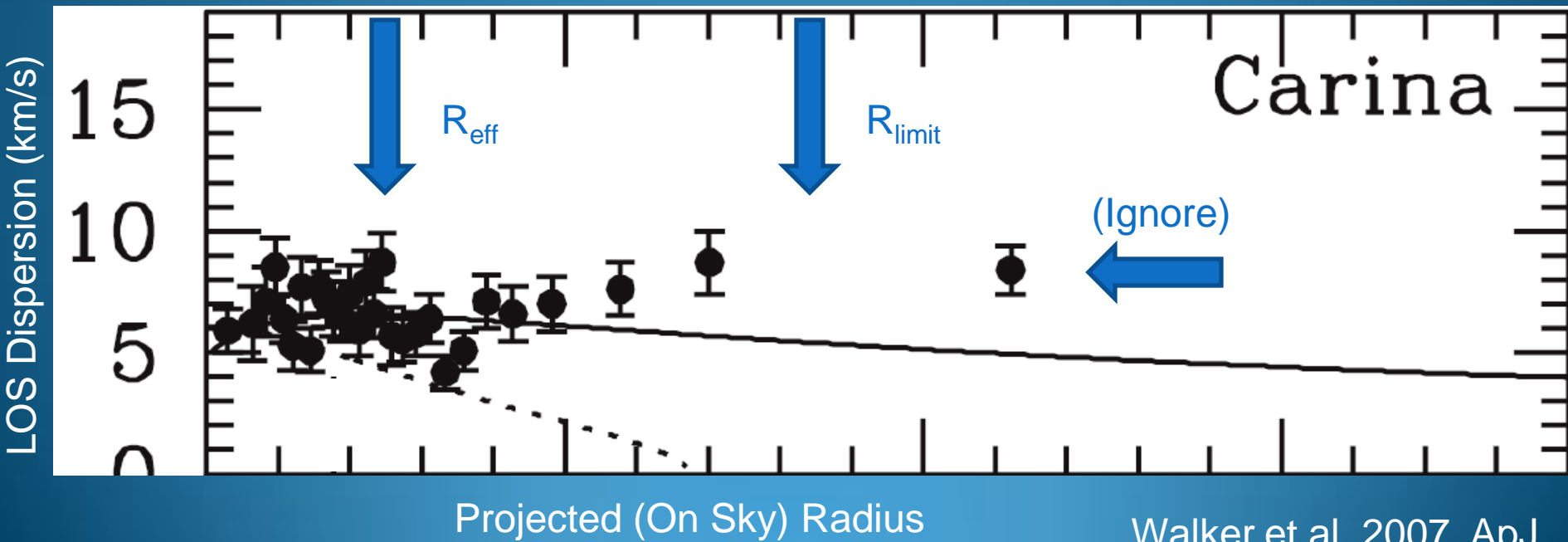
Mass Density
(6 parameters)

$$\rho(r) = \frac{\rho_s e^{-r/r_{cut}}}{(r/r_s)^c [1 + (r/r_s)^a]^{(b-c)/a}}$$

Using a Gaussian PDF for the observed stellar velocities, we marginalize over all free parameters (including photometric uncertainties) using a Markov Chain Monte Carlo (MCMC).

Thought Experiment

Given the following kinematics...



Thought Experiment



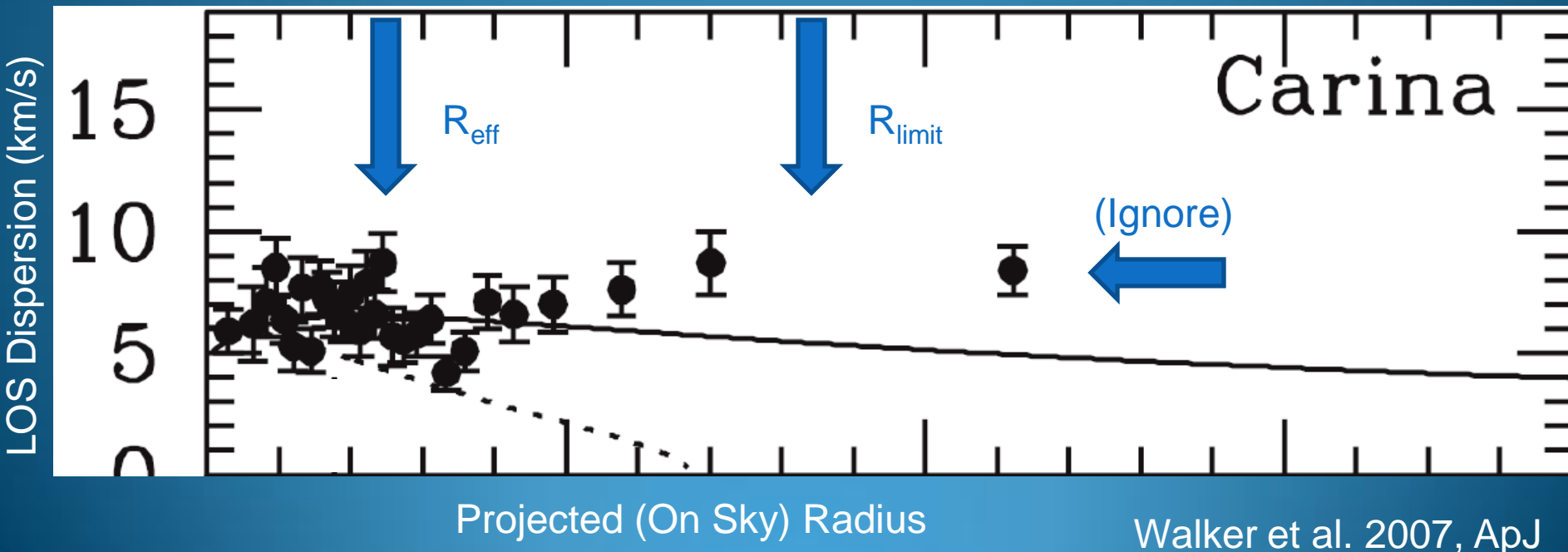
Given the following kinematics, will you derive a better constraint on mass enclosed within:

a) $0.5 * r_{1/2}$

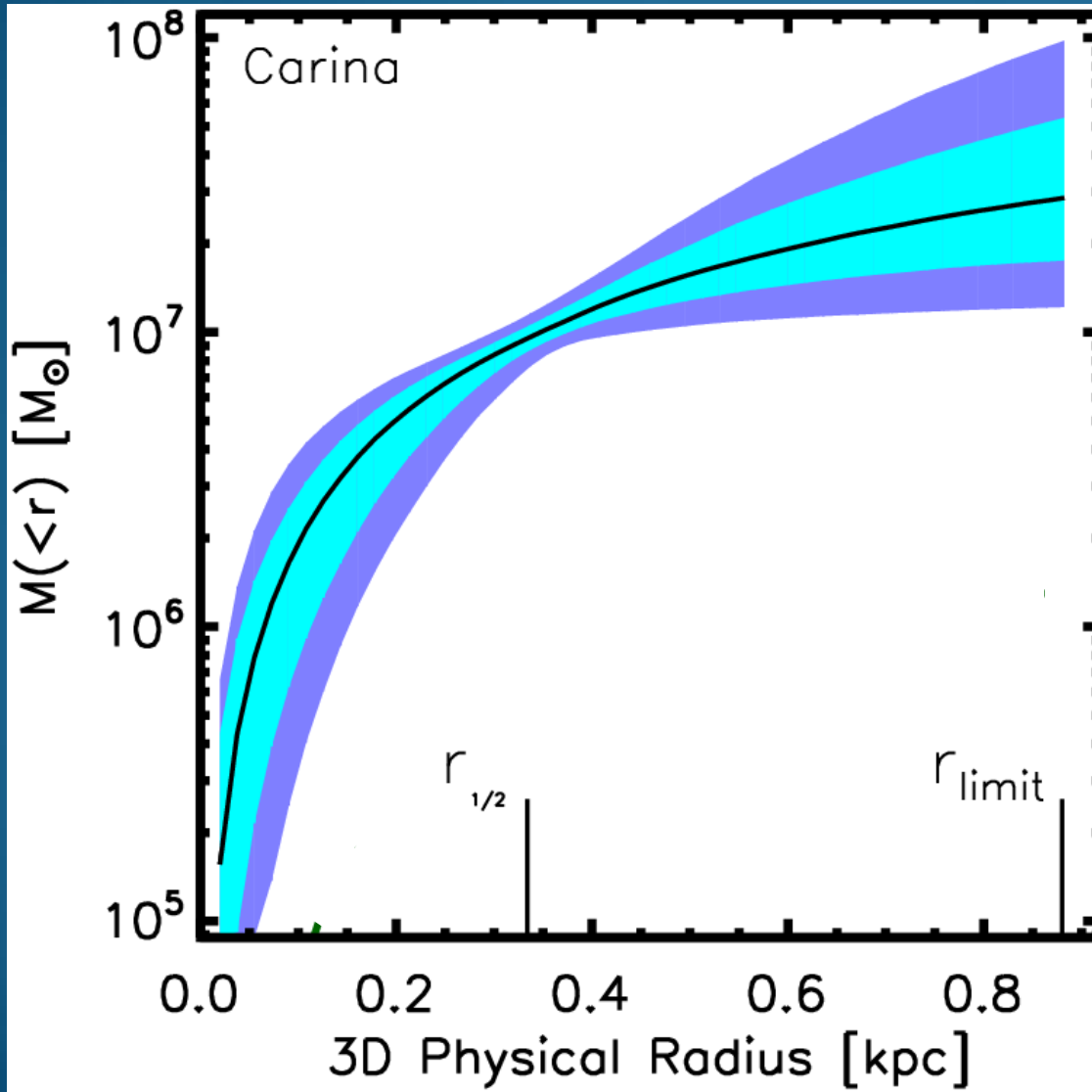
b) $r_{1/2}$

c) $1.5 * r_{1/2}$

Where $r_{1/2}$ is the derived 3D deprojected half-light radius of the system.
(The sphere within the sphere containing half the light).



Hmm...

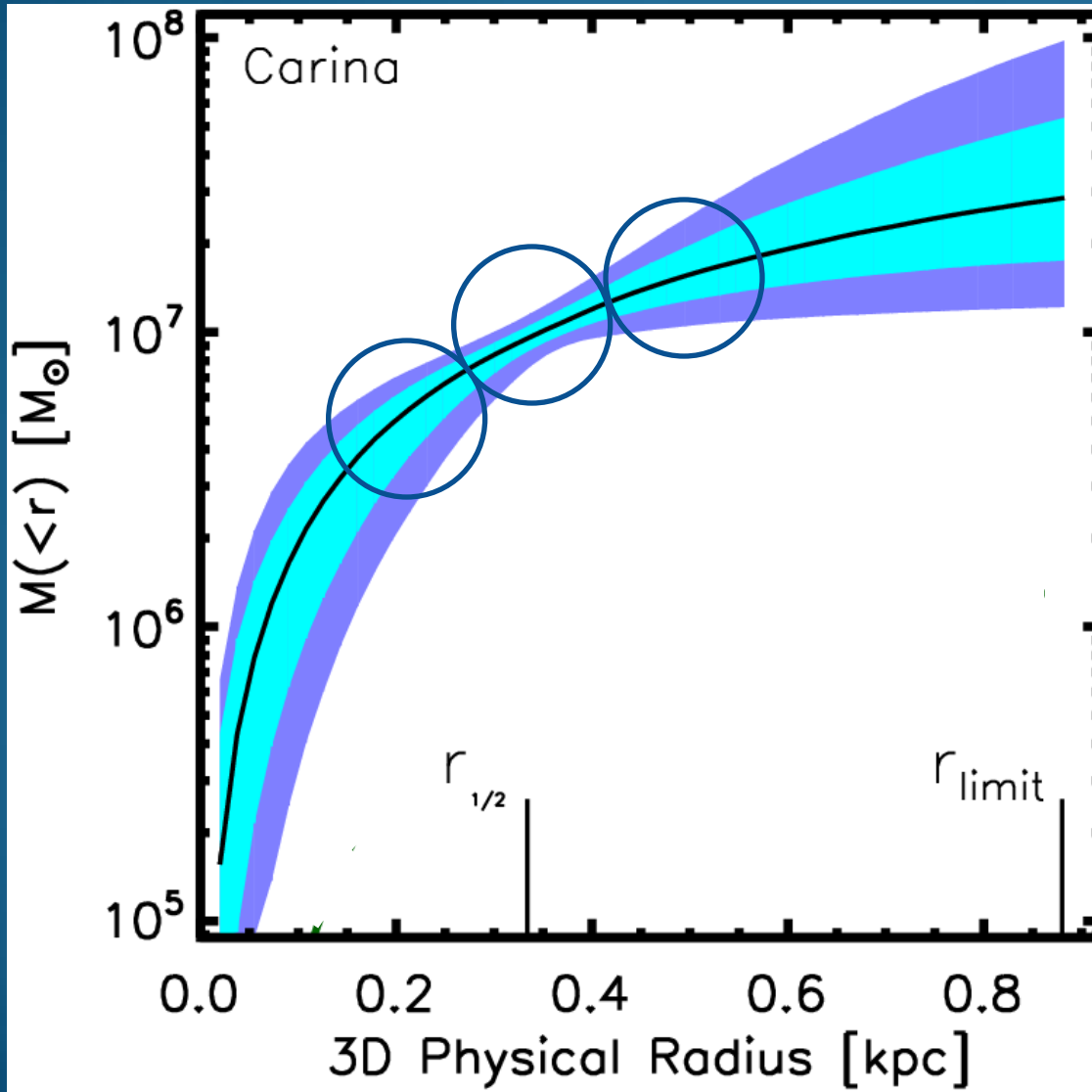


Confidence Intervals:
Cyan: 68%
Purple: 95%

Joe Wolf et al., in prep

Hmm...

It turns out that the mass is best constrained within $r_{1/2}$, and despite the given data, is less constrained for $r < r_{1/2}$ than $r > r_{1/2}$.

