

## Aspiration of high

## energies and intense lasers

APSA, University of Tokyo Tokyo, Japan Oct. 23, 2014

## T. Tajima, UCI, KEK

Acknowledgments for Collaboration: G. Mourou, N. Naumova, K. Nakajima, Y.M. Shin, S. Bulanov, A. Suzuki, T. Ebisuzaki, J. Koga, X. Q. Yan, U. Wienands, U. Uggerhoj, A. Chao, N.V. Zamfir, V. Shiltsev, M. Hogan, K. Ishikawa, Y. Tobita, M. Spiro, R. Kumar, N. Rostoker, Y. Kato, the late Professor H. Takuma



Eur. Phys. J. Spec. Top. 228, 1037 (2014)

## Content

- High intensity frontier of lasers I = E/τ: (A) large energy;
  (B)ultrashort; (C) high fluence (P = fE)
- Laser can accelerate particles: laser wakefield
- (A) Large energy laser: Higgs energy (>100GeV) over 10m
- (B) Ultrashort (zeptosecond) X-ray laser : TeV on a chip
  (cm)→ zeptoscience (τ = E/I)
- Nature finds wakefield acceleration: blackhole jets

 $\rightarrow$  ultrahigh energy cosmic rays (ZeV)

• (C) High fluence laser: CAN (Coherent Amplification Network) laser

## **Relativistic nonlinearity under intense laser**

Plasma free of binding potential, but its electron responses:

a) Classical optics : v << c,  $a_0 << l: \delta x$  only

b) Relativistic optics:  $v \sim c$  $a_0 >> 1: \delta z >> \delta x$ 



#### What is *collective force*? :Secret behind laser accelerator

How can a Pyramid have been built?





<u>Individual</u> particle dynamics  $\rightarrow$  <u>Coherent</u> and <u>collective</u> movement

Collective acceleration (Veksler, 1956; Tajima & Dawson, 1979) Collective radiation (N<sup>2</sup> radiation) Collective ionization (N<sup>2</sup> ionization)

 $\rightarrow$  Laser driven collective accelerating field

## Wakefield : a <u>Collective</u> Phenomenon + Relativistic Coherence

VS.

All particles in the medium participate = collective phenomenon



No wave breaks and wake <u>peaks at *v*≈*c*</u>



(electrons cohere; density cusps) → <u>Relativistic Coherence</u> Kelvin wake



#### Wave breaks at v < c





#### No relativistic coherence

Hokusai

## **Thousand-fold Compactification**

*Laser wakefield:* thousand folds gradient (and emittance reduction?)



## **Extreme Light Roadmap**

(modified from Tajima and Mourou, PR 2002)



## Suzuki's Challenge: "When can we reach 1 PeV ?"



V. Yakimenko (BNL) and R. Ischebeck (SLAC), AAC2006 Summary report of WG4



#### Theory of wakefield toward extreme energy



### IZEST proposes 100 GeV Ascent Experiment using 3.5 kJ, 500fs, 7 PW PETAL laser

① Large Energy Laser

K. Nakajima



## Single-Cycle Laser Compressor w/thin film



G. Mourou, S. Mironov, E. Khazanov and A. Sergeev, Eur. Phys. J. Special Topics, **223**, 1181(2014)



N. M. Naumova, J. A. Nees, I. V. Sokolov, B. Hou, and G. A. Mourou, Phys. Rev. Lett. **92**, 063902-1 (2004).

## Even, Isolated zeptosecond X-ray laser pulse possible

(preliminary simulation by N. Naumova, I. Sokolov, G. Mourou, 2014)



→ EW single osc. X-ray laser

## Earlier works of X-ray crystal acceleration

-X-ray optics and fields (Tajima et al. PRL,1987)

-Nanocrystal hole for particle propagation (Newberger, Tajima, et al. 1989, AAC; PR,..) -particle transport in the crystal (Tajima et al. 1990, PA)

