

Non-perturbative Simulation of Resonant Modes*

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OUTLINE

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- Motivation to develop the nonperturbative numerical code
- Experimental motivation
- Numerical procedure
- Results and comparison with theory

Introduction

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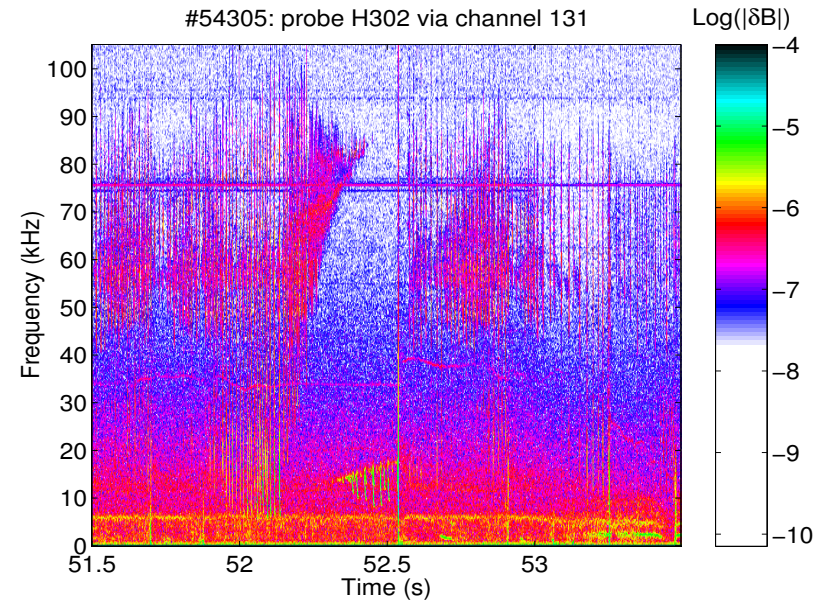
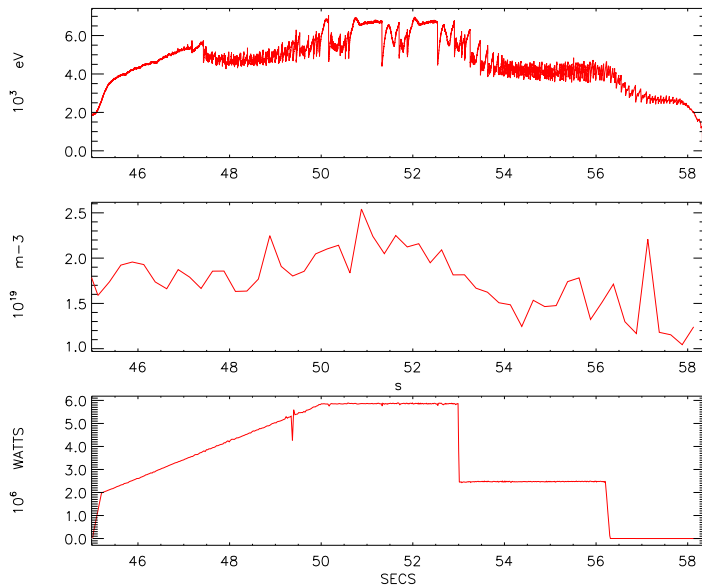
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- Resonant (or Energetic Particle (EPM)) Modes are often observed in tokamak experiments in the presence of hot energetic ions (fishbones, BAEs, R-TAEs).
- New types of low frequency $n = 1$ MHD activity were observed in JET discharges with high fast-ion energy content (Mantsinen00).
- New realistic numerical tools are required to address resonant, non-perturbative modes:
 - such as NOVA-KN (Kinetic Non-perturbative) hybrid code we are developing and apply
 - Carefull phase space integration to evaluate effect of ω_d on the resonant mode drive is required.
 - Important to know linear dispersion relation for perturbative nonlinear theory application (such as Berk-Breizman TAE saturation theory).
 - Complex plasma geometries require numerical treatment of particle-wave interaction.
 - Continuum mechanism may play major role in damping of some modes.

Experimental Observations in JET discharge #54305

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Two types of plasma oscillations are seen:

- 60 – 80 kHz and 10 – 20 kHz (around $t \simeq 52.5$ sec) with $n = 1$.
- Hot ion (ICRH H-minority) beta is large $\beta_H \sim 2\% \sim \beta_{pl}$.
- Both are chirping frequency - “fishbones”?

