

Summer School in Plasma Physics
The Abdus Salam International Centre for Theoretical Physics
8/17/09

Laser Acceleration and its Scope and Impact: 1979-2009

Toshi Tajima
LMU, MPQ
Garching

Acknowledgments for Advice and Collaboration: Iate-J. Dawson, N. Rostoker, F. Krausz, D. Habs, S. Karsch, L. Veisz, F. Gruener, G. Mourou, T. Esirkepov, M. Kando, K. Nakajima, A. Chao, A. Suzuki, F. Takasaki, S. Bulanov, A. Giullietti, F. Mako, X. Yan, J. Meyer-ter-Vehn, W. Leemans, T. Raubenheimer, A. Ogata, A. Caldwell, P. Chen, Y. Kato, Iate-A. Salam, M. Downer, S. Ichimaru, M. Tigner, V. Malka, A. Henig, H.C. Wu, K. Kondo, Y. Sano, M. Abe, S. Kawanishi, M. Hegelich, D. Jung, P. Shukla

Can the society continue to support ever escalating accelerators?



Accelerator = crown of 20th C science

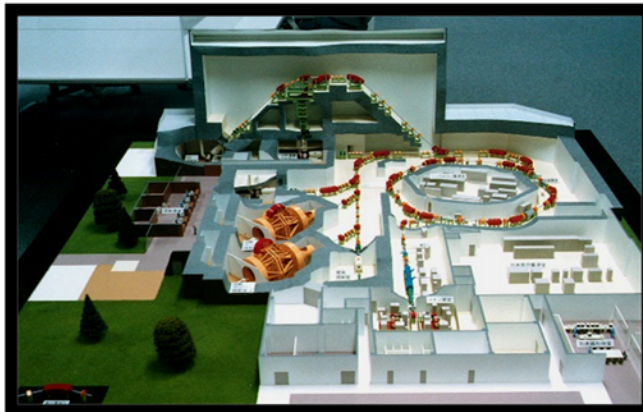
beam dump



LHC at CERN



supermagnets quench



hadron therapy accelerator and gantry



Terminated Texas tunnel. The SSC was abandoned after about 25% of the tunnel for the 87-kilometer-circumference large collider ring had been bored.

SSC tunnel



Dream Beams Symposium

MPQ Garching
Feb. 26 – 28, 2007

(given by F. Krausz and J. Meyer-ter-Vehn)

Demise of SSC (Super collider)



Terminated Texas tunnel The SSC was abandoned after about 25% of the tunnel for the 87-kilometer-circumference large collider ring had been bored.

By largest machine to probe smallest of structure of matter

| | |
|--------|--------------------|
| size | 10 ² km |
| energy | 20TeV |
| cost | \$10B |

US:

Texas site decided (1989)

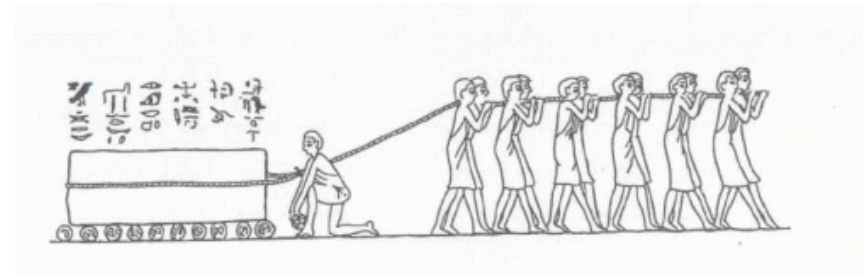
US Government decided to terminate its work: 1993

Tajima: 'Tamura Symposium' on the Future of Accelerator Physics @ UT Austin
(1995)

What is collective force ?



How can a Pyramid have been built?



Individual particle dynamics → Coherent and collective movement

Collective acceleration (Veksler,1956; Tajima & Dawson,1979)

Collective radiation (N^2 radiation)

Collective ionization (N^2 ionization)

Collective deceleration (Tajima & Chao,2008)



With giants of collective phenomena

Physics of **individual** particles;

Physics of **collection** of particles---collective phenomena



Professor Ryogo Kubo



Professor Iliya Prigogine

(Austin, ~1984)

Advent of collective acceleration(1956)



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CERN Symposium

ON HIGH ENERGY ACCELERATORS
AND PION PHYSICS

Geneva, 11th - 23rd June 1956

Proceedings

COHERENT PRINCIPLE OF ACCELERATION OF CHARGED PARTICLES

V. I. VEKSLER

Electrophysical Laboratory, Academy of Sciences, Moscow

This paper will include a very brief description of a new principle regarding the acceleration of charged particles.

In all existing accelerators of charged particles, the constant and varying electric field accelerating them is created by a powerful external source, and hence the strength of the field is independent, in the first approximation, of the number of particles which are being accelerated.

In resonance accelerators, the electromagnetic field has to be synchronized with the movement of the particles (this is of particular importance in linear accelerators). Finally, none of the existing methods permits the acceleration of neutral bunches of particles.

A new principle of particle acceleration is set forth below. Its distinctive feature lies in the fact that the particle-accelerating electric field is produced by the interaction of a geometrically small group of accelerated particles with another group of charges, plasma or an electromagnetic wave. This method has a number of important features. It appears, in the first place, that the magnitude of the accelerating field produced by this interaction and acting on each particle depends on the number

Theoretical studies of various aspects of the coherent acceleration method have been made by M. S. Rabinovich, A. A. Kolomenski, B. M. Bolotovskii, L. V. Kovrizhnikh and I. V. Iankov, as well as by A. I. Akhiezer, Ia. Fainberg and their collaborators. The calculations made by these theoretical workers shed light on a number of complicated problems connected with the development of the different variants of this new acceleration principle, and it therefore seems appropriate to describe the new method despite the fact that a great many problems involved still await solution.