#### APPLICATION OF NOVEL MATERIAL IN CRYSTAL ACCELERATOR CONCEPTS

B. Newberger, T. Tajima, The University of Texas at Austin, Austin, Texas 78712

F. R. Huson, W. Mackay, Texas Accelerator Center, The Woodlands, Texas

B. C. Covington, J. R. Payne, Z. G. Zou, Sam Houston State University, Huntsville, Texas

N. K. Mahale, S. Ohnuma, University of Houston, Houston, Texas 77004

which incorporate regular macroscopic features on the underlying crystal lattice are of potential ne application to crystal accelerators and coherent nrces. We have recently begun an investigation of iterial, porous Si, in which pores of radii up to a attice spacings are etched through finite volumes rystal. The potential reduction of losses to partianneled along the pores makes this a very interial in crystal accelerators for relativistic, positively icles. Our results on material properties which are this context will be presented. The consequences ransport will be discussed. and  $k = v_0/m_I c^2$ ,  $v_0$ , is the "spring constant of th channel well. Its specific form depends on the moconstruct the continuum potential of a string of aton purposes it suffices to take a typical value of  $2 \times 10^1$ is the multiple scattering velocity space "diffusion" We have used<sup>10</sup>

$$D = z\pi r_e^2 N Z_{\rm val} \left(\frac{m_e}{m_I}\right)^2 L_R,$$

where  $r_E$  is the classical electron radius,  $Z_{\text{val}}$  is t of valence electrons, and N is the number density of tal. Logarithmic dependencies on particle energy neglected throughout:  $L_P$  is a constant with a ty

Particle Accelerators, 1990, Vol. 32, pp. 235-240 Reprints available directly from the publisher Photocopying permitted by license only © 1990 Gordon and Breach, Science Publishers, Inc. Printed in the United States of America

#### BEAM TRANSPORT IN THE CRYSTAL X-RAY ACCELERATOR

T. TAJIMA, B. S. NEWBERGER University of Texas-Austin, Austin TX 78712 U.S.A. F. R. HUSON, W. W. MACKAY Texas Accelerator Center, The Woodlands, TX 77381 U.S.A. B. C. COVINGTON, J. PAYNE Sam Houston State University, Huntsville, TX 77341 U.S.A. N.K. MAHALE, S. OHNUMA University of Houston, Houston, TX 77204 U.S.A.

<u>Abstract</u> A Fokker-Planck model of charged particle transport in crystal channels which includes the effect of strong accelerating gradients has been developed<sup>1</sup> for application to VOLUME 59, NUMBER 13

PHYSICAL REVIEW LETTERS Crystal X-Ray Accelerator 28 SEPTEMBER 1987

T. Tajima Department of Physics and Institute for Fusion Studies, The University of Texas, Austin, Texas 78712

and

M. Cavenago Department of Physics, University of California, Irvine, California 92717 (Received 18 November 1986)

An ultimate linac structure is realized by an appropriate crystal lattice (superlattice) that serves as a "soft" irised waveguide for x rays. High-energy ( $\approx$ 40 keV) x rays are injected into the crystal at the Bragg angle to cause Bormann anomalous transmission, yielding slow-wave accelerating fields. Particles (e.g., muons) are channeled along the crystal axis.