Numerical procedure in NOVA-KN

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NOVA-KN - hybrid code, integro-differential problem:

Second order ideal MHD + hot ions, coupled through perturbed pressure (Cheng, '92, see also next talk by C.Z.Cheng):

$$\delta p_{\perp} = m_h \int d\mathbf{v} \mu B \hat{g}, \quad \delta p_{\parallel} = m_h \int d\mathbf{v} 2 (\mathcal{E} - \mu B) \hat{g}$$

$\hat{g}$ - nonadiabatic perturbed distribution function (Gorelenkov, '99):

$$\hat{g} = \mathcal{E} \sum_{m,p} \frac{(\omega - \omega_*) G_{m,p}}{\omega - \bar{\omega}_d - p\omega_b} \frac{\partial F_h}{\partial \mathcal{E}} e^{-it\omega + im\theta - in\varphi}, \quad (1)$$

$G_{m,p}$ is the matrix of wave-particle interaction (Cheng '92).

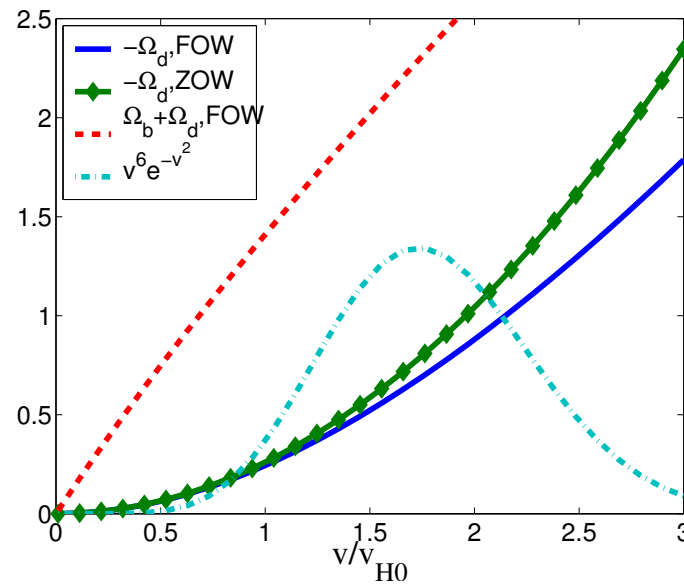
- Present version uses ZOW for G matrix.
- Accurate estimate of drift $\bar{\omega}_d$ and bounce ω_b hot ion frequencies in FOW.
- NOVA-KN solves integro-differential equations and finds mode frequency in complex plane.

Characteristic hot ion drift motion frequencies

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Consider deeply trapped hot ions at $r/a = 1/3$ and $E_{H0} = 800keV$ (agrees with ORBIT).



Characteristic frequencies are determined by the ions with $v/v_{H0} \simeq 1.7$ ($E_{H0} = 2.3MeV$), where the perturbed pressure integrand is peaked $v^6 e^{-v^2}$. $\Rightarrow$

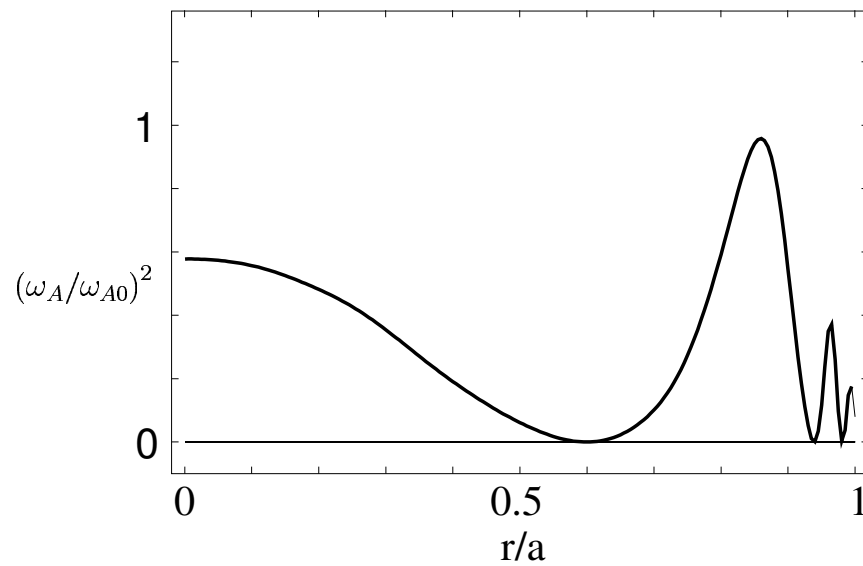
1. $\Omega_d \equiv \omega_d/\omega_{A0} \sim 0.7$ (~ 0.35 for ICRH H-minority, $\sim 60kHz$) and
2. $\Omega_b \equiv \omega_b/\omega_{A0} \sim 2$, ($\omega_{A0}/2\pi = 160kHz$).