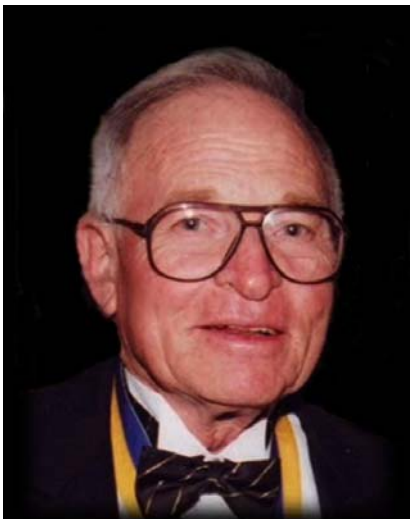
1. *Acceleration of charged bunches by means of the medium*

It was pointed out in a paper by Tamm that the loss of energy by particles due to Čerenkov radiation could be reversed, i.e. the medium travelling at a great velocity past charged particles should be able to convey energy to the latter. Up to now, however, no attention has been paid to the possibility of developing an acceleration process of this kind by using a high density electron beam (plasma) as the moving medium. Of course, if a single charge e is

Prehistoric activities(1973-75,...84)



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N. Rostoker

Collective ion acceleration by a reflexing electron beam: Model and scaling

F. Mako
Naval Research Laboratory, Washington, D. C. 20375

T. Tajima
Institute for Fusion Studies, University of Texas, Austin, Texas 78712

(Received 21 June 1983; accepted 2 April 1984)

Analytical and numerical calculations are presented for a reflexing electron beam type of collective ion accelerator. These results are then compared to those obtained through experiment. By constraining one free parameter to experimental conditions, the self-similar solution of the ion energy distribution agrees closely with the experimental distribution. Hence the reflexing beam model appears to be a valid model for explaining the experimental data. Simulation shows in addition to the agreement with the experimental ion distribution that synchronization between accelerated ions and electric field is phase unstable. This instability seems to further restrict the maximum ion energy to several times the electron energy.

I. INTRODUCTION

Experiments on collectively accelerating ions utilizing a reflexing intense relativistic electron beam in a plasma have been carried out.^{1,2} These experiments began to reveal sever-

chronous fashion. Thus, energetic ions would be expected. The ion energy would, of course, be bounded above by the ion to electron mass ratio times the initial electron energy; that is, the energy is bounded when the ions reach the initial electron energy.

Collective acceleration suggested:

Veksler (1956)

(ion energy)~ (M/m)(electron energy)

Many experimental attempts (~'70s):

led to no such amplification

(ion energy)~ (several)x(electron)

Mako-Tajima analysis (1978;1984)

sudden acceleration, ions untrapped,
electrons return

→ #1 **gradual acceleration necessary**

→ #2 **electron acceleration** possible
with **trapping** (with Tajima-
Dawson field), **more tolerant** for
sudden process

Path once trodden



Collective acceleration of ions by electron beam

R.Mako / T. Tajima

Ions left out, while electrons shoot backward

- laser electron acceleration (1979)
- laser ion acceleration of limited ion mass (2009)

The electric field is

$$\epsilon = \frac{\phi_0}{v_0 t} \frac{5}{36} \left(\frac{6}{\sqrt{3}} - \frac{x}{v_0 t} \right),$$

where the conservation of energy was used as a boundary condition, i.e.,

$$U^2/2 + \psi = 0 \quad \text{at} \quad \xi = 0.$$

The maximum ion energy can now be obtained by setting $n_i = 0$, i.e.,

$$E_{\text{max}} = 6q\phi_0 \quad \text{at} \quad \xi = 6/\sqrt{3}.$$

In the experiment the diode voltage was 0.8 MV and the ions were doubly ionized helium,⁶ thus the maximum ion energy predicted by theory is

$$E_{\text{max}} = 9.6 \text{ MeV}.$$

The experimental result⁶ for the maximum helium ion energy was 9.6 MeV and therefore is in good agreement with the theory.

The ion number as a function of energy is calculated to be

$$N_i(E_i) = \frac{n_0 A}{\beta} \left[\left(\frac{6}{5} \right)^{1/2} - \left(\frac{E_i}{5q\phi_0} \right)^{1/2} \right]^6, \quad (15)$$

where

$$n_0 A = \frac{16 J_0 A}{5 e} \left(\frac{2m}{e\phi_0} \right)^{1/2},$$

$$\beta = (3)^{1/2} (1/v_0 t),$$

$$A = \pi r_0^2, \quad r_0 = \text{electron beam radius},$$

and

$$v_0 = (q\phi_0/M)^{1/2}.$$

Equation (15) is our main result. The natural logarithm of Eq. (15) is plotted in Fig. 2 along with the experimental

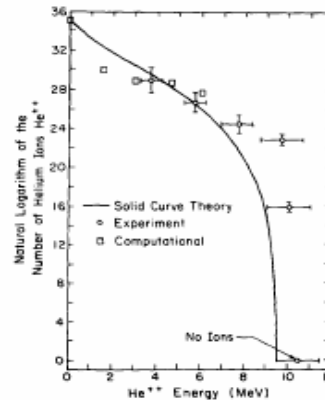


FIG. 2. Comparison between theory, experiment, and simulation, of the natural logarithm of the ion number versus energy.

data. The following experimental values were used: $J_0 = 40$ kA, $\phi_0 = 0.8$ MV, $q = 2e$ (doubly ionized helium), $t = 100$ ns and $r_0 = 2.5$ cm. The agreement between Eq. (15) and the experiment⁶ is reasonable. The relation in Ref. 3 does not provide such a good fit: it has too weak a slope.