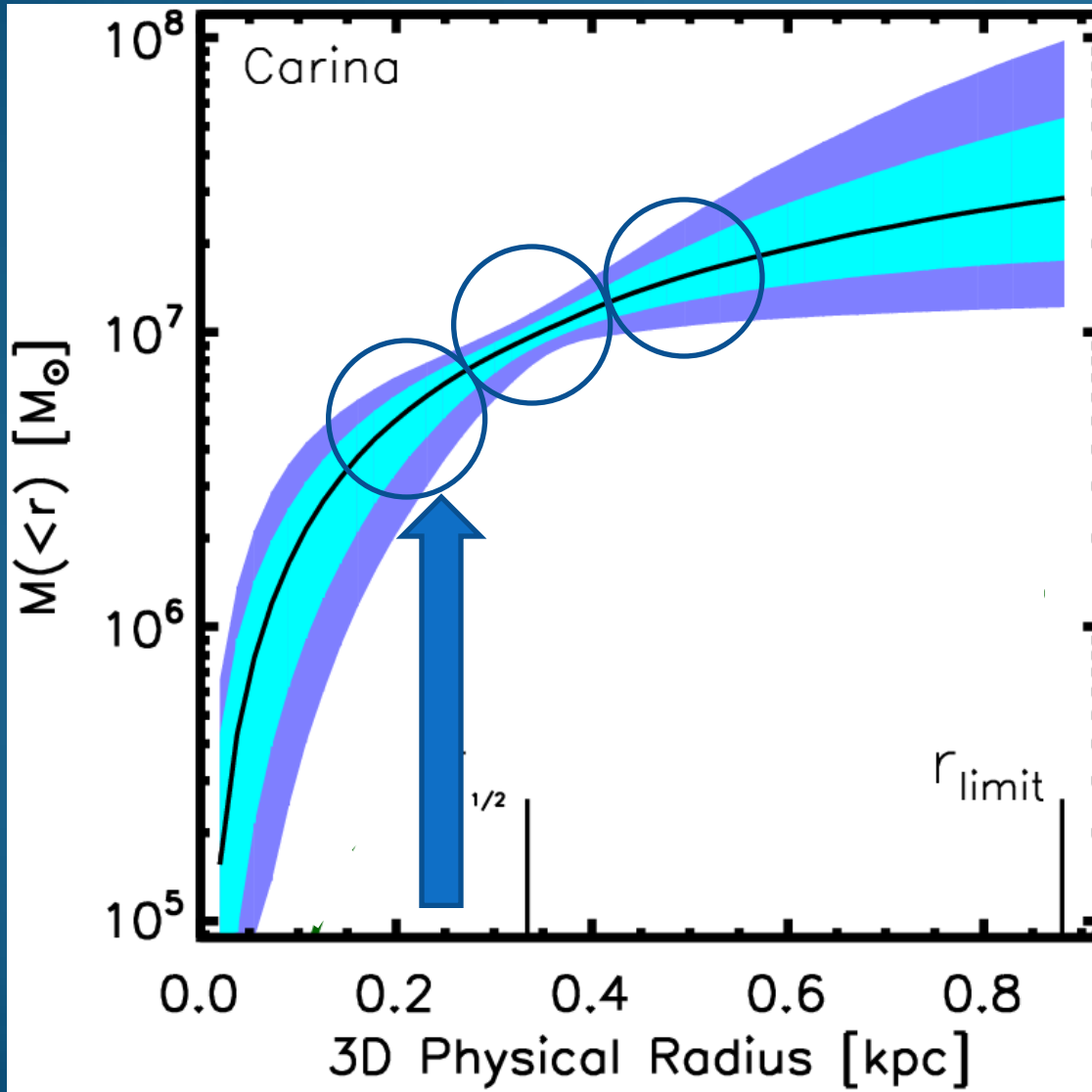


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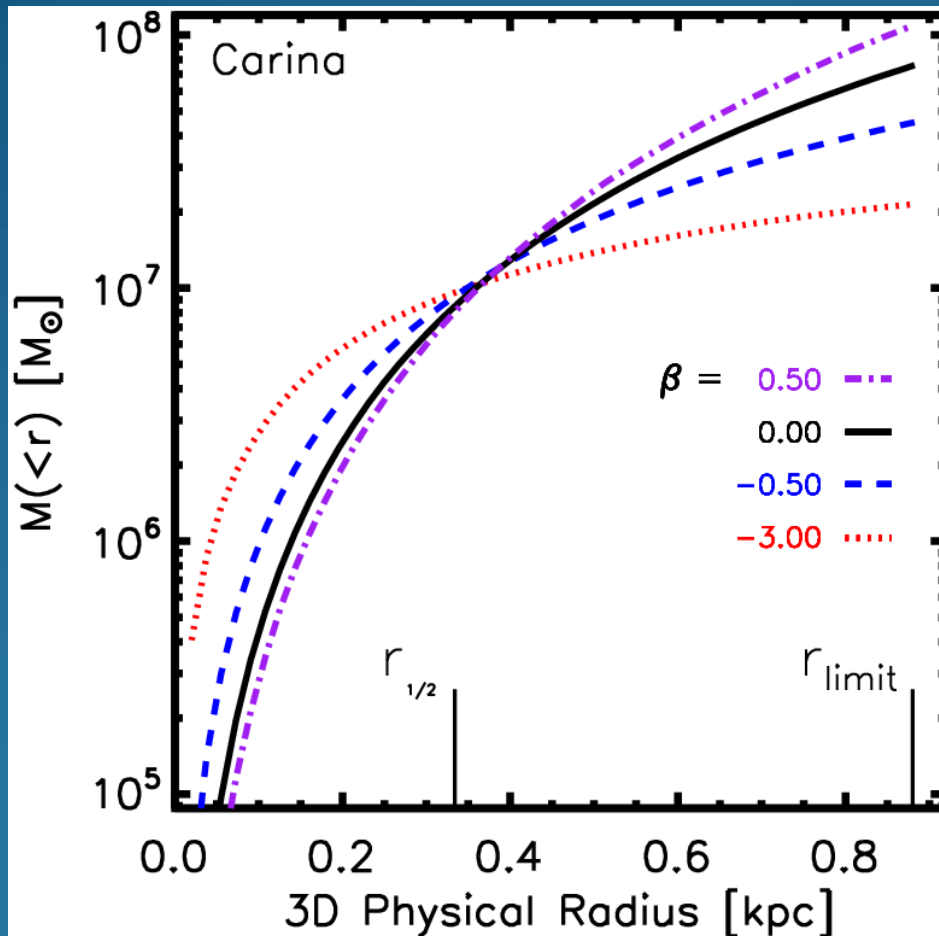
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Anisotrwhat?



Radial Anisotropy

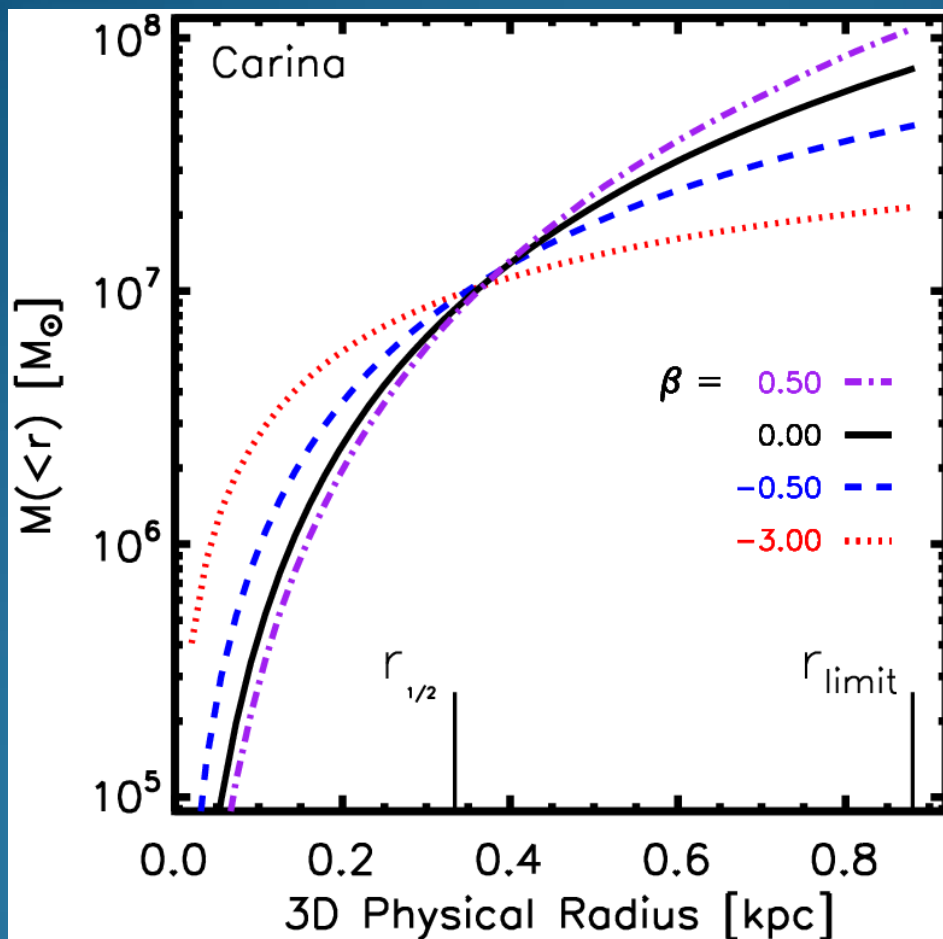
Isotropic

Tangential

Joe Wolf et al., in prep

Center of system:
Observed dispersion is radial

Anisotrwhat?



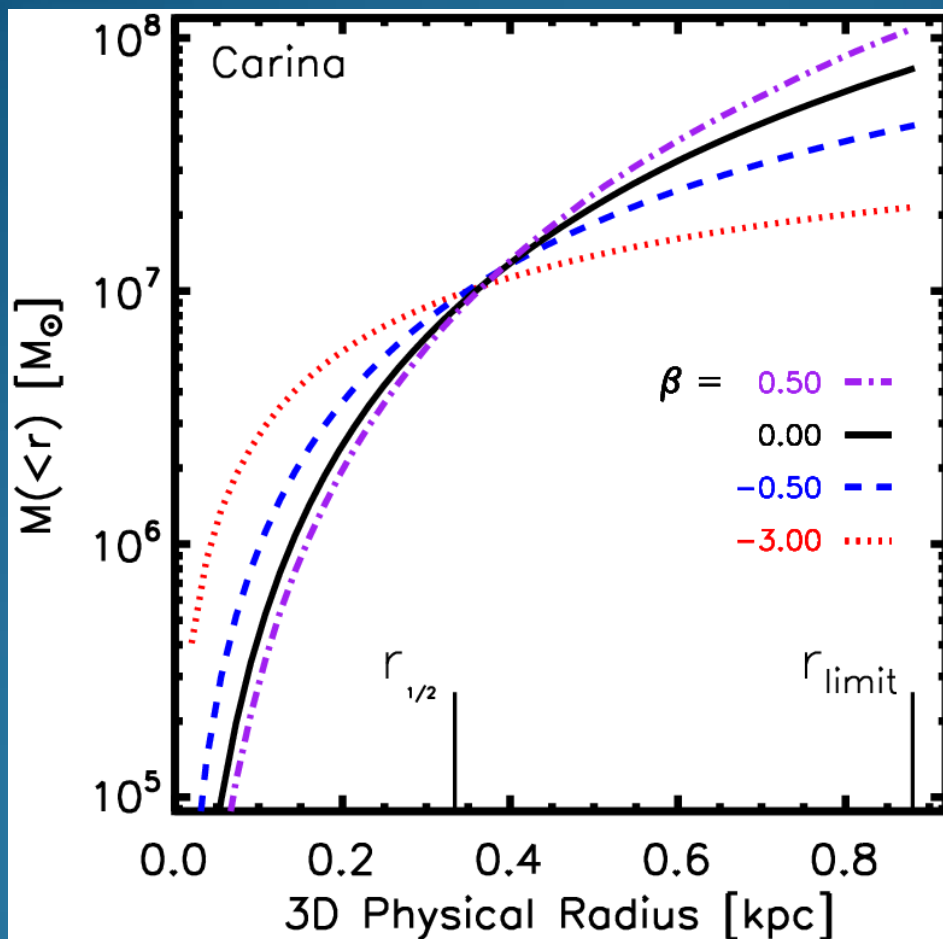
Edge of system: Observed dispersion is tangential

- Radial Anisotropy
- Isotropic
- Tangential

Joe Wolf et al., in prep

Center of system:
Observed dispersion is radial

Anisotrwhat?



Edge of system: Observed dispersion is tangential

- ← Radial Anisotropy
- ← Isotropic
- ← Tangential

Newly derived analytic equations **predict** that the effect of anisotropy is minimal $\sim r_{1/2}$. E.g.:

$$M(< r; 0) - M(< r; \beta) = \frac{\beta(r) r \sigma_r^2(r)}{G} \left(\frac{d \ln \rho_\star}{d \ln r} + \frac{d \ln \sigma_r^2}{d \ln r} + \frac{d \ln \beta}{d \ln r} + 3 \right)$$

Joe Wolf et al., in prep

Mass-anisotropy degeneracy has effectively been *terminated* at $r_{1/2}$:

Derived equation under several simplifications:

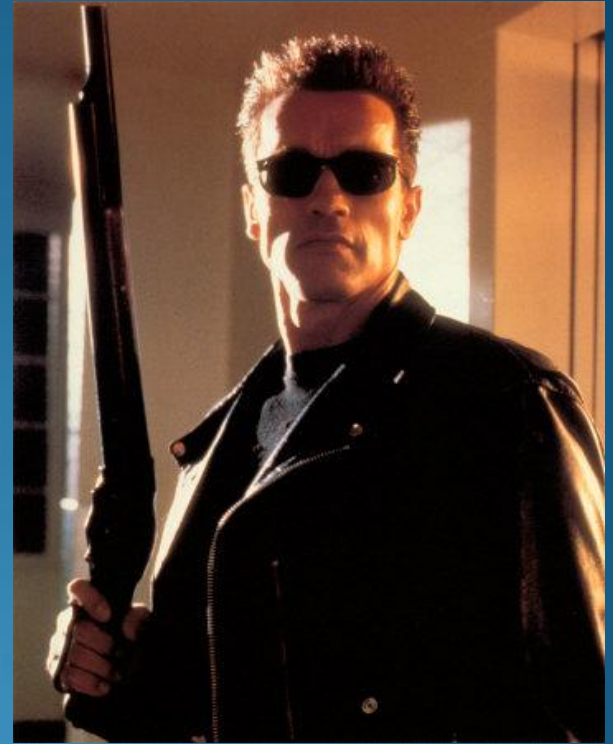
$$M_{1/2} = 3 G^{-1} r_{1/2} \langle \sigma_{\text{los}}^2 \rangle$$



Mass-anisotropy degeneracy has effectively been *terminated* at $r_{1/2}$:

Derived equation under several simplifications:

$$M_{1/2} = 3 G^{-1} r_{1/2} \langle \sigma_{\text{los}}^2 \rangle$$



$$\frac{M_{1/2}}{M_{\odot}} \simeq 930 \frac{R_{\text{eff}}}{\text{pc}} \frac{\langle \sigma_{\text{los}}^2 \rangle}{\text{km}^2 \text{ s}^{-2}}$$

$$r_{1/2} \simeq \frac{4}{3} * R_{\text{eff}}$$

Wait a second...

Isn't this just the scalar virial theorem (SVT)?

$$M_{1/2} = 3 G^{-1} r_{1/2} \langle \sigma_{\text{los}}^2 \rangle$$

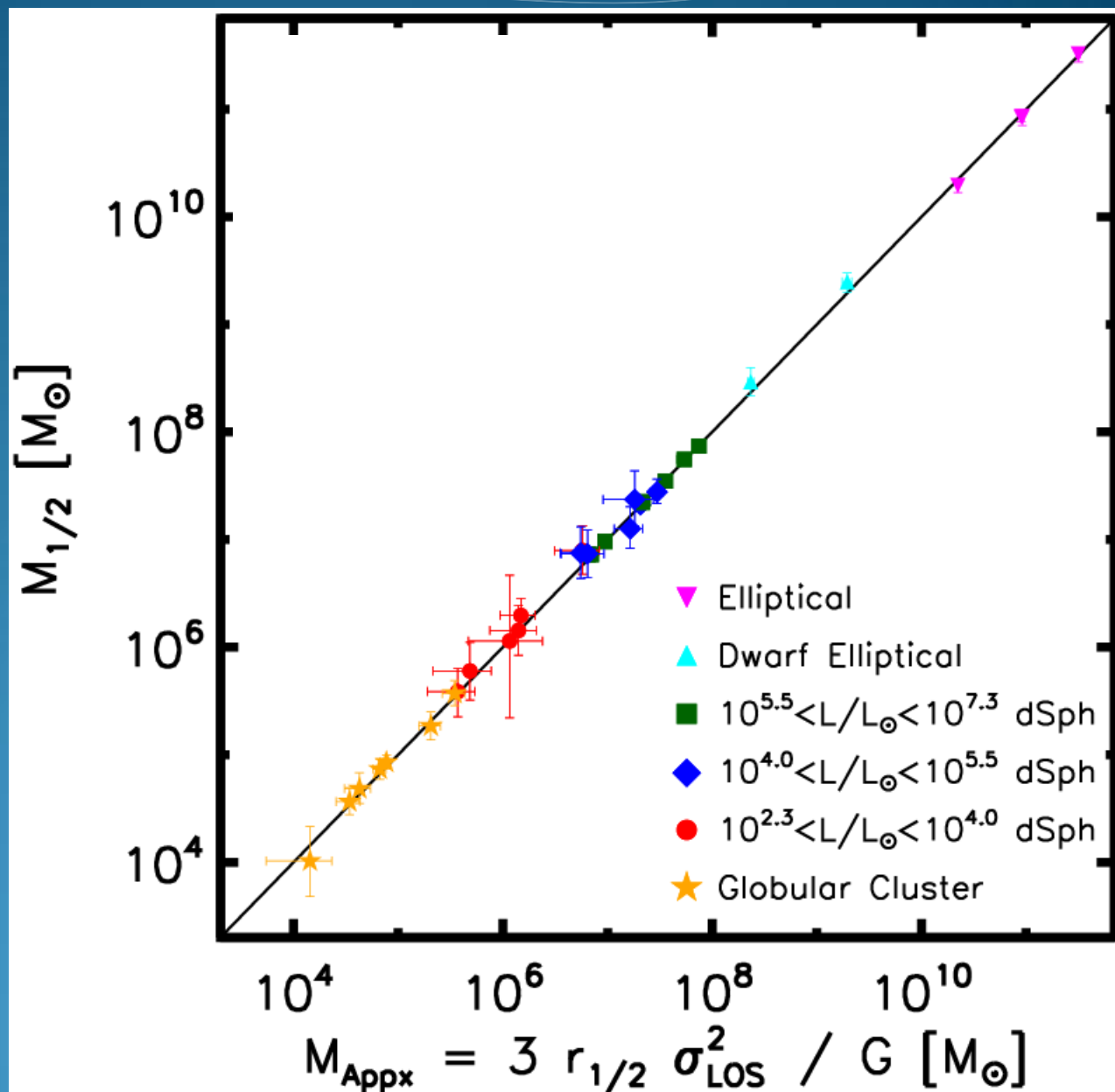
Nope! The SVT only gives you limits on the total mass of a system.

This formula yields the mass within $r_{1/2}$, the 3D deprojected half-light radius, and is accurate independent of our ignorance of anisotropy.

Really?

Boom!

Equation tested on systems spanning almost **eight** decades in half-light mass after lifting simplifications.