Alfvén continuum

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$q = 1$ surface is at $r/a \equiv \sqrt{\psi/\psi_0} = 0.6$.



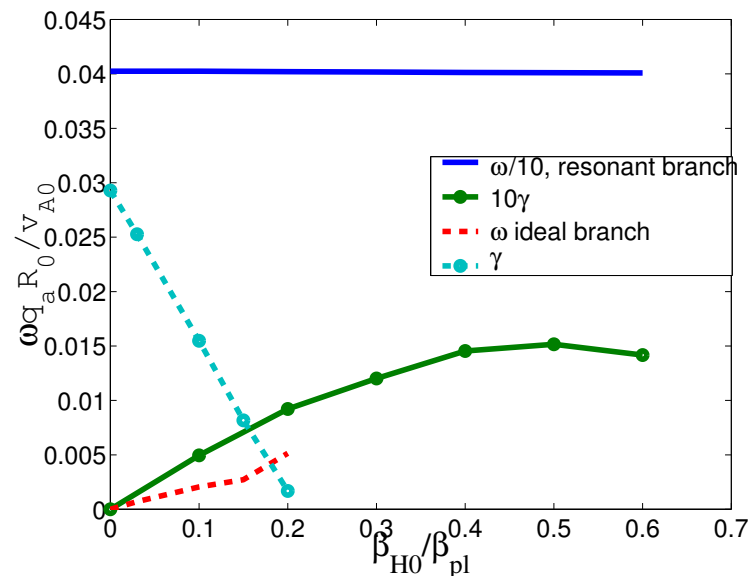
- Any mode with non-zero frequency is reflected from the continuum.
- Drive should be strong enough to overcome the continuum damping.

Two branches are found with NOVA-KN

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- **Lower frequency** mode is a continuation of the ideal mode.
It is strongly stabilized by hot ions at $\beta_{H0}/\beta_{pl} > 0.2$.
- **Higher frequency** mode (higher frequency fishbone -resonant mode) has frequency comparable to the measured at H-minority temperature $800keV$.



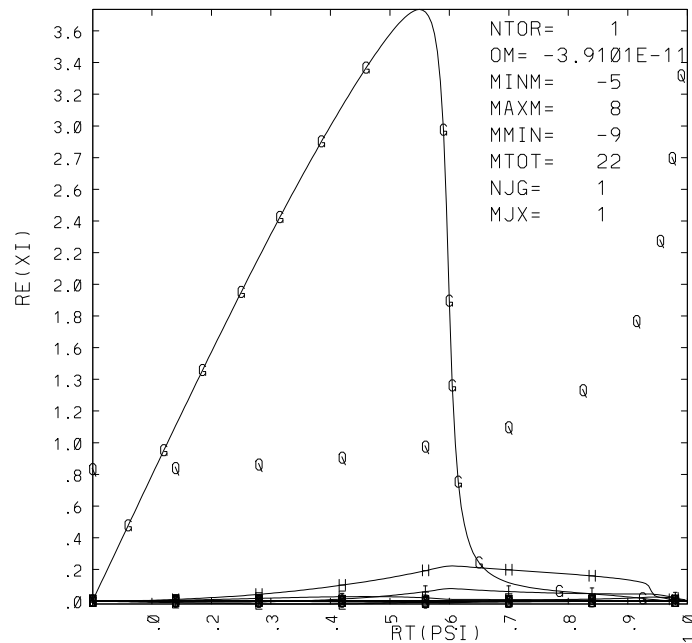
Ideal and Resonant mode radial structures

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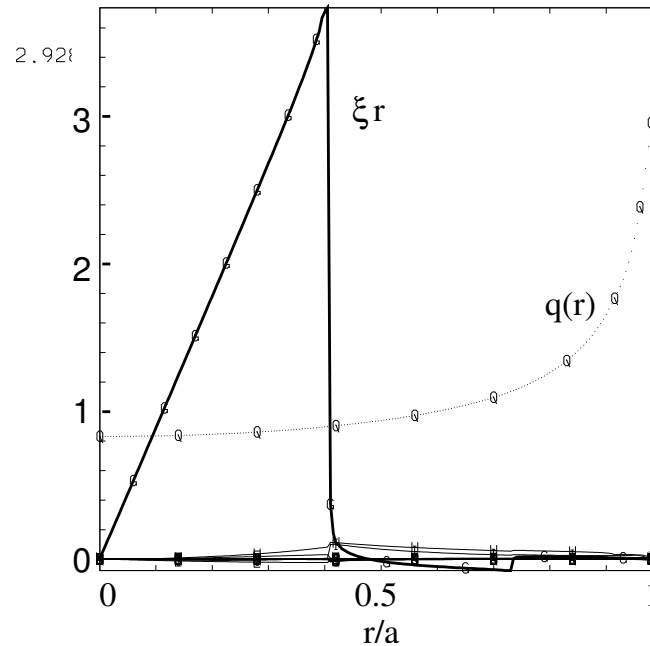
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Shown are the radial structures of the dominant poloidal harmonics ($m = 1$).