III. SCALING AND ACCESSIBILITY OF THE MODEL

In the preceding section, the analysis assumed that a self-similar state could be reached. To address the question of whether a self-similar state can be attained, a detailed analysis of the initial value problem is required. This detailed analysis should include a self-consistent treatment of the dy-

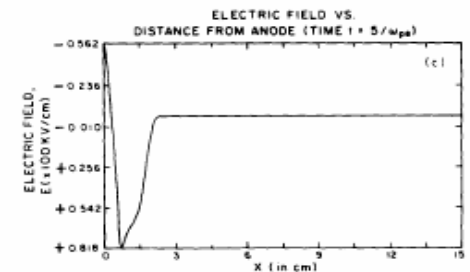
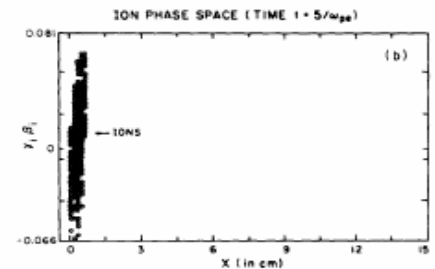
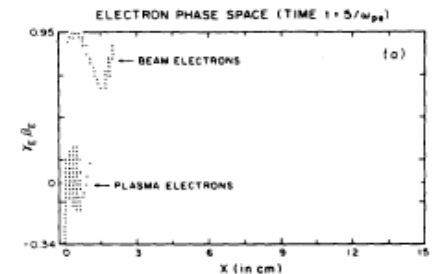


FIG. 3. Simulation phase space at early time $t = 5/\omega_{pe}$. (a) Electron phase space (beam and plasma), (b) ion phase space, (c) electron field versus position.

Laser Acceleration of Electrons

← Lesson #2 of collective ion acceleration



Gradient limit : breakdown threshold for microwave ($< 100\text{MeV/m}$)

E. Lawrence: cyclotron (c. 1932)

SSC: 10^2 km circumference (\dagger 1993); Linear Collider: $> 10\text{km}$ ($\sim 2020?$)

Plasma : already ‘broken’ matter. No breakdown threshold.

‘collective ion acceleration’ (Veksler, 1956): ion trapping difficult

Introduction of laser acceleration (Tajima and Dawson, 1979)

Linear EM field: cannot accelerate: *Woodward-Lawson Theorem*

Strong nonlinear fields

longitudinal acceleration (rectification of laser fields; $\mathbf{v} \times \mathbf{B}/c \sim O(E)$)

laser plays master, plasma slaves----- provides hard structure

electron trapping possible (revisit of ion acceleration now)

→ High Field Science

Ultrafast pulses

fs regime: ions immobile; enhanced with collective electron resonance

absence of ‘notorious’ hydrodynamical plasma instabilities; controllability;

relatively small laser energy (e.g. ELI)

Large gradient ($> 10\text{GeV/m}$, leap by > 3 orders of magnitude)

Low emittance ($< \text{mm mrad}$ regime)

Wakefield: a Collective Phenomenon



All particles in the medium participate = collective phenomenon



Kelvin wake

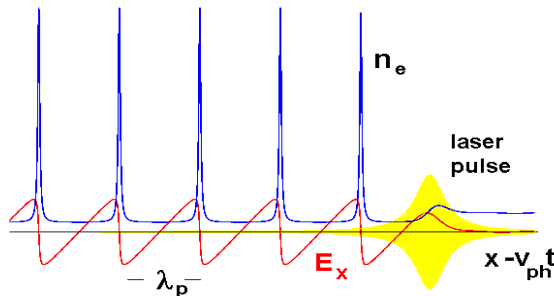
$$\omega = \sqrt{kg}$$

$$x = X_1 \cos \theta \left(1 - \frac{1}{2} \cos^2 \theta \right)$$

$$y = X_1 \cos^2 \theta \sin \theta$$

$$-\pi/2 < \theta < \pi/2$$

No wave breaks and wake **peaks** at $v \approx c$



$$\lambda_p = 2\pi / k_p \quad k_p v_{ph} = \omega_{pe}$$

$$\omega_{pe} = \left(4\pi n e^2 / m_e \right)^{1/2} \quad \text{(The density cusps. Cusp singularity)}$$