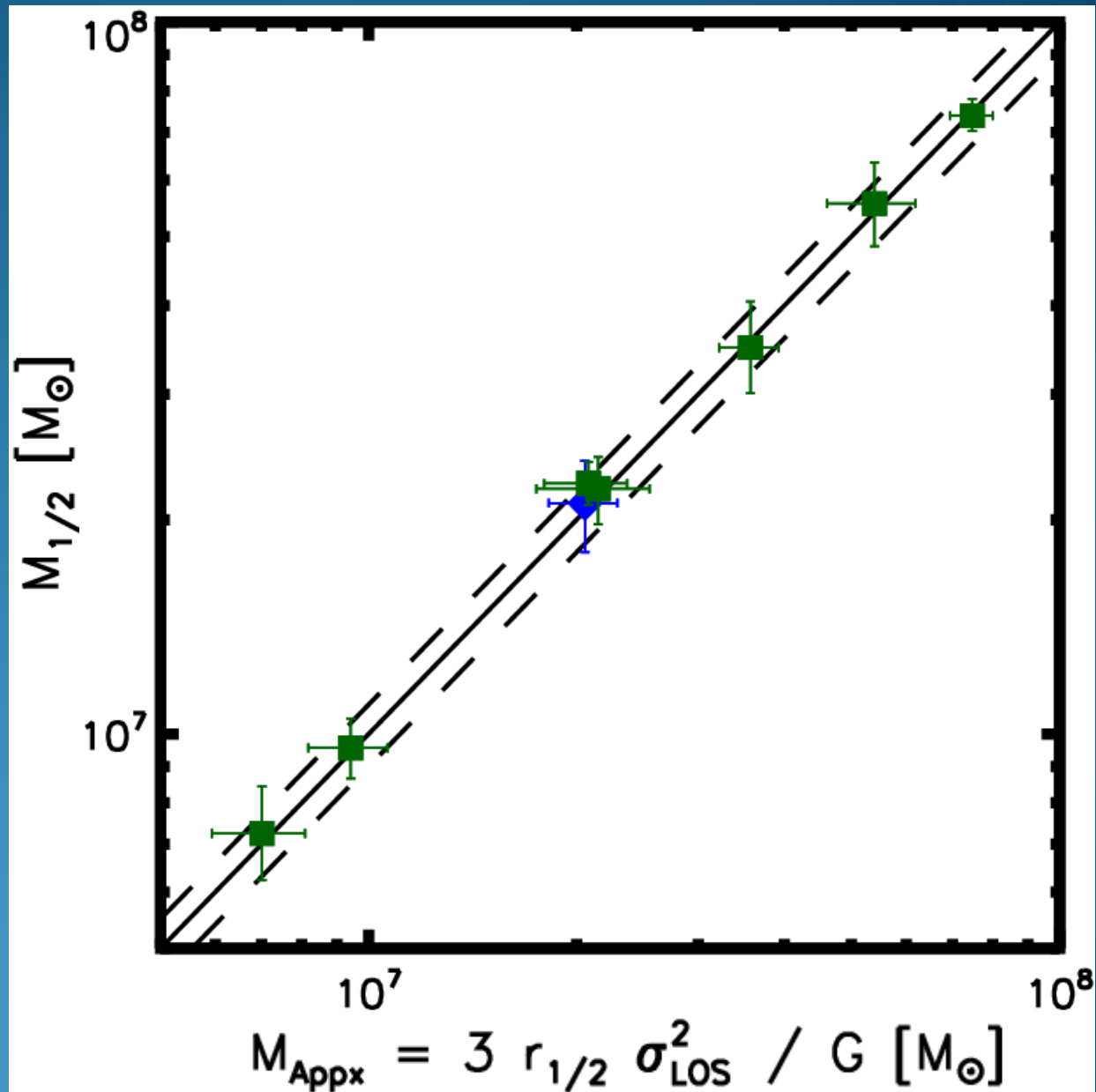


Boom!

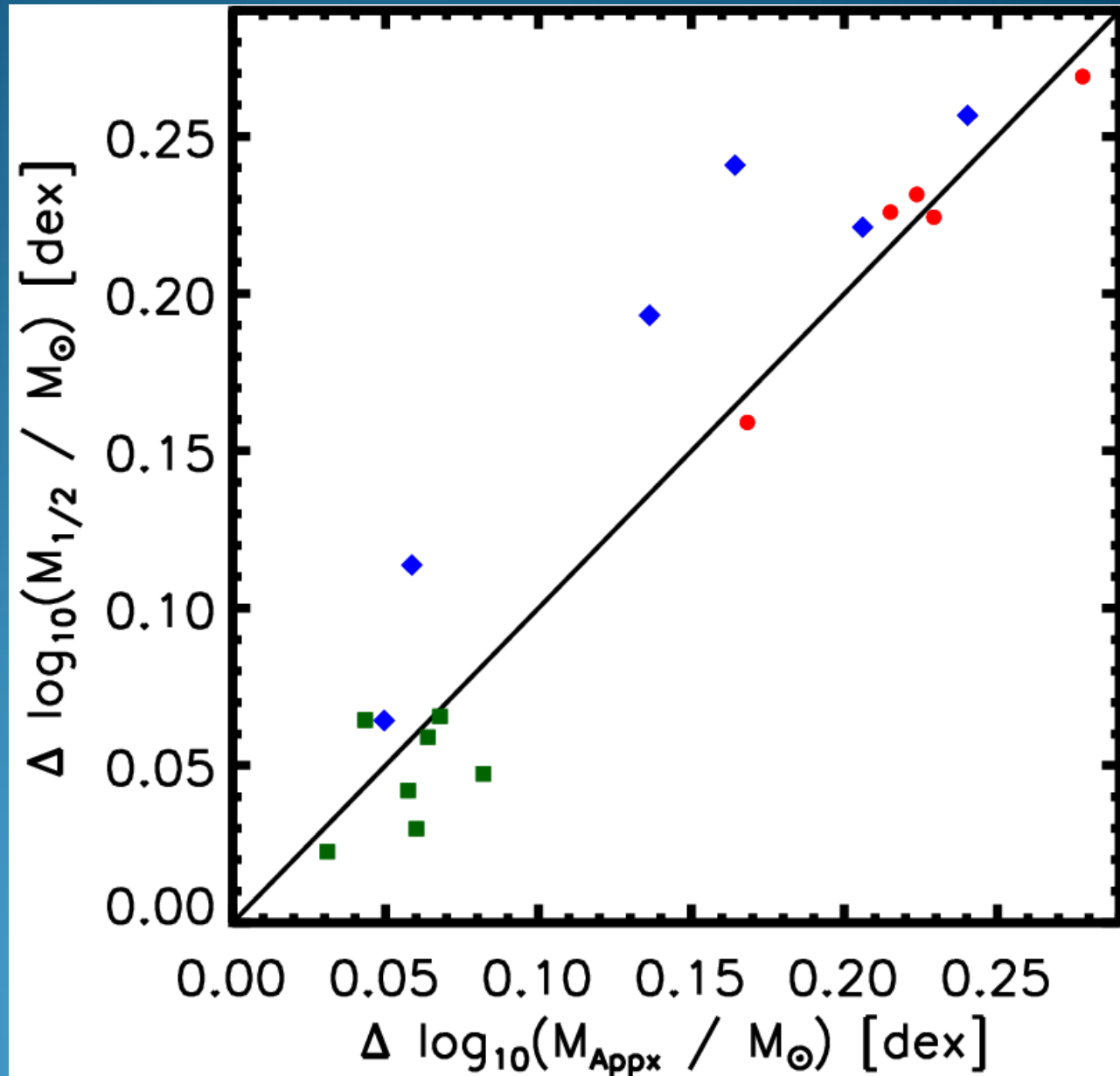
“Classical” MW dwarf spheroidals



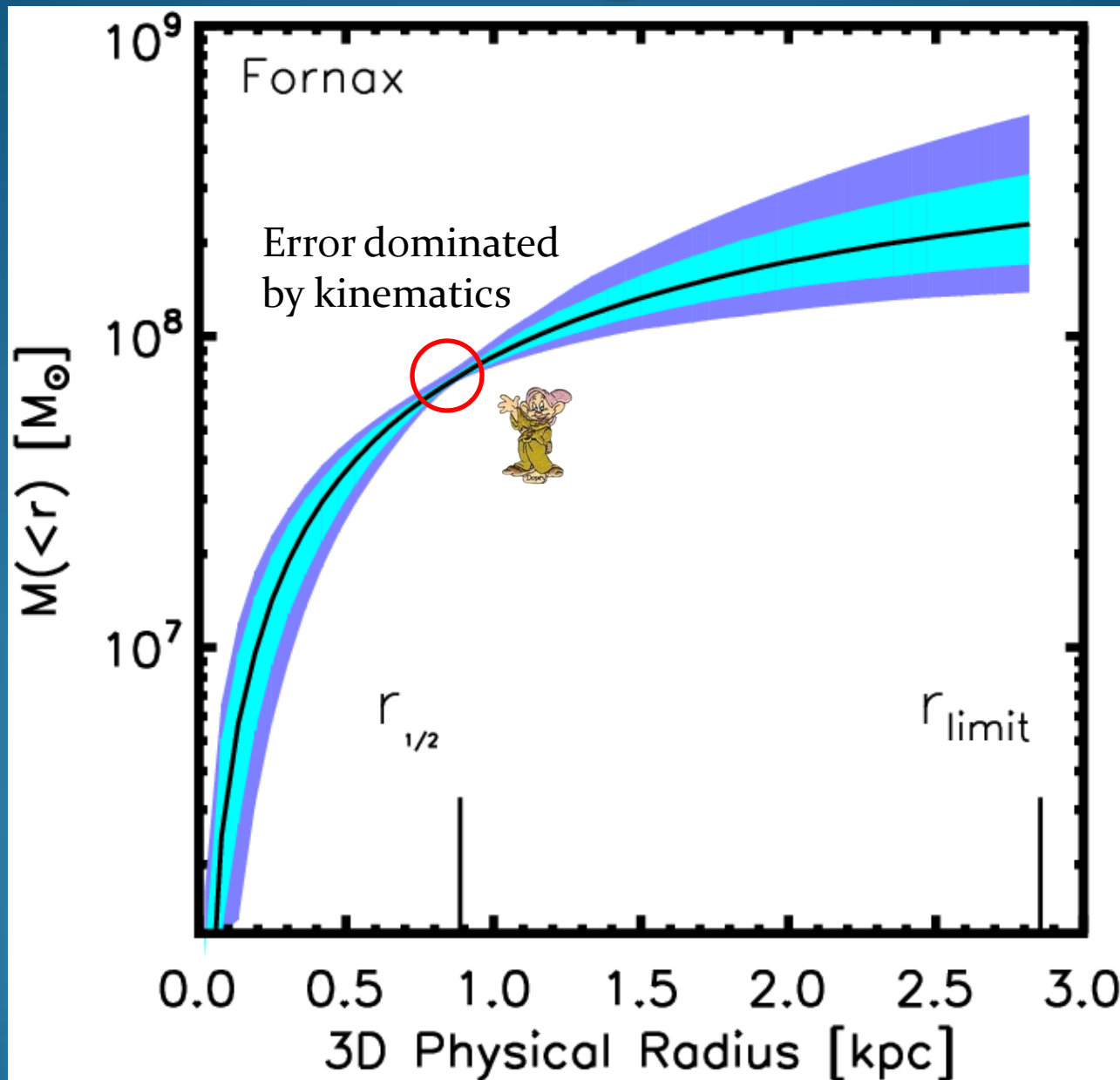
Dotted lines:
10% variation in
factor of 3 in M_{Appx}



For Oleg (how you doin'?)