 $m/(k_{-}+k_{-})$ 

PACS numbers: 52.75.Di, 41.80.-y, 61.80.Mk

An approach to the attainment of ever higher energies by extrapolating the linac to higher accelerating fields, higher frequencies, and finer structures is prompted by several considerations, including the luminosity requirement which demands the radius of the colliding-beam spot be proportionately small at high energies: an  $=\pi^{-1/2}hc(f\mathcal{N})^{-1/2}P\epsilon^{-2}$ , where f, N, P, and  $\epsilon$  are the duty cycle, total number of events, beam power, and beam energy, respectively. This approach, however, encounters a physical barrier when the photon energy becomes of the order  $\hbar \omega \simeq \hbar \omega_{\sigma} \simeq mc^2 a^2 \simeq 30 \text{ eV}$  (a=the fine-structure constant), corresponding to wavelength (scale length)  $\lambda \simeq 500$  Å: The metallic wall begins to absorb the photon strongly, where  $\omega_n$  is the plasma frequency corresponding to the crystal electron density. In addition, since the wall becomes not perfectly conducting for  $\hbar \omega \ge mc^2 a^2$ , the longitudinal component of fields becomes small and the photon goes almost straight into the wall (a soft-wall regime). As the photon energy  $\hbar \omega$ much exceeds  $mc^2a^2$  and becomes  $\geq mc^2a$ , however, the metal now ceases to be opaque. The mean free path of the photon is given by Bethe-Bloch theory as  $l_i = (3/2^8 \pi)$  $\times a_{\rm B}^{-2} \alpha^{-1} n^{-1} (\hbar \omega / Z_{\rm eff}^2 \mathcal{R})^{7/2}$ , where  $a_{\rm B}$  is the Bohr radius, n the electron density,  $Z_{eff}$  the effective charge of the lattice ion, and  $\mathcal{R}$  the Rydberg energy.

In the present concept the photon energy is taken at the hard x-ray range of  $\hbar\omega = mc^2 a$  and the linac structure is replaced by a crystal structure, e.g., silicon or GaAs-AIAs. (A similar bold endeavor was apparently undertaken by Hofstadter already in 1968.<sup>1</sup>) Here the crystal axis provides the channel through which accelerated particles propagate with minimum scattering (channeling<sup>-1</sup>) and the x rays are transmitted via the Bormann effect (anomalous transmisson.<sup>3</sup>) when the x rays (wavelength  $\lambda$ ) are injected in the xz plane with a where b is the transverse lattice constant and later a the longitudinal lattice constant ( $a \simeq b$ ) (see Fig. 1). The row of lattice ions (perhaps with inner-shell electrons) constitutes the "waveguide" wall for x rays, while they also act as periodic rises to generate slow avaes. A superlattice<sup>3</sup> such as Ge<sub>c</sub>Si<sub>1-c</sub>S<sub>i</sub> (in which the relative concentration c ranges from 0 to 1 over 100 Å or longer in the longitudinal z direction) brings in an additional freedom in the crystal structure and provides a small Brillouin wave number  $k_{-} = 2n/s$  with s being the periodicity length. We demand that the x-ray light in the crystal channel walls becomes a slow wave and satisfies the high-energy acceleration condition

$$(z_{s}) = c_{s}$$
 (2)

where  $\omega$  and  $k_z$  are the light frequency and longitudinal wave number.

The energy loss of moving particles in matter is due to ionization, bremsstrahlung, and nuclear collisions. We can show<sup>6</sup> that a channeled high-energy particle moving fast in the z direction oscillates in the xy plane according to the Hamiltonian

$$H = \frac{1}{2m} (p_x^2 + p_y^2) + V(x,y), \quad (3)$$



## X-ray LWFA in crystal suggested

#### X-ray Laser Wakefield Accelerator in crystal:

LWFA pump-depletion length:

$$L_{acc} \sim a_{\chi} (c/\omega_{\rho}) (\omega_{\chi}/\omega_{\rho})^{2}, \qquad (a_{\chi} = eE_{\chi}/mc\omega_{\chi})$$

LWFA energy gain

 $\varepsilon_{\chi}=2a_{\chi}^{2}mc^{2}(n_{cr}/n_{e}),$ 

Here,  $n_{cr} = 10^{29}$ ,  $n_e = 10^{23}$ ,  $a_{\chi} \sim 30$  (pancake laser pulse with the Schwinger intensity, with focal radius assumed the same as optical laser radius. Could be greater if we further focus by optics, or nonlinearity, or if we not limit the intensity at Schwinger. see below)

The vacuum self-focus power threshold

 $P_{cr} = (45/14) c E_S^2 \lambda^2 \alpha^{-1}, \qquad (E_S: \text{Schwinger field})$ 



#### **Schwinger** fiber acceleration in vacuum:

(no surface, no breakdown)