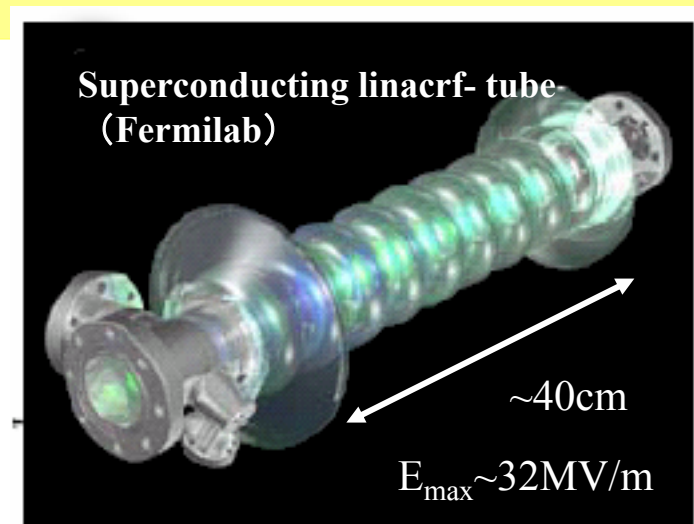
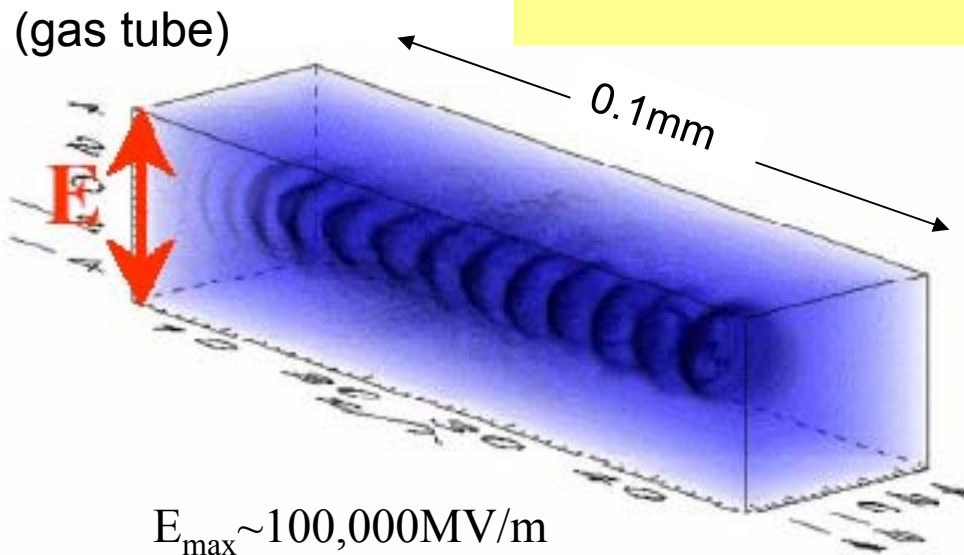
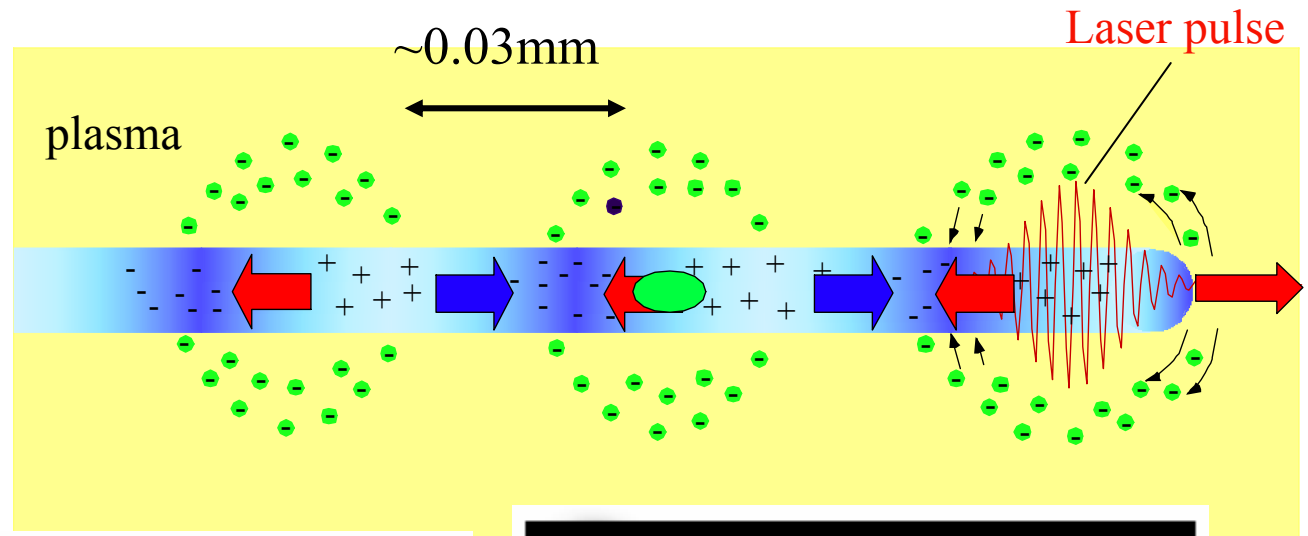
Wave **breaks** at $v < c$



Thousand-fold Compactification



Laser wakefield: thousand folds gradient (and emittance reduction?)





The late Prof. Abdus Salam



At ICTP Summer School (1981), Prof. Salam summoned me and discussed about laser wakefield acceleration.

Salam: ‘Scientists like me began feeling that we had less means to test our theory. However, with your laser acceleration, I am encouraged’. (1981)

He organized the Oxford Workshop on laser wakefield accelerator in 1982.

Effort: many scientists over many years to realize his vision / dream
High field science: spawned

Laser technology invented (1985)



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Chirped pulse amplification (CPA) invented:

to overcome the gain medium nonlinearities in spatially expanded amplification to temporal expansion:

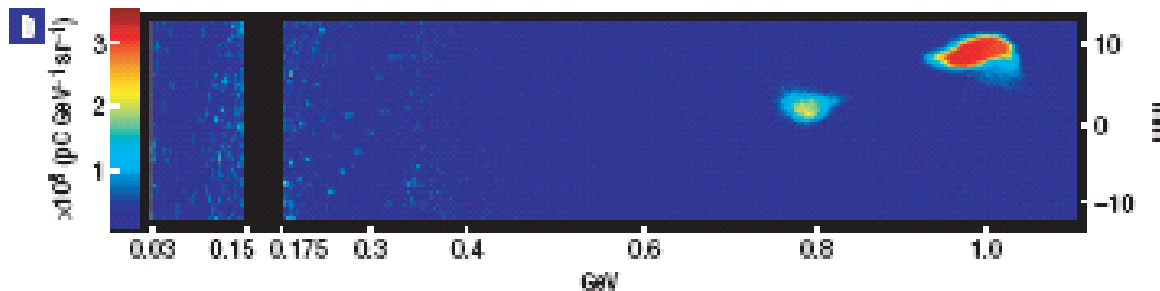
smaller, shorter pulse, more intense,
higher replate,
all simultaneous.

(Professor Gerard Mourou)

→ many table-top TW and PW lasers world-wide
first Chair, ICUIL (International Committee for
Ultra Intense Lasers)
toward EW laser (*Extreme Light Infrastructure*)

→ First LWFA experiments
(Nakajima et al 1994; Modena et al 1995)
→ derives High Field Science

GeV electrons from a centimeter accelerator (a slide from S. Karsch)



310- μ m-diameter
channel capillary

$P = 40$ TW

density $4.3 \times 10^{18} \text{ cm}^{-3}$.