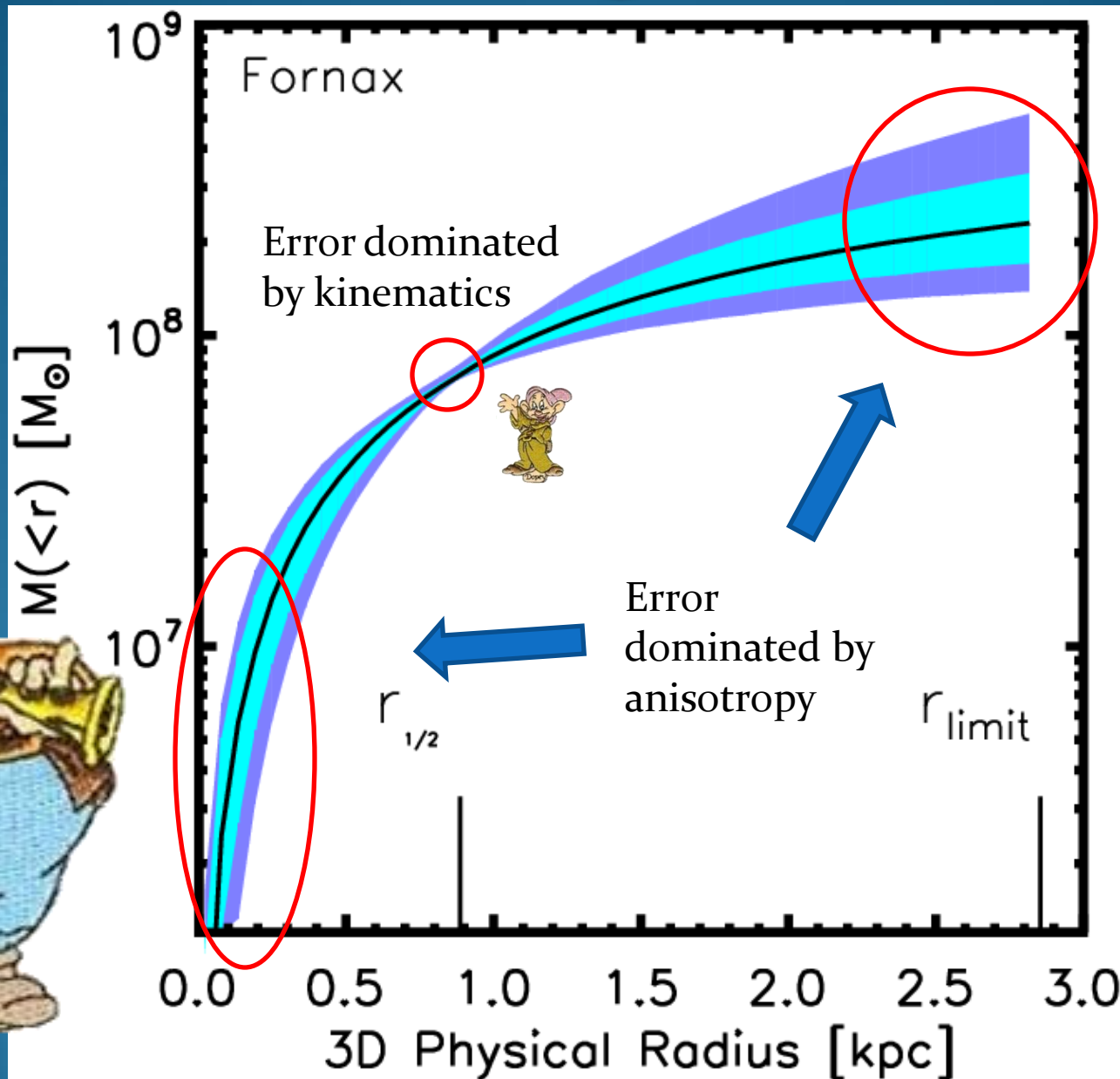


Mass Errors: Origins

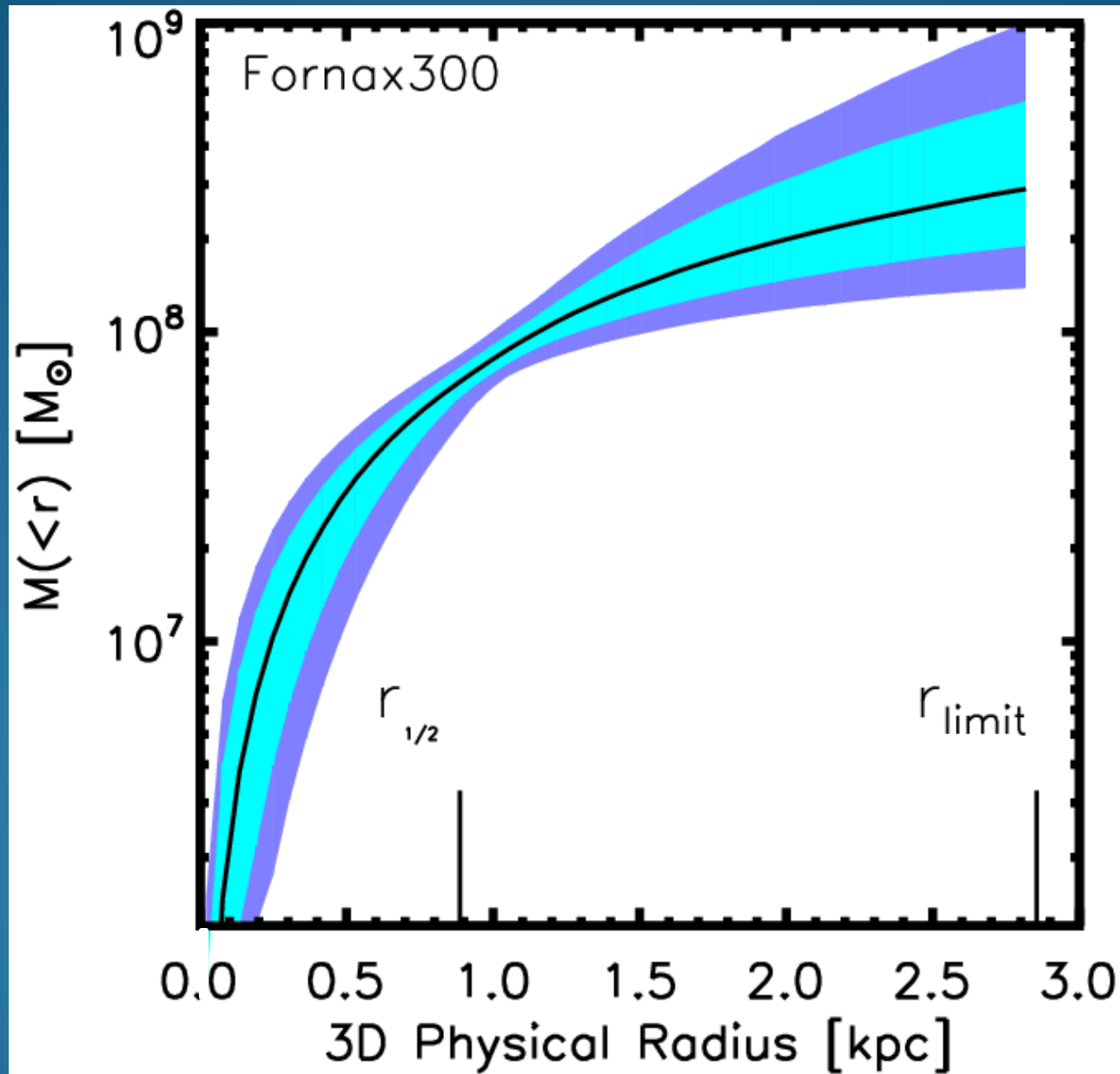


Joe Wolf et al., in prep

Mass Errors: Origins

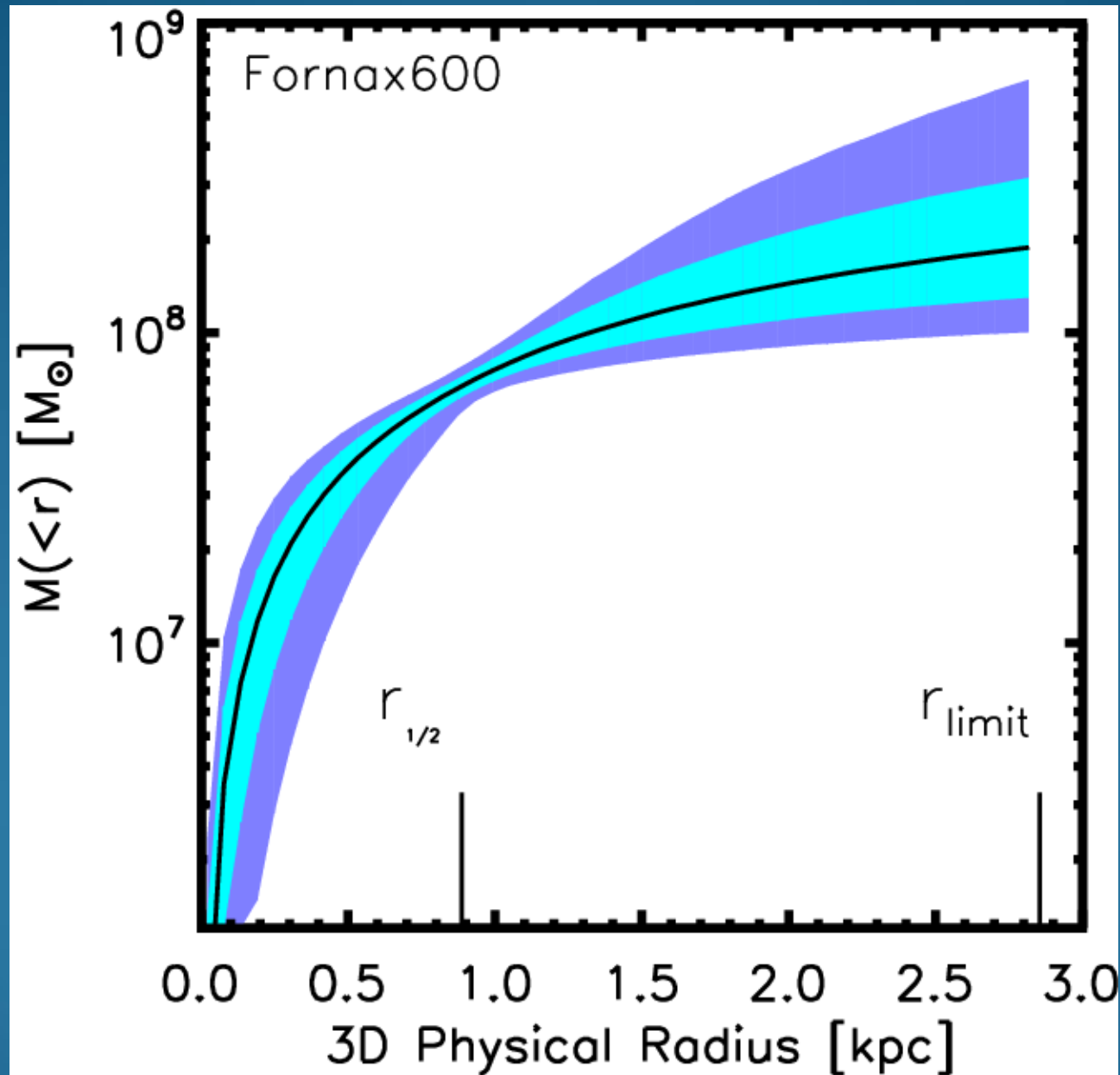


Mass Errors: 300 stars



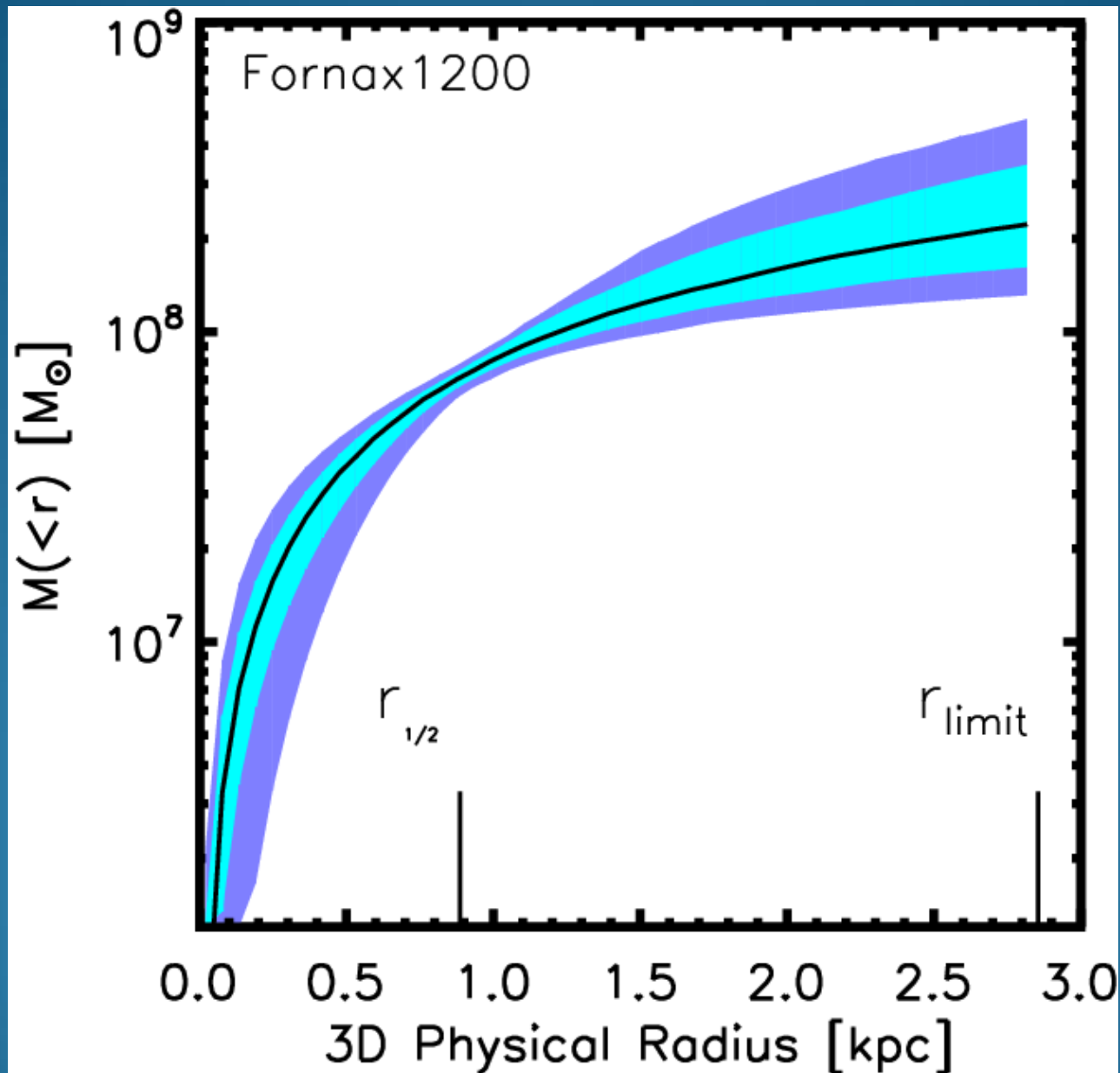
Joe Wolf et al., in prep

Mass Errors: 600 stars



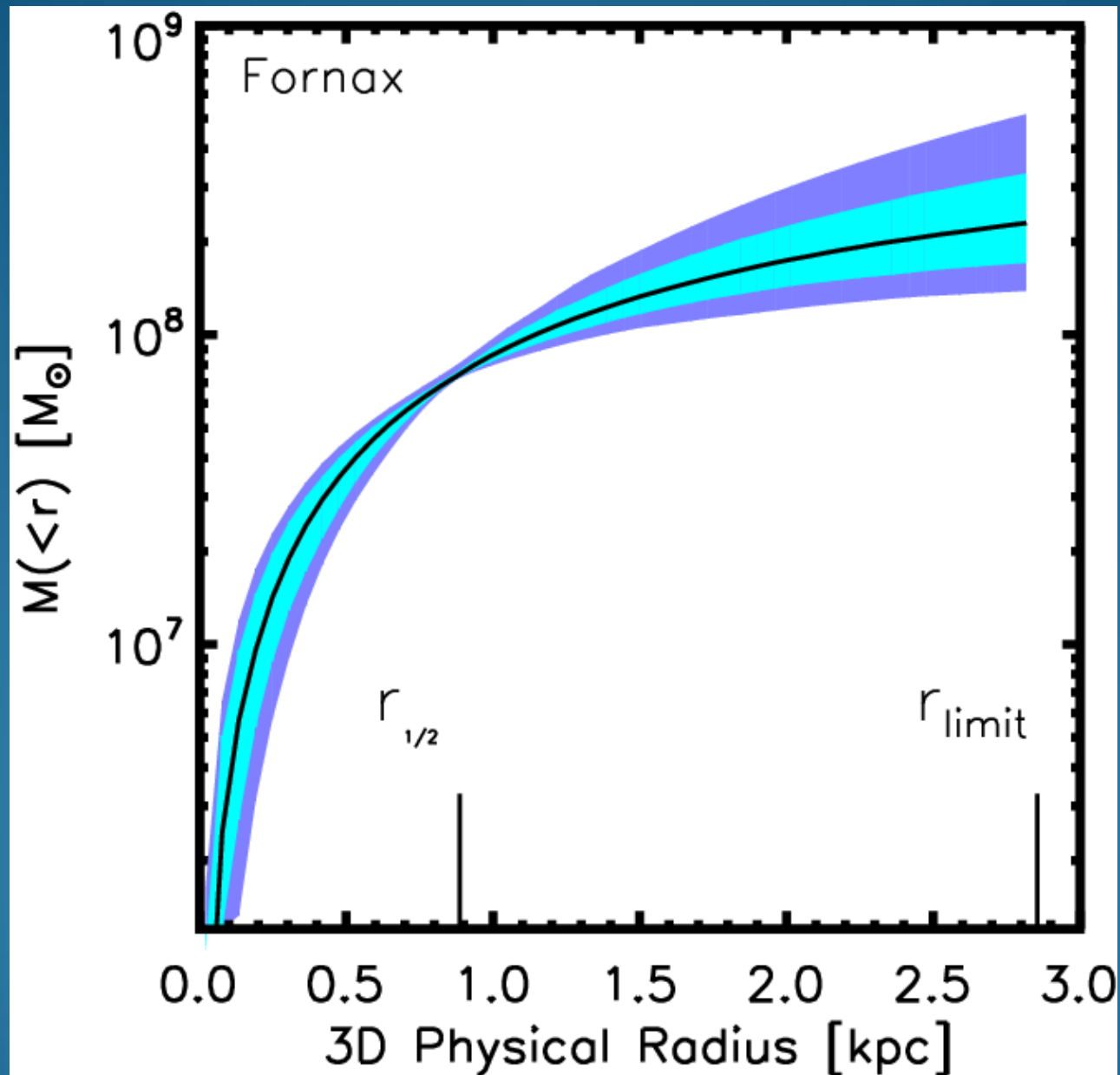
Joe Wolf et al., in prep

Mass Errors: 1200 stars



Joe Wolf et al., in prep

Mass Errors: 2400 stars



Joe Wolf et al., in prep

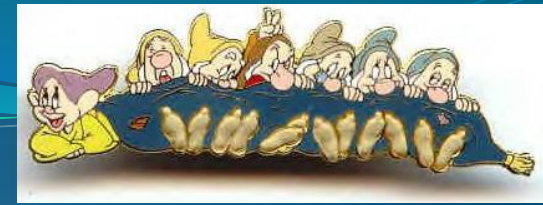
Outline



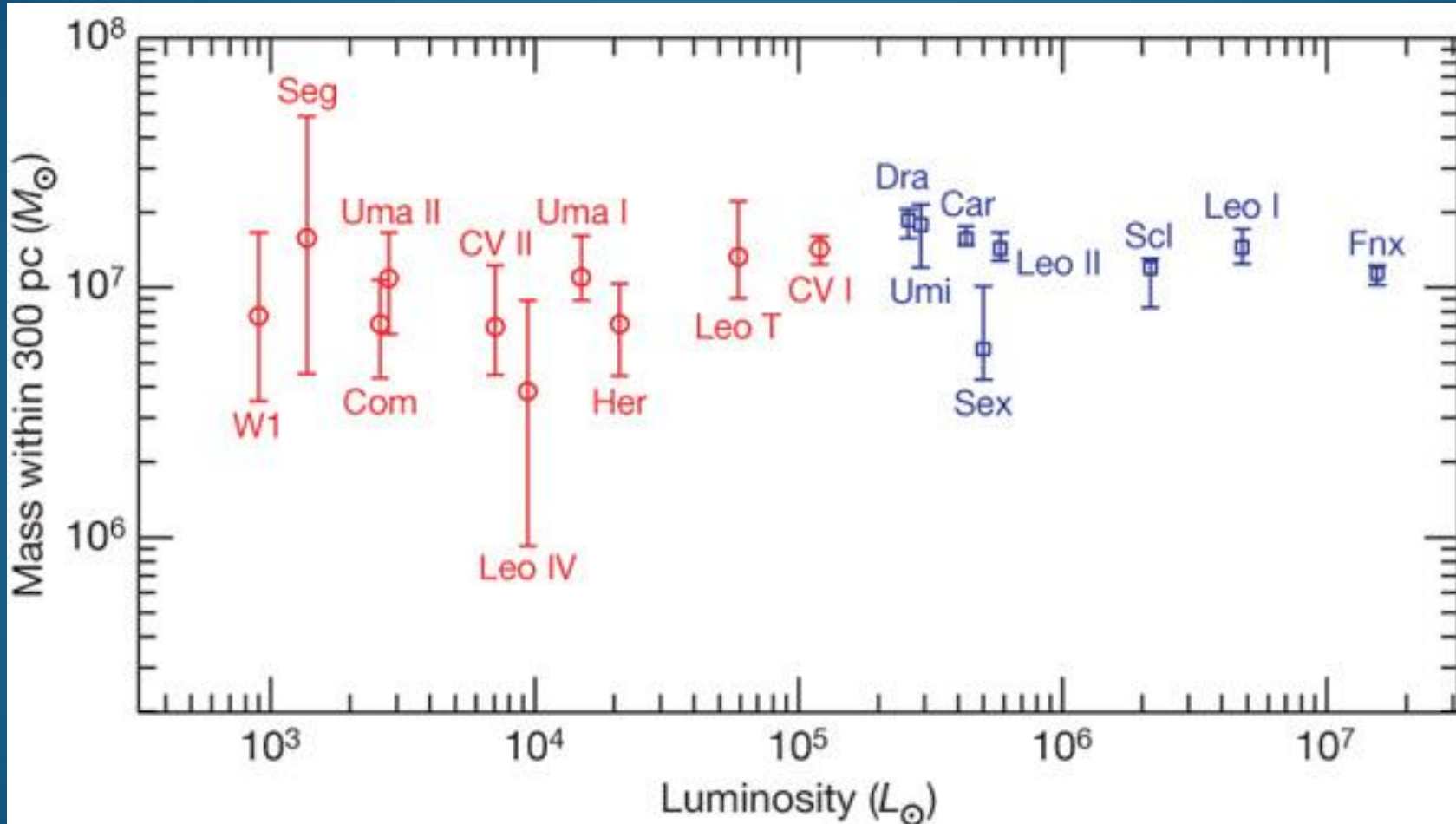
1. A new mass estimator: accurate without knowledge of anisotropy/ β
2. Applications of new mass determinations for MW dSphs
3. Comparison between MW and M_{31} dSphs
4. The skinny on slope determinations