Ideal mode



Resonant mode



- Ideal mode is strongly driven $\Rightarrow$ mode structure is effected.
- Resonant mode is reflected from the continuum due to finite frequency.
 - Predicted growth rate is probably too low for such nonperturbative mode.

Comparison with theory

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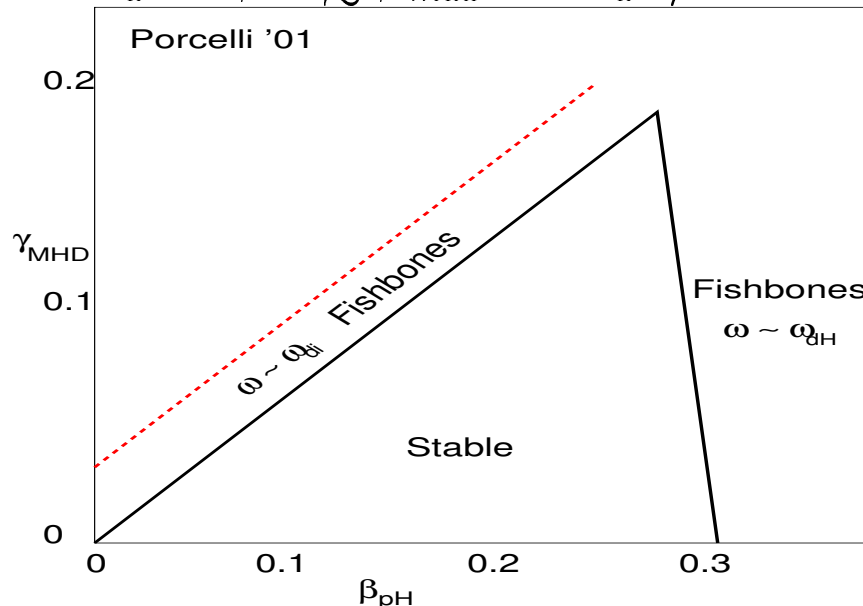
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Two types (branches) of **fishbones** are known Chen, et.al. '84, Coppi, et.al. '86, '88, Porcelli '01: $i [\omega (\omega - \omega_{di})]^{1/2} = \lambda_{mhd} + \lambda_H$.

- Low frequency $\omega \sim \omega_{di}$, if $\omega_{di} \ll \omega_{dH} (\sim 1:250)$ and

$$\beta_H \lesssim \beta_{min} \simeq \pi \left[\beta^2 - (\beta^{MHD})^2 \right]$$

- High frequency $\omega \simeq \omega_{dH}$, if $\beta_H \gtrsim \beta_{max} \simeq \epsilon \omega_{dH} / \omega_A \pi$.



NOVA-KN predicts $\omega / \omega_{dH} \sim 1/100$ for low frequency branch.

H-minority distribution function (factorizing?) may be a key factor for ω_I .

Summary

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- A non-perturbative code NOVA-KN is being developed to account for realistic calculations of the non-perturbative resonant modes with FOW effects. It is fast and useful tool to study physics of particle-wave interaction.
- More new physics needs to be included in computations such as: nonisotropic distribution function of hot ions, nonisotropic plasma equilibrium, radial coupling due to FOW.
- Comparison with JET experiments shows that the possible explanation for the $10 - 70 kHz$ frequency range MHD $n = 1$ oscillations is the excitation of the two types of modes.
 - At low frequency $f = 10 - 20 kHz$ the unstable mode is hot ion modified ideal MHD mode excited when beta of hot ions is not so strong, but frequency in experiment is higher.
 - At higher frequencies $f = 50 - 70 kHz$ NOVA-KN predicts instability of the resonant modes having frequency equal to characteristic toroidal precession frequency of H-minority ions.

Summary (continued)

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- Expected hot ion energy is 800keV in the discharge studied, which is slightly higher than typically observed 500keV , but agrees with the preliminary results by V. Kiptily.