Leemans et al., Nature Physics, september 2006

Laser Electron Accelerator

T. Tajima and J. M. Dawson

Department of Physics, University of California, Los Angeles, California 90024

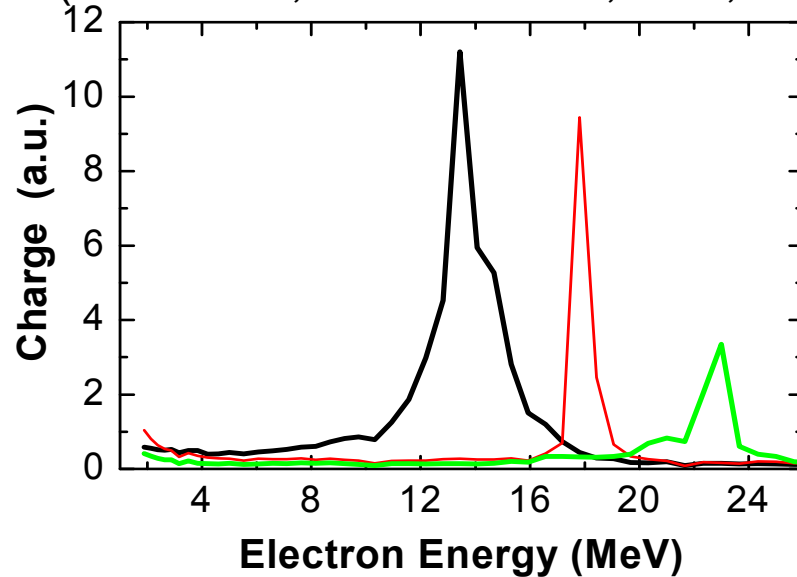
(Received 9 March 1979)

An intense electromagnetic pulse can create a weak of plasma oscillations through the action of the nonlinear ponderomotive force. Electrons trapped in the wake can be accelerated to high energy. Existing glass lasers of power density 10^{18} W/cm^2 shone on plasmas of densities 10^{18} cm^{-3} can yield gigaelectronvolts of electron energy per centimeter of acceleration distance. This acceleration mechanism is demonstrated through computer simulation. Applications to accelerators and pulsers are examined.

MPQ Laser Acceleration Effort (1)

Monoenergy electron spectra: from **few-cycle laser** (LWS-10)

(K. Schmid, L. Veisz et al., PRL, 2009)

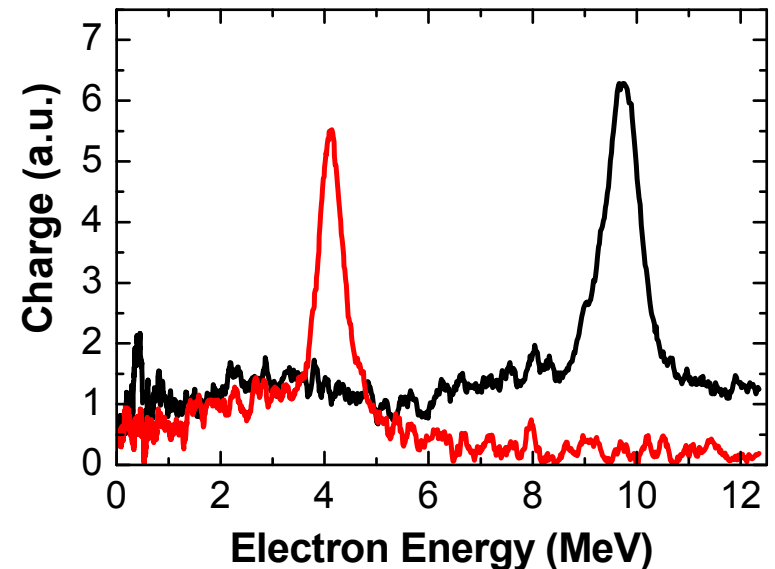


Large electron spectrometer 2 – 400 MeV

- No thermal background !
- Energies: 13.4 MeV, 17.8 MeV, 23 MeV
- FWHM energy spread: 11%, 4.3%, 5.7 %
- ~ 10 pC charge

Small electron spectrometer:

- Electron energies below 500keV
- No thermal background !
- 4.1 MeV (14%); 9.7 MeV (9.5%)



MPQ Laser Acceleration Effort (2)



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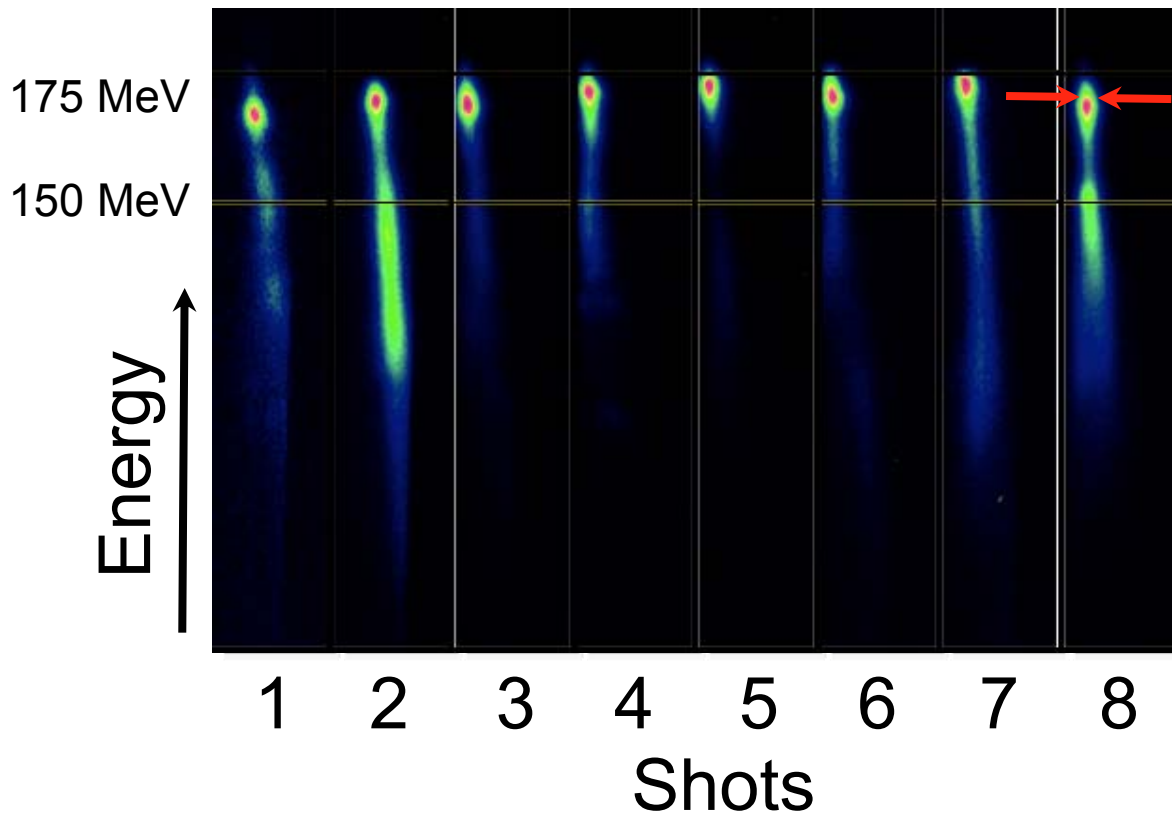
Reproducible acceleration conditions