Applications: dSphs

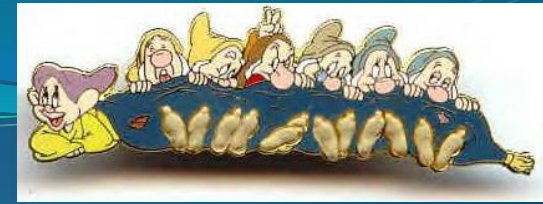


A common mass scale? $M(<300) \sim 10^7 M_{\text{sun}} \rightarrow M_{\text{halo}} \sim 10^9 M_{\text{sun}}$

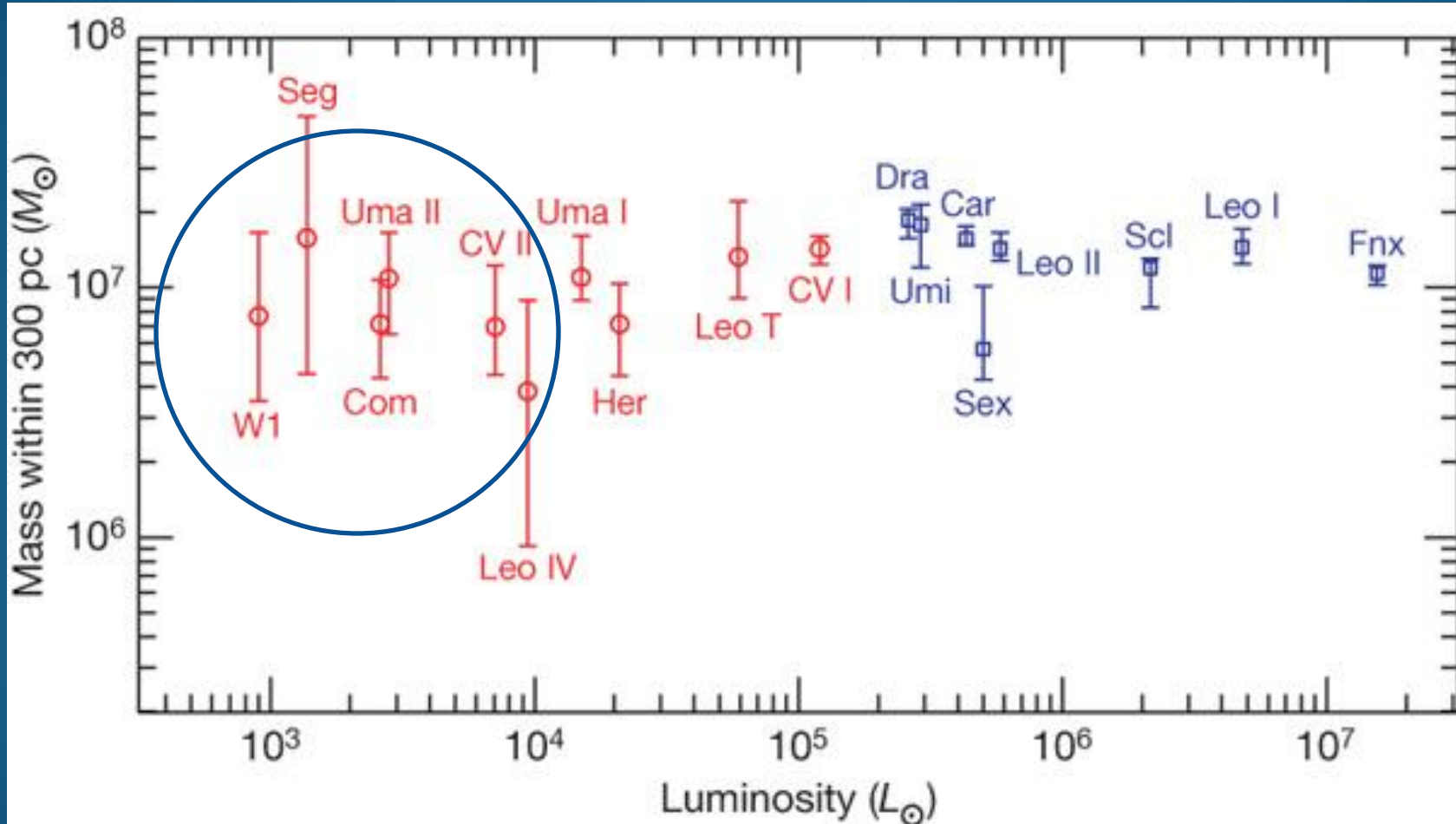


Strigari, Bullock, Kaplinghat, Simon, Geha, Willman, Walker 2008, Nature

Applications: dSphs

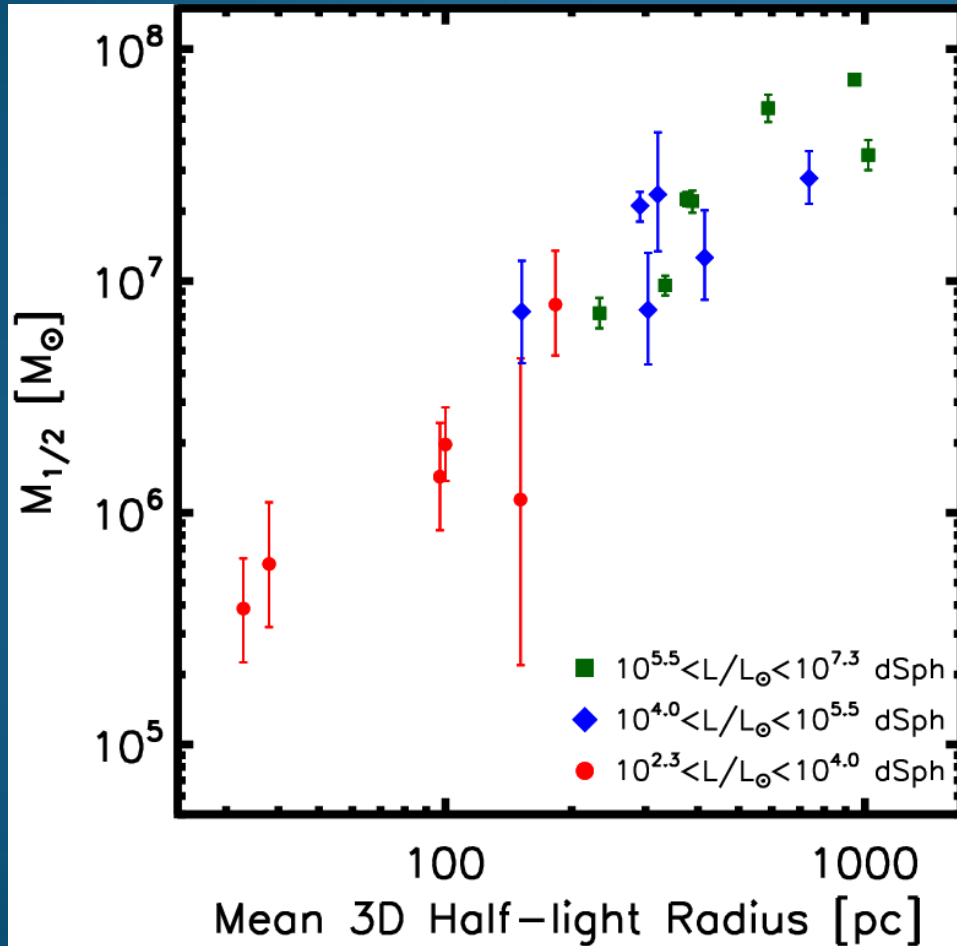


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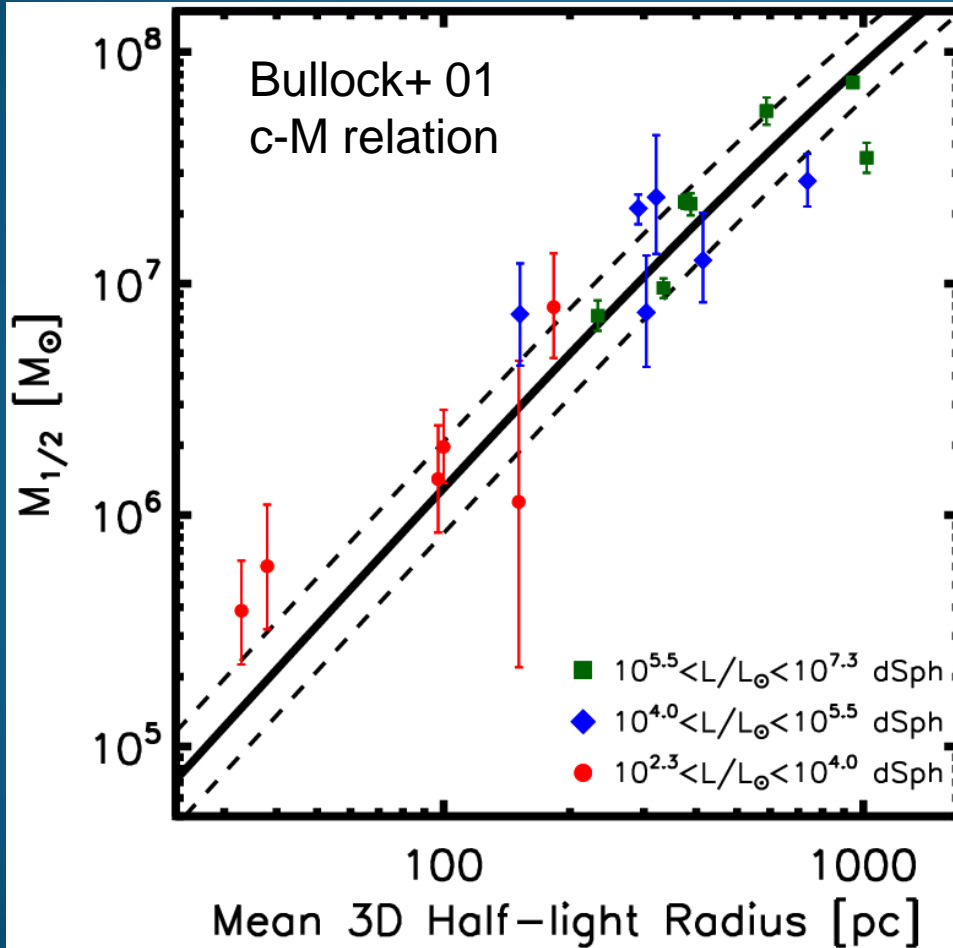
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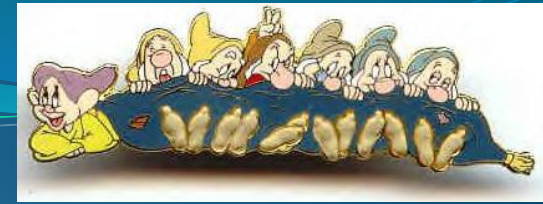
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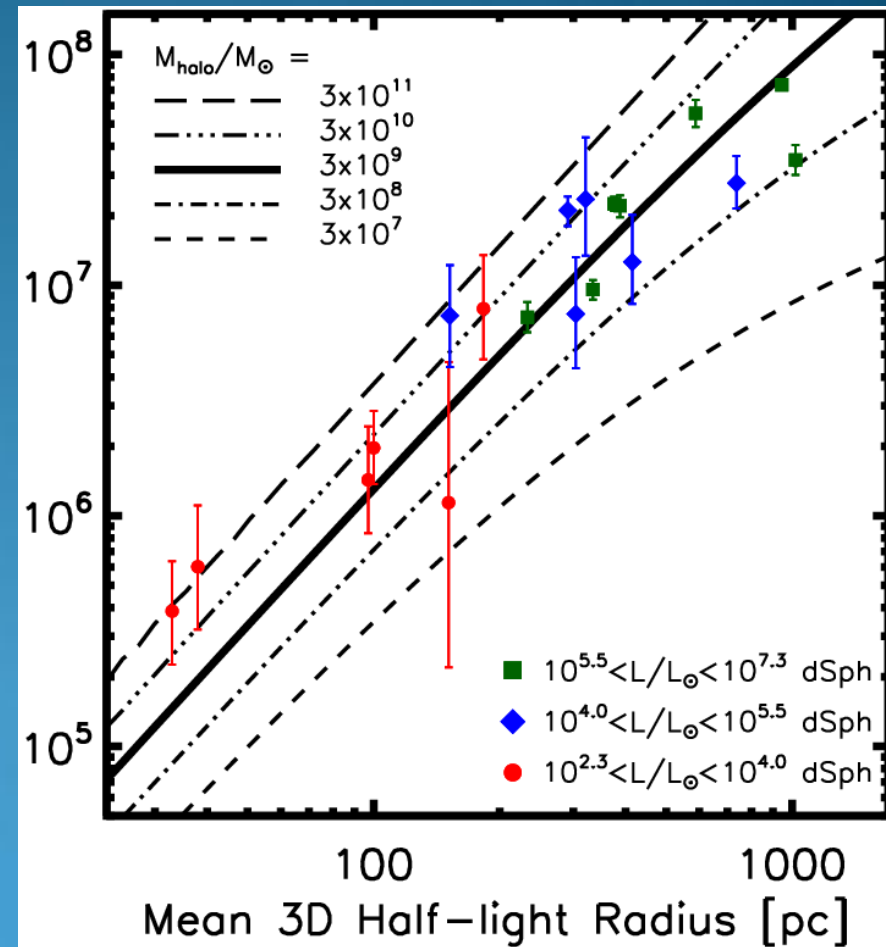
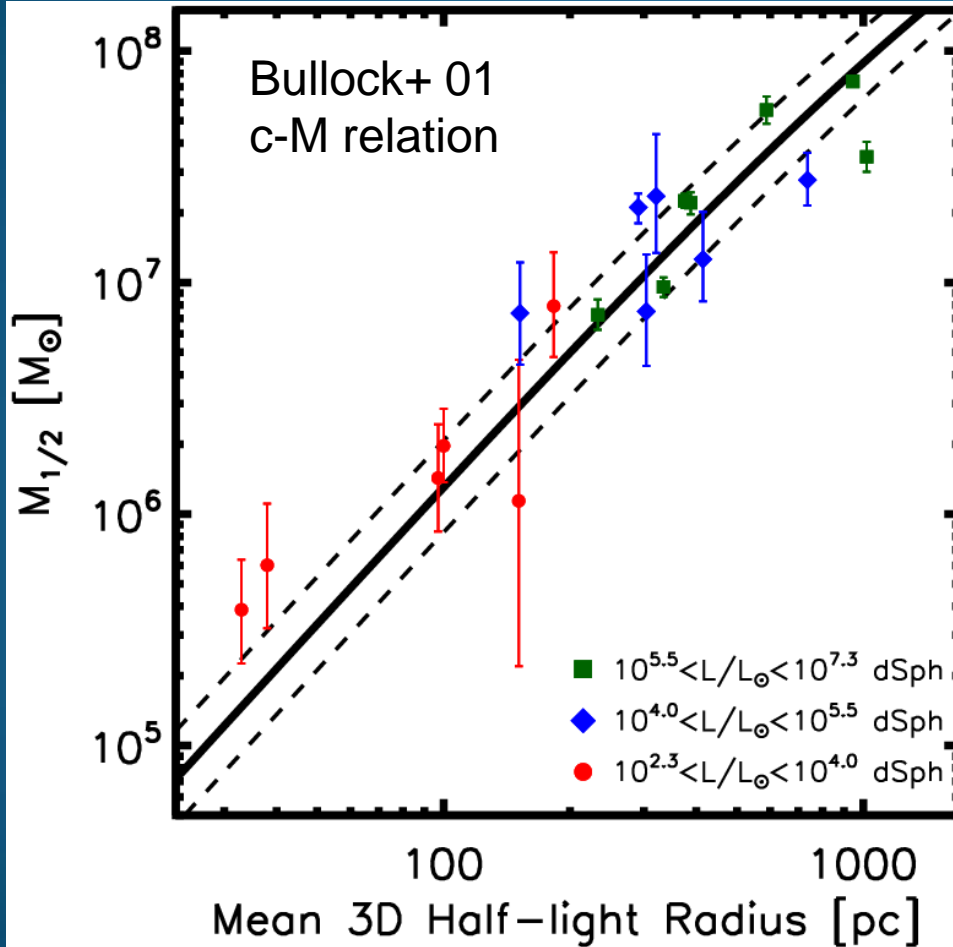
A common mass scale? Plotted: $M_{\text{halo}} = 3 \times 10^9 M_{\text{sun}}$



Applications: dSphs



A common mass scale? Plotted: $M_{\text{halo}} = 3 \times 10^9 M_{\text{sun}}$
Minimum mass threshold for galaxy formation?