Vacuum photon dispersion relation with focus

$$\omega = c \vee (k_z^2 + < k_{perp}^2 >),$$



The vacuum dispersion relation with fiber self-modulation  $\omega / (k_z + k_s) = c, \quad (k_s = 2\pi / s)$ 

(Tajima and Cavenago, PRL, 1987)

# Wakefield on a chip toward TeV over cm



## With Fermilab: carbon nanotube acceleration

#### 16th Advanced Accelerator Concept Workshop (AAC2014)



### **TeV/m Nano-Accelerator**

#### **Current Status of CNT-Channeling Acceleration Experiment**



#### Y. M. Shin<sup>1,2</sup>, A. H. Lumpkin<sup>2</sup>, J. C. Thangaraj<sup>2</sup>, R. M. Thurman-Keup<sup>2</sup>, P. Piot<sup>1,2</sup>, and V. Shiltsev<sup>2</sup>

Thanks to X. Zhu, D. Broemmelsiek, D. Crawford, D. Mihalcea, D. Still, K. Carlson, J. Santucci, J. Ruan, and E. Harms

<sup>1</sup>Northern Illinois Center for Accelerator and Detector Development (NICADD), Department of Physics, Northern Illinois University

<sup>2</sup>Fermi National Accelerator Laboratory (FNAL)

## Ultrahigh Energy Cosmic Rays ບໍ່ອີ (UHECR)

Fermi mechanism runs out of steam beyond 10<sup>19</sup> eV due to synchrotron rad Wakefield acceleration comes in rescue prompt, intense, linear acceleration



## Cen A



- Distance: 3.4Mpc
- Radio Galaxy
  - Nearest
  - Brightest radio source
- Elliptical Galaxy
- Black hole at the center w/ relativistic jets



### Astrophysical wakefield acceleration: Superintense Alfven Shock in the Blackhole Accretion Disk toward ZeV Cosmic Rays ( $a_0 \sim 10^6 - 10^{10}$ , large z)

Magnetic field lines

ZEST



Ebisuzaki and Tajima, Astropart. Phys. (2014)



# comic ray acceleration and gamma-ray emission





#### HOW A "CAN" LASER AMPLIFIER WORKS

Producing High Peak Power and High Average Power, Mitigating Heat

(CAN: Coherent Amplifier Network)

The stretched pulse, is coupled to a multiplicity of single mode fiber amplifiers. Each fiber will amplify the input pulse to the *mJ* level. An oscillator produces a short pulse of ~*30fs* duration.



3

After the last amplifier the pulse is focused by spherical or paraboloidal mirror.

6

The resulting pulse is short (*30 fs*), but the energy is enormous (*30 Joules*)



The same operation is repeated in a second and third amplifier stage where each fiber amplifier of the first stage feeds a multiplicity (10-100-10000) of single mode amplifiers. In turn each fiber will amplify its input to the *mJ* level.



The pulse is first fed into a single mode optical fiber amplifier and passes through a pair of diffraction gratings, which stretch the pulses by around 10^5 times. The pulse after stretching is at the *mJ* level. The stretching separates the various components of the stretch pulse, producing a rainbow in time.



## Conclusions

- A new direction of ultrahigh intensity: zeptosecond lasers
- EW 10keV X-rays laser from 1PW optical laser
- X-ray LWFA in crystal: accelerating gradient 1-10TeV/cm, TeV on a chip
- **Crystal nanoengineering**: s.a. nanoholes, arrays, focus optics
- **Zeptosecond nanobeams** of electrons, protons (ions), muons (neutrinos), **coherent** γ-rays
- Start of **zeptoscience**

#### In dedication to the late Professor Hiroshi Takuma:

#### 「道を問ひ 学に魅入りし 木津の地に 御霊残れり 微笑みととも」

"Inquiring the Tao Lured in beauty of physics Your Spirit remains Here in the Kizu 'Shrine' You founded, with smile on us....'

T. Tajima

## ありがとうございました!