$$E \approx 169.7 \pm 2.0 \text{ MeV}$$

1.1% peak energy
fluctuation !

$$\Delta E/E \approx 1.76 \pm 0.26\% \text{ RMS}$$

→ Essential property for
future table-top FEL operation



Source size image: provides
emittance measurement,
given the resolution can be
improved

(J. Osterhoff, ... S. Karsch, PRL 2008 et al.)

MPQ Laser Acceleration Effort (3)



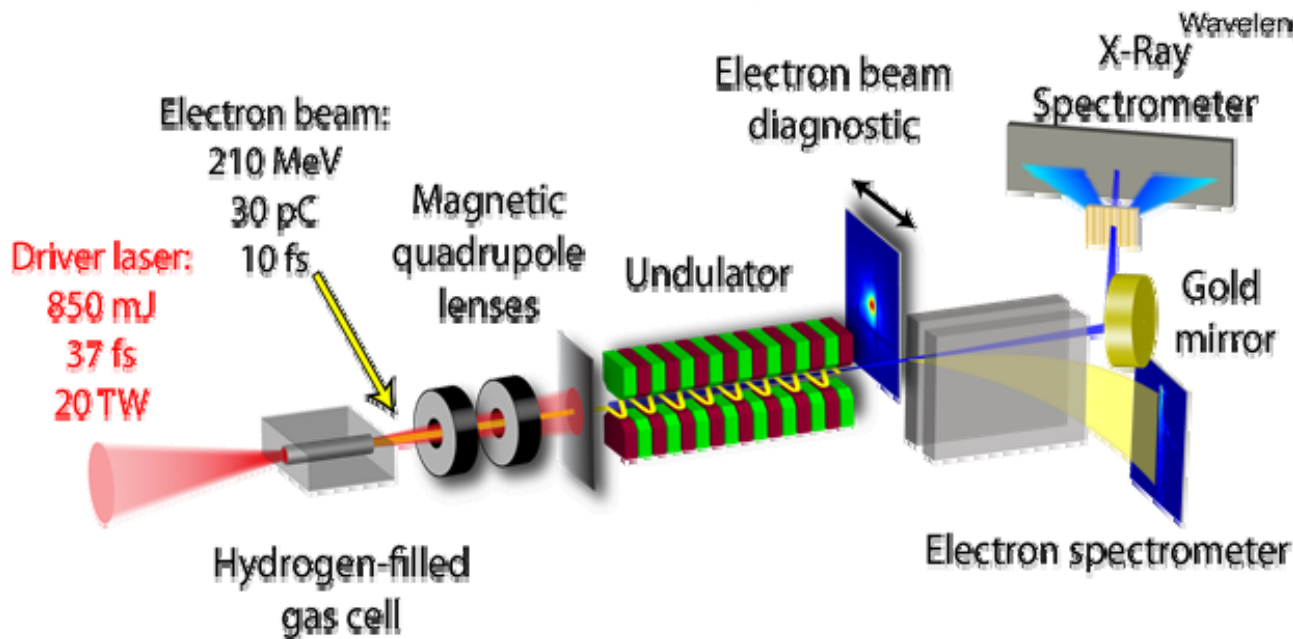
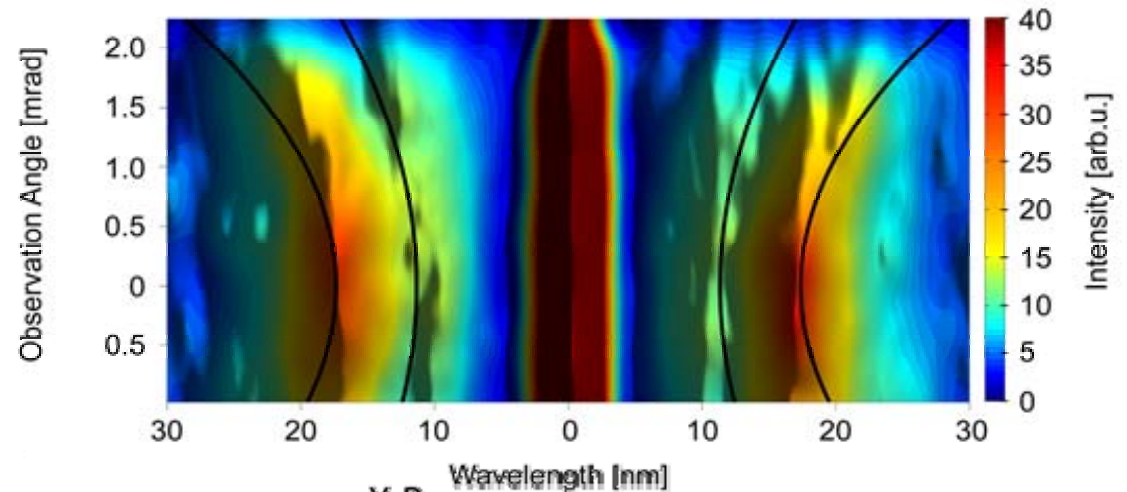
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Laser-driven Soft-X-Ray Undulator Radiation

(F. Gruener, S. Karsch, et al., accepted by Nature Phys.)