Notice: No trend with luminosity, as might be expected! Joe Wolf et al., in prep

Another dataset: M31

UC Irvine: James Bullock, Manoj Kaplinghat, Erik Tollerud, Joe Wolf, Basilio Yniguez

UC Santa Cruz: Raja Guhathakurta (SPLASH PI)

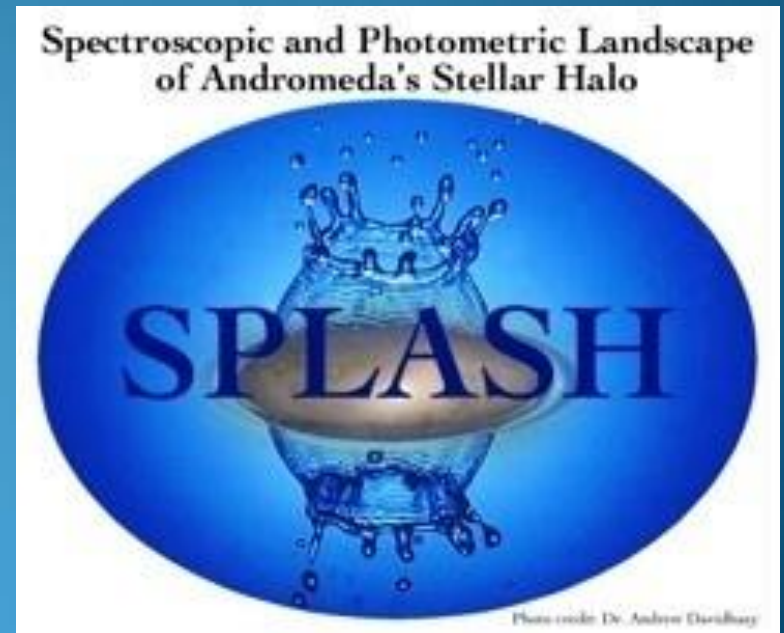
STScI: Jason Kalirai

Yale: Marla Geha

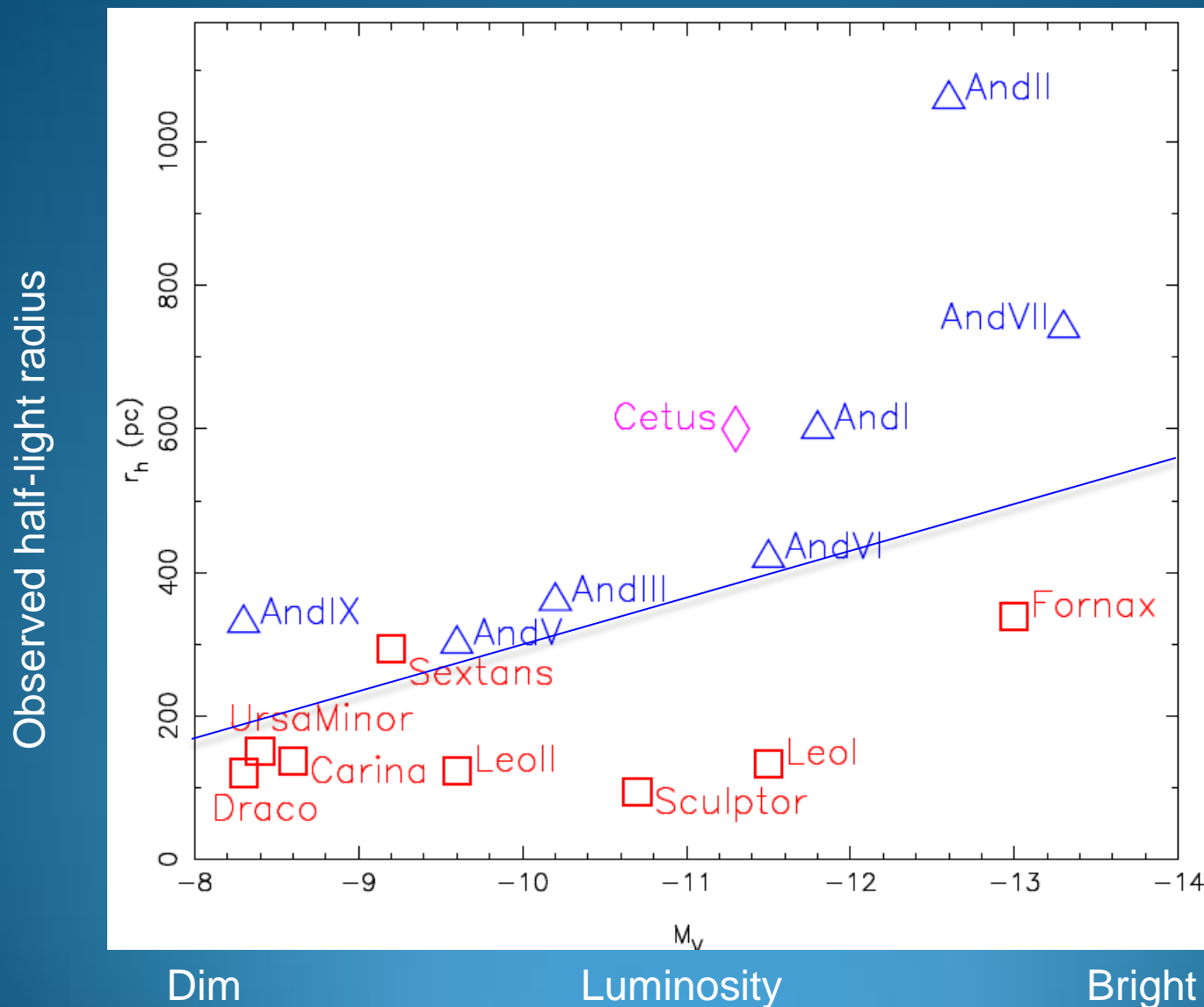
U. Washington: Karrie Gilbert

Caltech: Evan Kirby

And others involved in SPLASH →

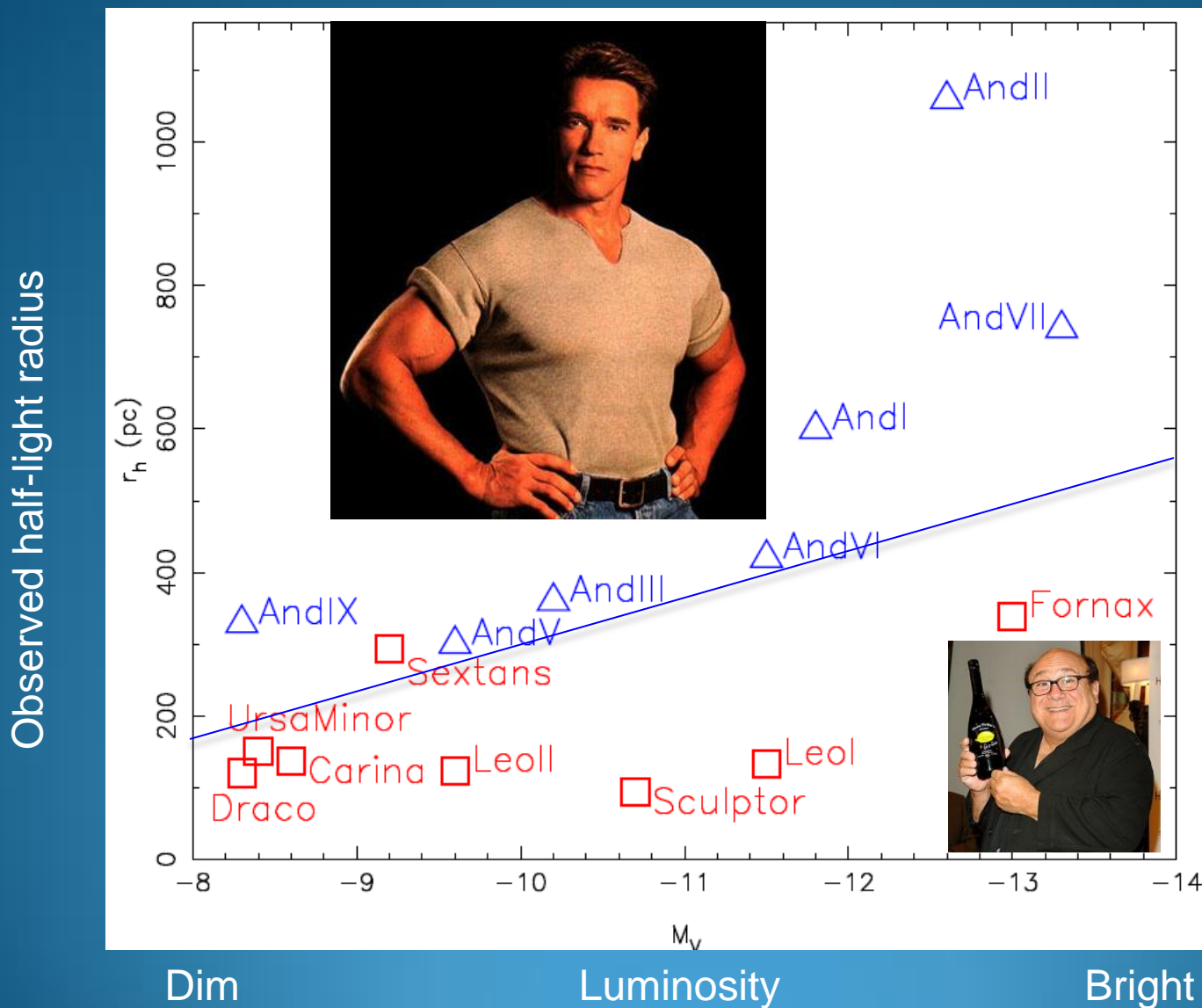


M31 dSphs: Larger than MW dSphs



McConnachie & Irwin 2006, MNRAS

M31 dSphs: Larger than MW dSphs

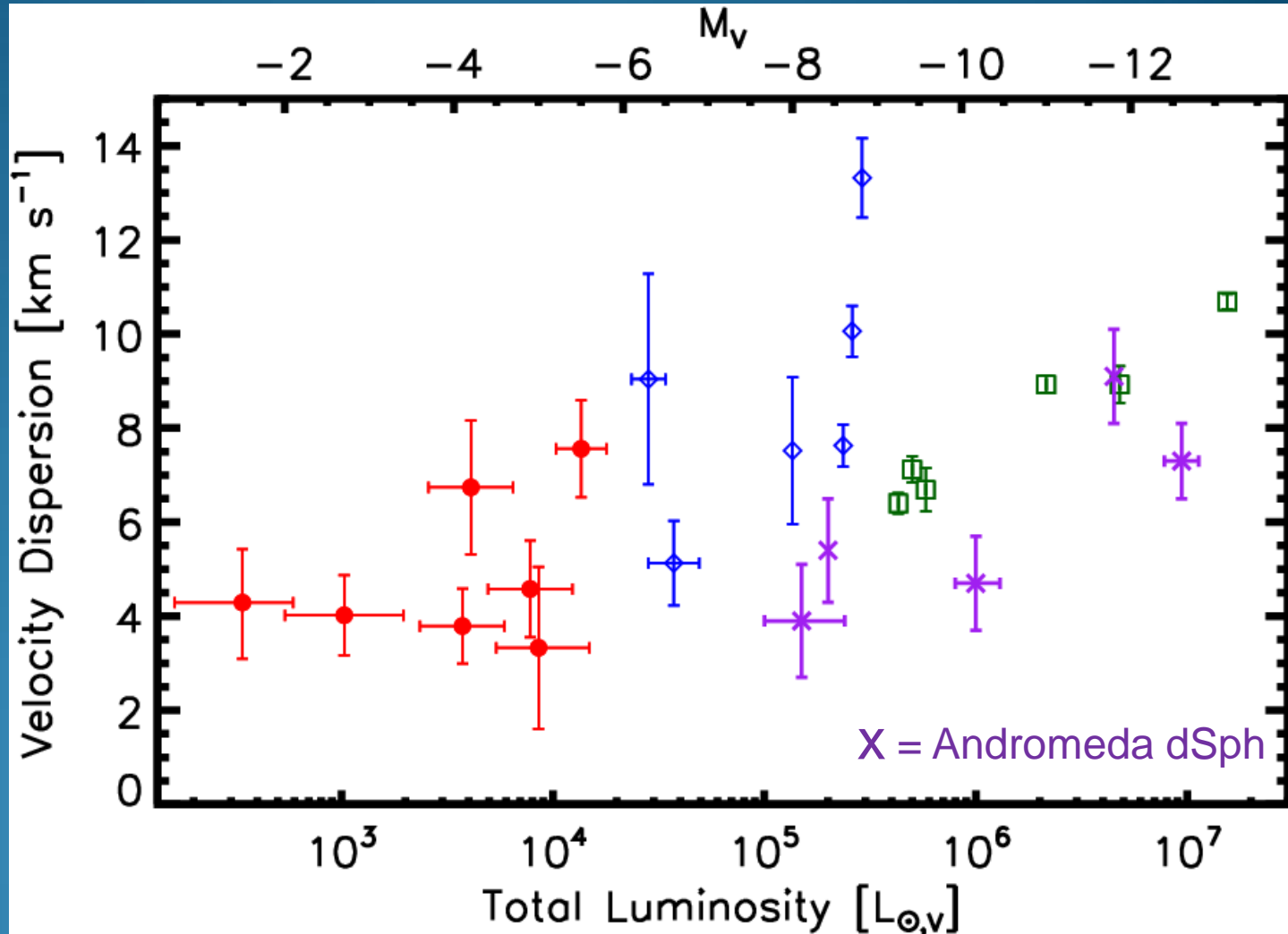


McConnachie & Irwin 2006, MNRAS

Dispersion vs Luminosity

Keck/DEIMOS

And	#	σ km/s
I	76	9.1 ± 1.0
II	95	7.3 ± 0.8
III	43	4.7 ± 1.0
X	22	3.9 ± 1.2
XIV	38	5.4 ± 1.1



Dispersion data from Kalirai et al 2009, in prep

M31 dSphs:

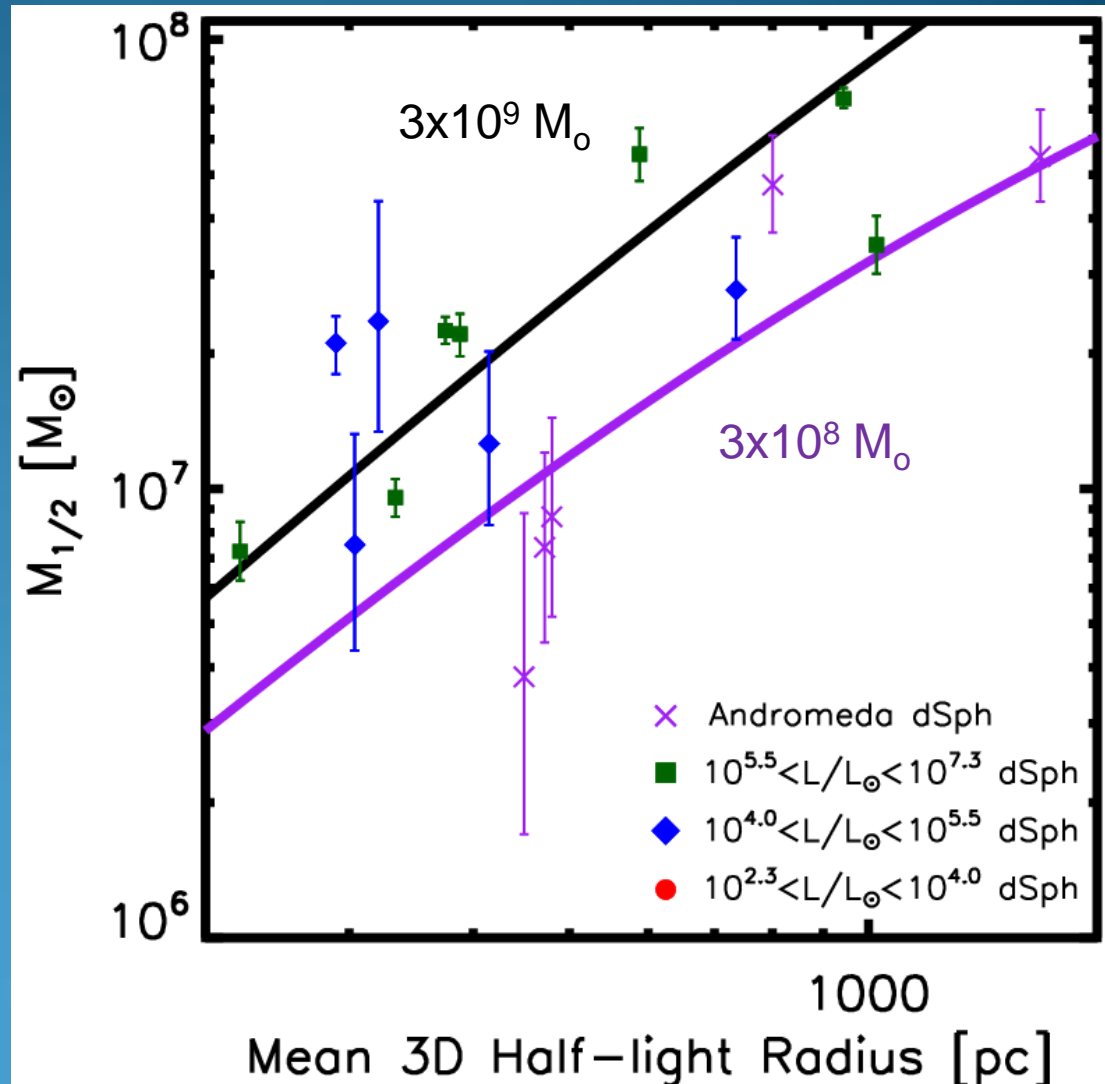
Bigger but less massive!

Spectroscopic data from Keck/DEIMOS.

DM halo mass offset by ~ 10 .
 $M(<300 \text{ pc})$ offset by ~ 2 .



Joe Wolf et al., in prep



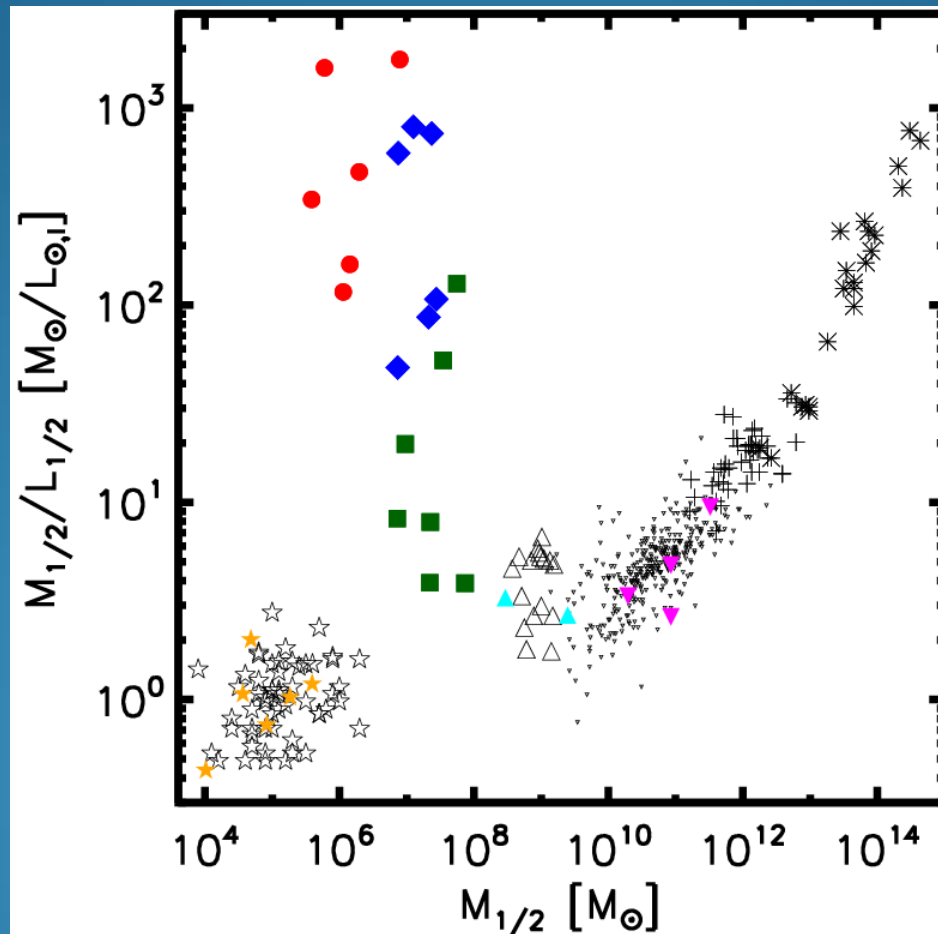
M31: Different Environment?

If M₃₁'s DM halo collapsed later → Less dense substructure & later forming star formation.

Interesting:

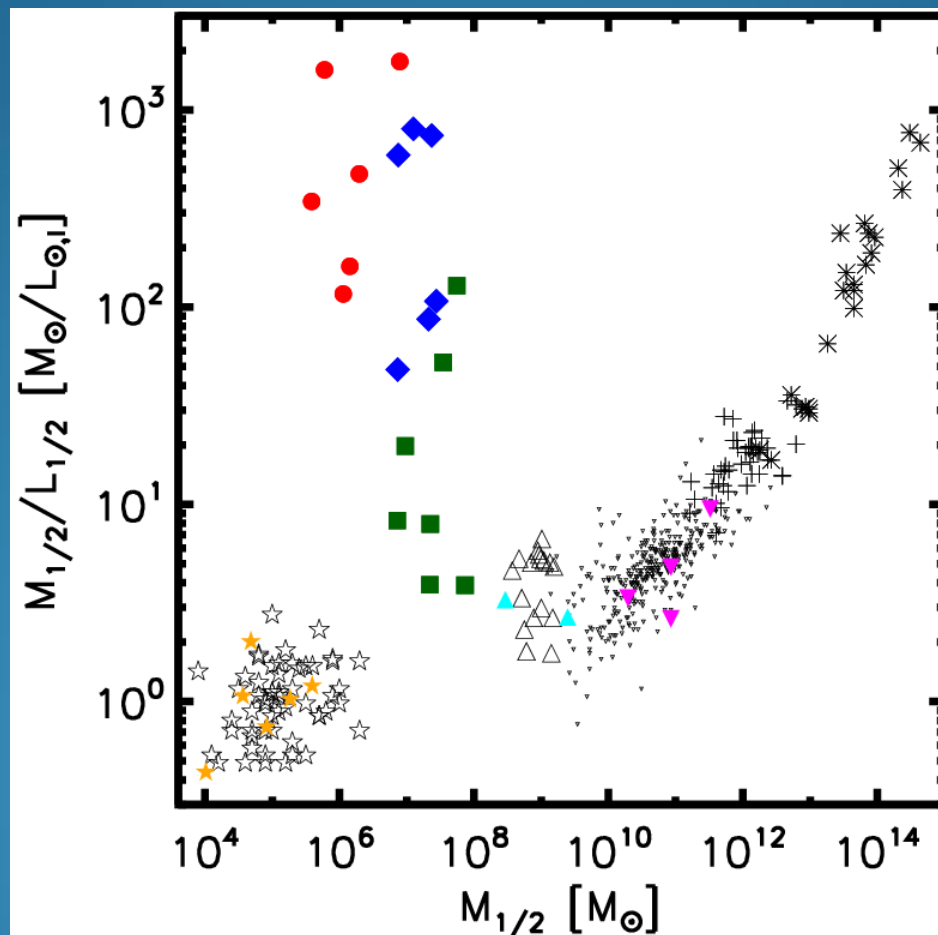
Brown et al. 2008 find that portion of investigated M₃₁ stellar halo is younger (on average) than MW's.

Applications: Global



Applications: Global

Much information about feedback & galaxy formation can be summarized with this plot. Also note similar trend to number abundance matching.



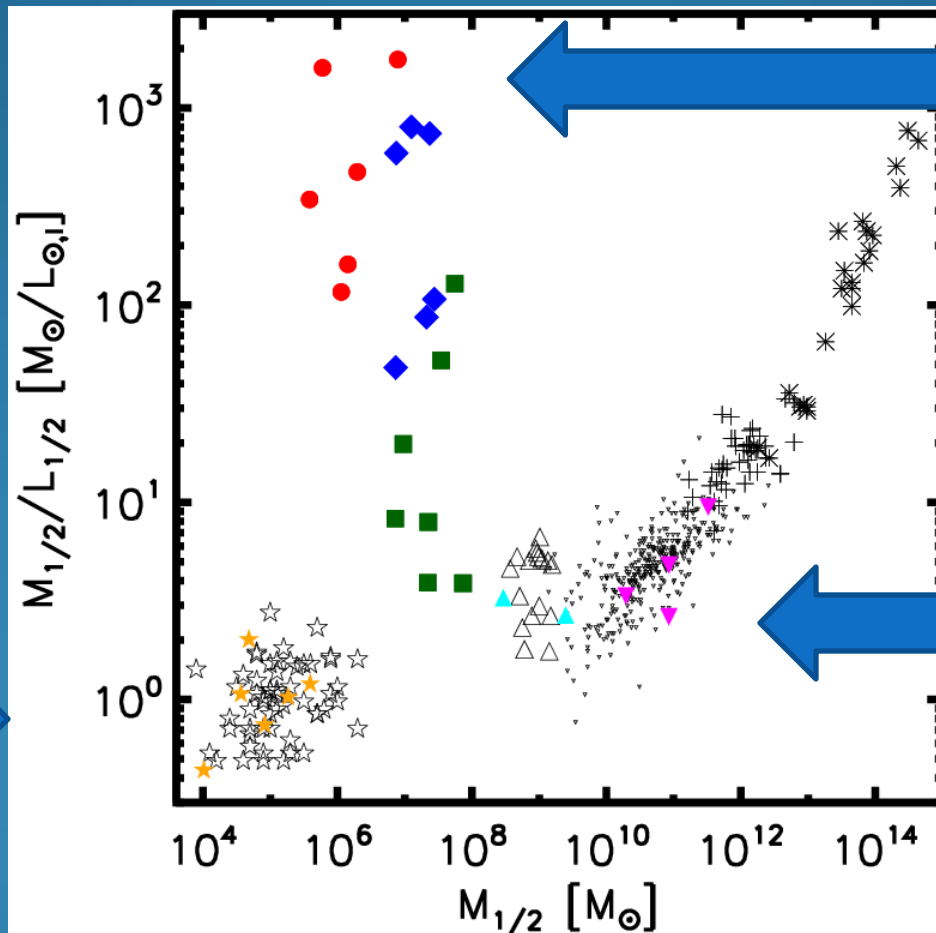
Applications: Global

Much information about feedback & galaxy formation can be summarized with this plot. Also note similar trend to number abundance matching.

Ultrafaint dSphs:
Most DM
dominated
systems known!

Globulars:
Offset from L^*
by factor of
three

(What the
\$%#*?!)

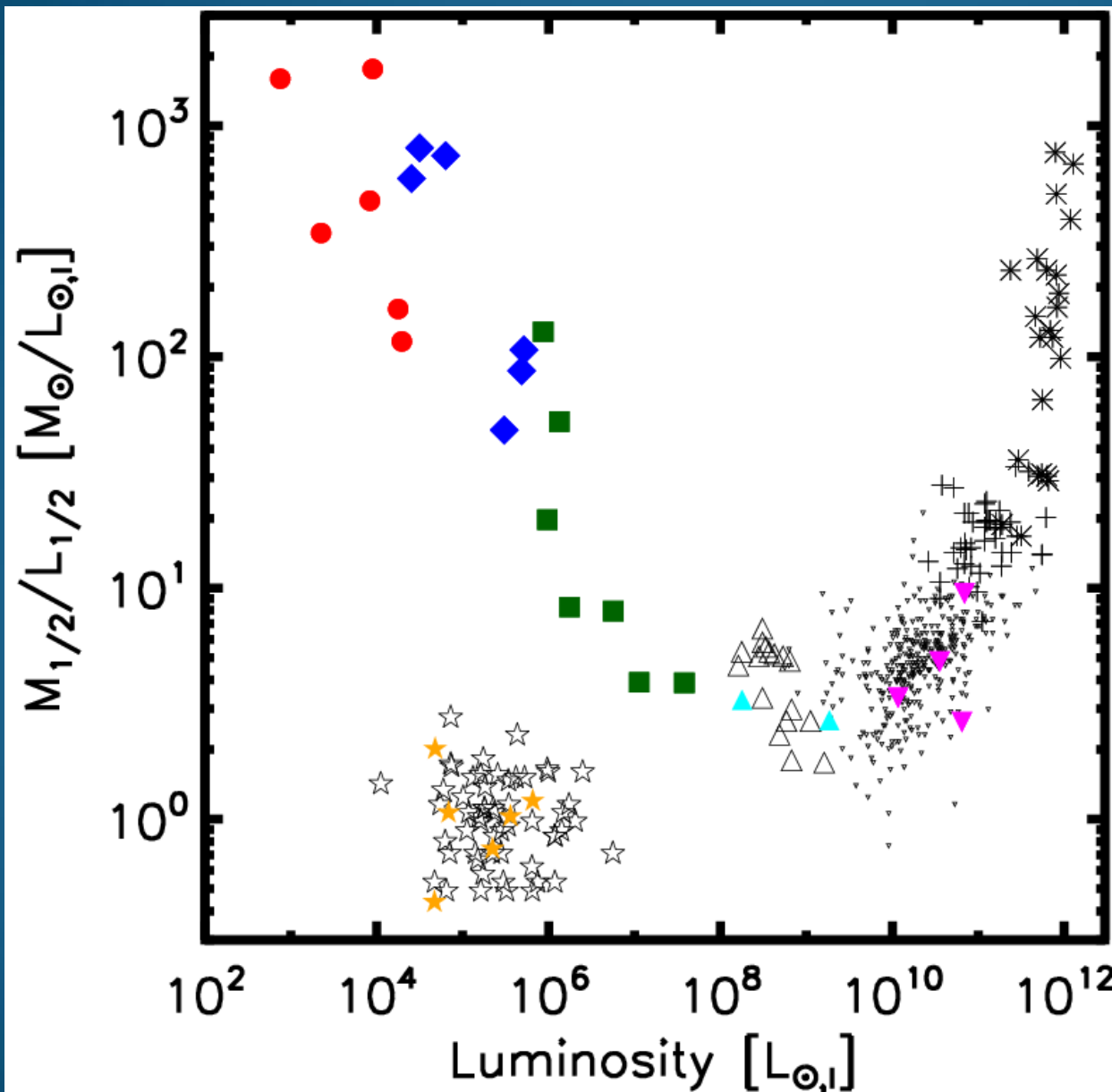


Inefficient at
galaxy formation

L^* : Efficient at
galaxy
formation

Joe Wolf et al., in prep

Applications: Global



Last plot:
Mass floor

This plot:
Luminosity ceiling

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2. ~~Applications of new mass determinations for MW dSphs~~
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4. The skinny on slope determinations



Slopes & Priors – Oh my!

“Can the observed or potentially measurable velocity dispersions tell apart a cusp vs. a core in their centers?” – Conference Website

Slopes & Priors – Oh my!

“Can the observed or potentially measurable velocity dispersions tell apart a cusp vs. a core in their centers?” – Conference Website

No.

Slopes & Priors – Oh my!

“Can the observed or potentially measurable velocity dispersions tell apart a cusp vs. a core in their centers?” – Conference Website

No.
(At least not with LOS kinematics alone.)



Slopes & Priors – Oh my!