Characteristic undulator radiation spectrum



Intra-Operatory Radiation Therapy (IORT)



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LWFA electron sources: *technology transferred to co.*

NOVAC7

(HITESYS SpA)
RF-based

VS.

CEA-Saclay
experim. source
Laser-based

El. Energy < 10 MeV
(3, 5, 7, 9 MeV)

El. Energy > 10 MeV
(10 - 45 MeV)

Peak curr. 1.5 mA

Peak curr. > 1.6 KA ←

Bunch dur. 4 μs

Bunch dur. < 1 ps ←

Bunch char. 6 nC

Bunch char. 1.6 nC

Rep. rate 5 Hz

Rep. rate 10 Hz

Mean curr. 30 nA

Mean curr. 16 nA

Releas. energy (1 min)
@9 MeV (≈dose)
18 J

Releas. energy (1 min)
@20 MeV (≈dose)
21 J



(A. Giulietti et al., Phys. Rev. Lett., 2008)





Collective deceleration

Beam dump: harder to stop and more hazardous radioactivation



Gas (plasma) **collective force** to shortstop the HE beams

- the shorter the bunch is, the easier to stop
(ideally suited for laser wakefield accelerated beams)
- little radioactivation (good for environment)
example of **Toilet Science** that tends impact of own produce
(as opposed to 'Kitchen Science' of 20th C)
- possible energy recovery

Beam Stopping and its Energy Recovery Using Plasma

February 25, 2008
Toshiki Tajima and Alexander W. Chao

1 Motivations

1.1 Beam Stopping

In the effort to make a high energy accelerator system as compact as possible, it is necessary not only to make the accelerator compact, but also to make the beam stopping system compact. With this motivation, we introduce the concept of passive plasma decelerator at the end of the use of the high energy beam by immersing the beams to be decelerated into an appropriately designed plasma.

Stopping power due to collective force



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Bethe-Bloch stopping power in matter

Plasma stopping power due to individual force

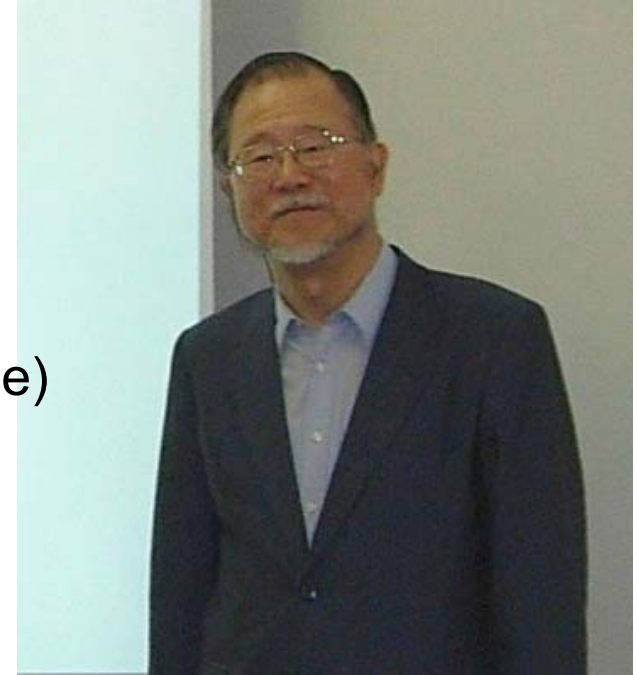
$$-(dE / dx)_{ind} = (F / \beta^2) \ln(m_e v^2 / e^2 k_D)$$

That due to collective force (perturbative regime)

$$-(dE / dx)_{coll} = (F / \beta^2) \ln(k_D v / \omega_{pe})$$

$$F = 4\pi e^4 n_{e,m} / m_e c^2 = e^2 k_{pe,m}^2$$

(Ichimaru, 1973)



Professor Setsuo Ichimaru

Plasma stopping power due to short-bunch wakefield (wavebreak regime)

$$-(dE / dx)_C = m_e c \omega_{pe} (n_b / n_e)$$

(Wu et al, 2009)

Greater by **several orders** over Bethe-Bloch in solid

Keys issues of future colliders



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(T. Raubenheimer, SLAC, 2008)

Beam Acceleration

- * Largest cost driver for a linear collider is the acceleration
 - ILC geometric gradient is ~ 20 MV/m \rightarrow 50km for 1 TeV
- * Size of facility is costly \rightarrow higher acceleration gradients
 - High gradient acceleration requires high peak power and structures that can sustain high fields
 - Beams and lasers can be generated with high peak power
 - Dielectrics and plasmas can withstand high fields
- * Many paths towards high gradient acceleration
 - High gradient microwave acceleration } ~ 100 MV/m
 - Acceleration with laser driven structures } ~ 1 GV/m
 - Acceleration with beam driven structures } ~ 1 GV/m
 - Acceleration with laser driven plasmas } ~ 10 GV/m
 - Acceleration with beam driven plasmas } ~ 10 GV/m

Challenge posed by DG Suzuki



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Frontier science driven by advanced accelerator

Table-top X-ray FEL

PeV=10¹⁵ eV

$\alpha = \frac{\hbar^2}{e c}$ PeV Accelerator

“New paradigm”

1000 times higher energy

3rd-generation Synchrotron Light Source

Leptogenesis

SUSY breaking

Extra dimension

Dark matter

Supersymmetry

TeV=10¹² eV

“Standard model”