An e-mail from my advisor when I told him what I was planning on presenting:

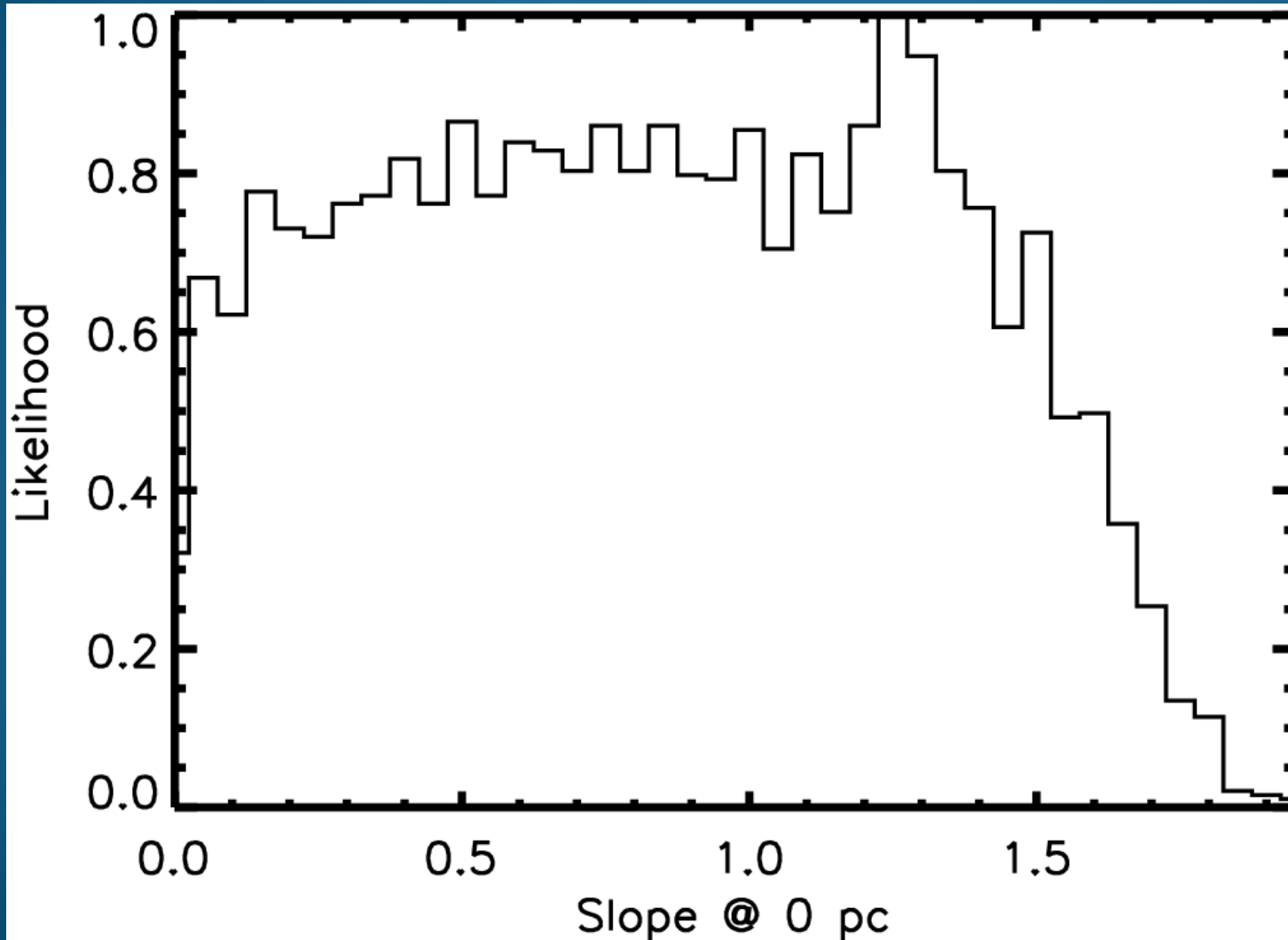
★👤 James Bullock to me

What radius are you quoting the log slope at? $r^{1/2}$ I hope. Please don't tell me you're using gamma.

Slopes & Priors – Oh my!

Beta prior #1: Constant beta that is flat from -10 to 0.91.

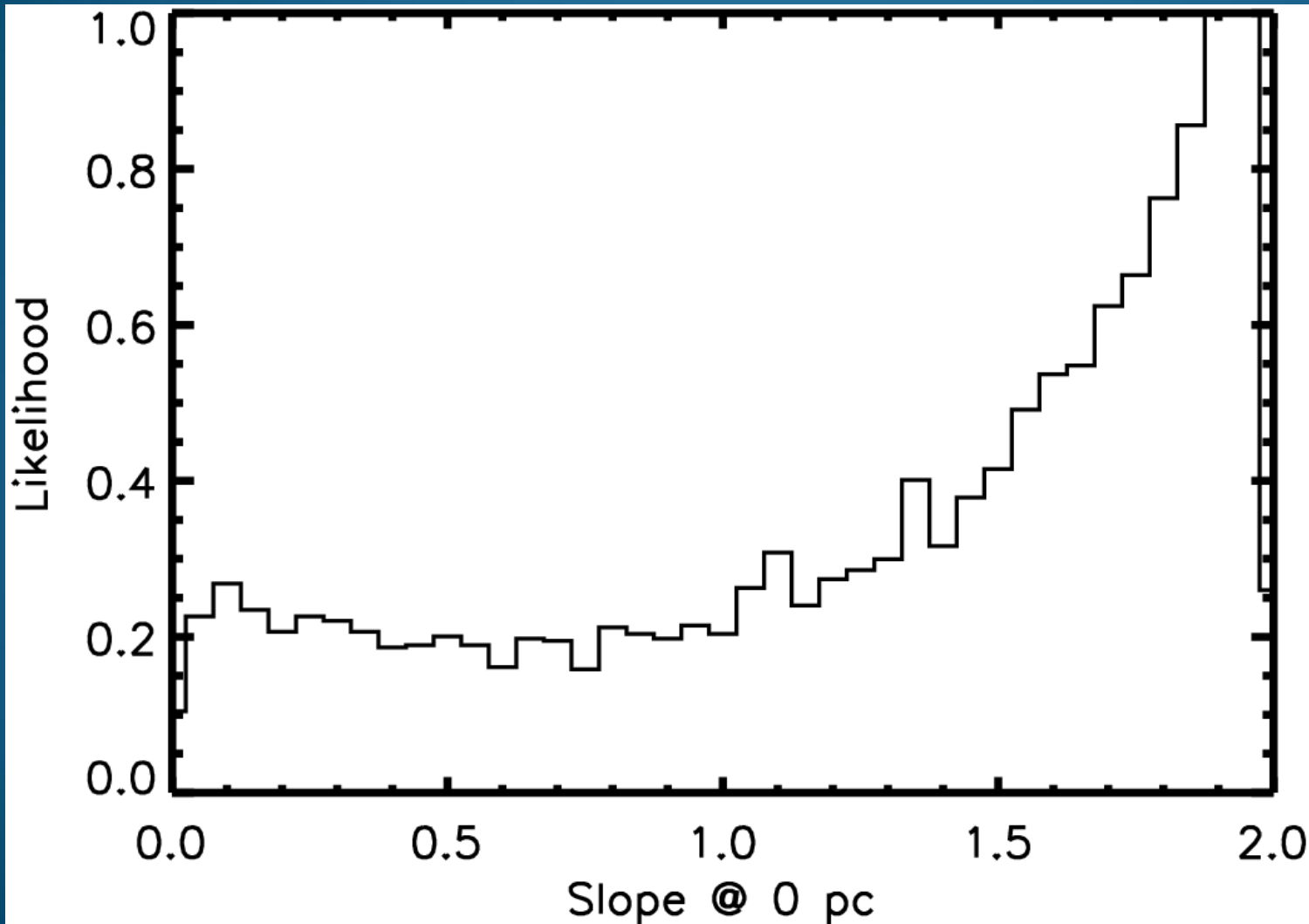
Gamma = Log slope of Carina at 0 pc



Joe Wolf et al.,
in prep

Slopes & Priors – Oh my!

Beta prior #2: Constant beta that is as likely to be negative as positive (ranging from -10 to 0.91).



Joe Wolf et al.,
in prep



Story Time

The next two slides are copied directly from G. Gilmore's 2007 Ann Arbor presentation (slides 14 and 15)



Story Time

The next two slides are copied directly from G. Gilmore's 2007 Ann Arbor presentation (slides 14 and 15)

-
-
-

Except for his pink backgrounds

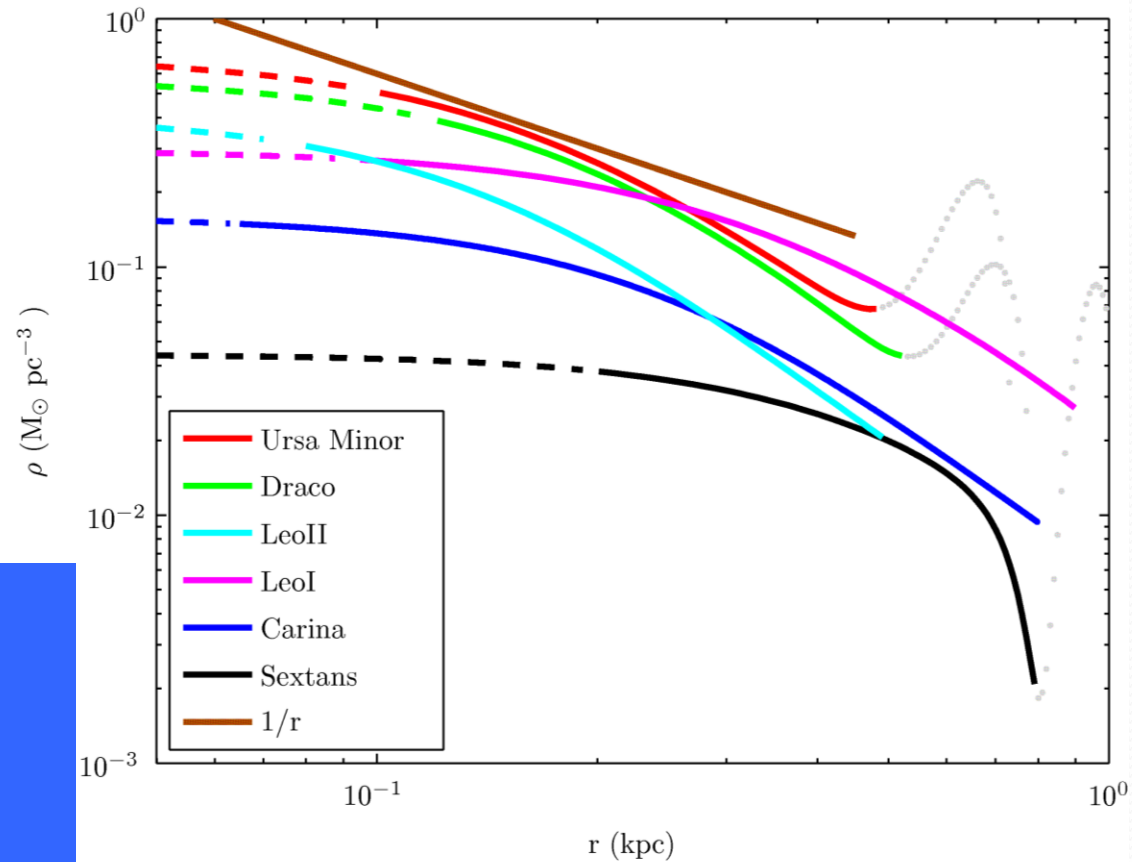


Derived mass density profiles:

Jeans' equation with assumed isotropic velocity dispersion:
All consistent with

cores (similar results from other analyses)

CDM predicts slope of -1.3 at 1% of virial radius and asymptotes to -1 (Diemand et al. 04)



Need different technique at large radii, e.g. full velocity distribution function modelling..
And understand tides.

Conclusion two:

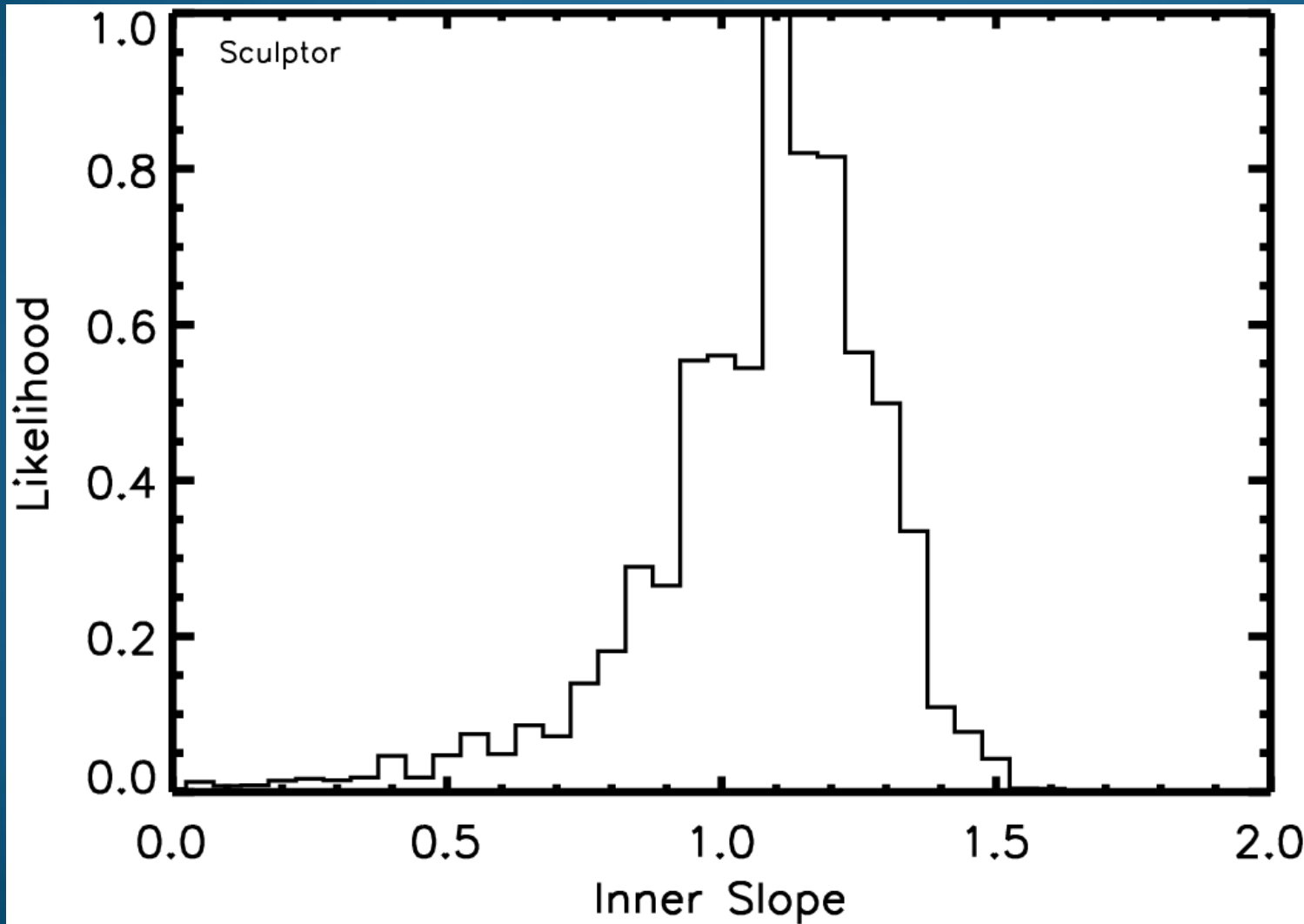
- High-quality kinematic data exist
- Jeans' analyses → prefers cored mass profiles
- Mass-anisotropy degeneracy *allows* cusps
- Substructure, dynamical friction → prefers cores
- Equilibrium assumption is valid inside optical radius
- More sophisticated DF analyses underway

- **Cores always preferred, but not always required**
- **Central densities always similar and low**
- **Consistent results from available DF analyses**

- Extending analysis to lower luminosity systems difficult due to small number of stars
- Integrate mass profile to enclosed mass:

Story Time: A New Ending

Forcing isotropy: 4 of the 8 classical dSphs show no preference for either cores or cusps, and Sculptor strongly prefers a cusp



Joe Wolf et al.,
in prep

- 
- **When assuming isotropy, “cores always preferred”**

- When assuming isotropic pores is “always preferred”

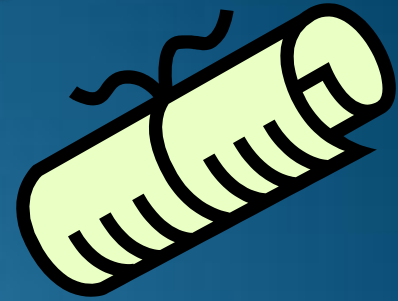


Take-Home Messages



$$M_{1/2} = 3 G^{-1} r_{1/2} \langle \sigma_{\text{los}}^2 \rangle$$

$$\frac{M_{1/2}}{M_{\odot}} \simeq 930 \frac{R_{\text{eff}}}{\text{pc}} \frac{\langle \sigma_{\text{los}}^2 \rangle}{\text{km}^2 \text{s}^{-2}}$$



- M₃₁ dSphs: Offset mass scale. What the *&%#?!
- Knowing M_{1/2} accurately without knowledge of anisotropy gives new constraints for galaxy formation theories to match
- Future simulations must be able to reproduce these results
- GCs vs L* M/L ratios...hmm?
- Inner slopes of dSphs cannot be determined with only LOS kinematics.
- Jeans modeling w/ isotropy does not *always* prefer cores