Higgs

Quarks

Leptons

100 GV/m

Plasma Acceleration Technology

10/39

1 fs = 10⁻¹⁵ s

1000 times shorter time resolution

Rhodopsin ~200 fs

Photosynthetic reaction in leaves ~100 fs

Photo-switching of metal-to-insulator

1 ps = 10⁻¹² s

FEMTO-SEC BEAM TECHNOLOGY

13/39

compact, ultrastrong a

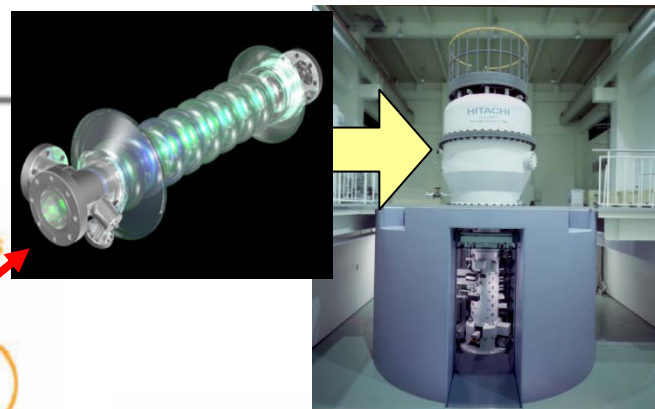
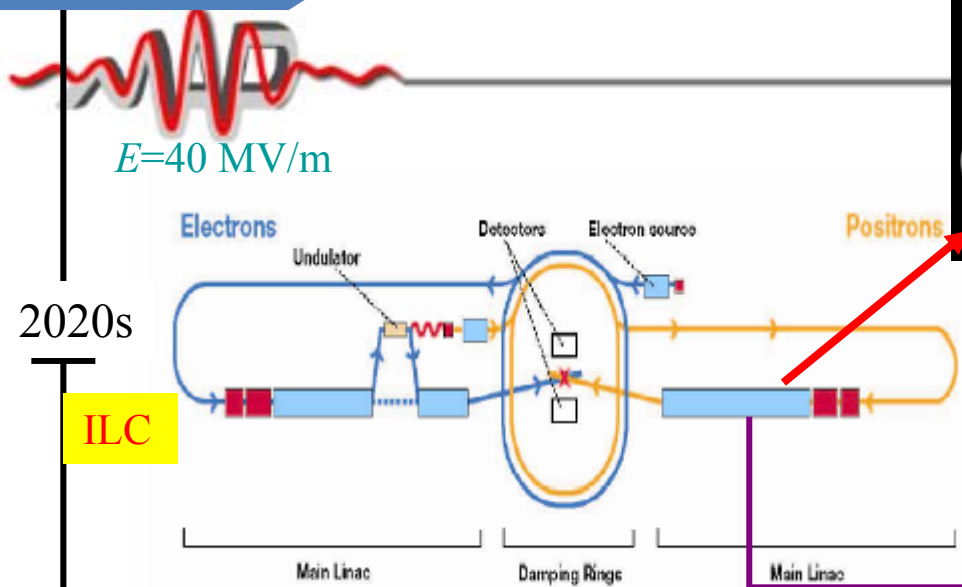
atto-, zeptosecond

Can we meet the challenge?

A. Suzuki @KEK(2008)

Accelerator

Evolution of Accelerators and their Possibilities (Suzuki, 2008)



Ultra-High Voltage STEM with Superconducting RF cavity

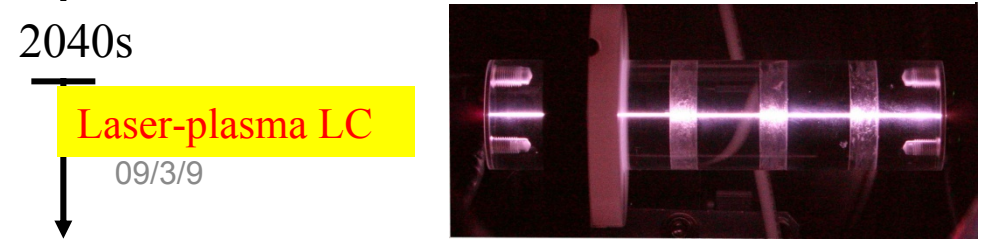
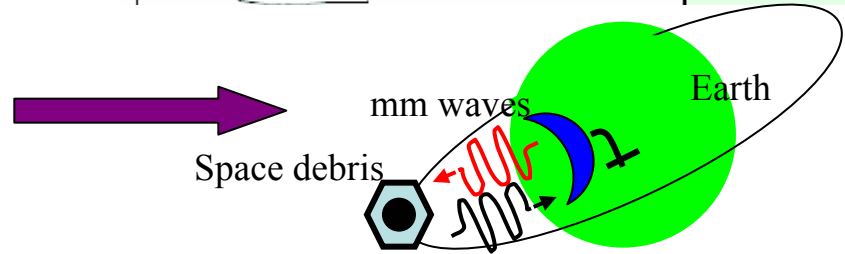
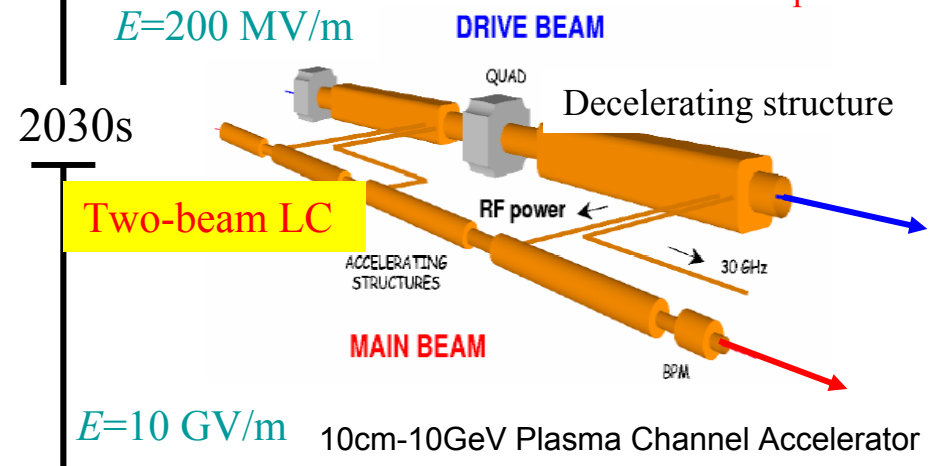
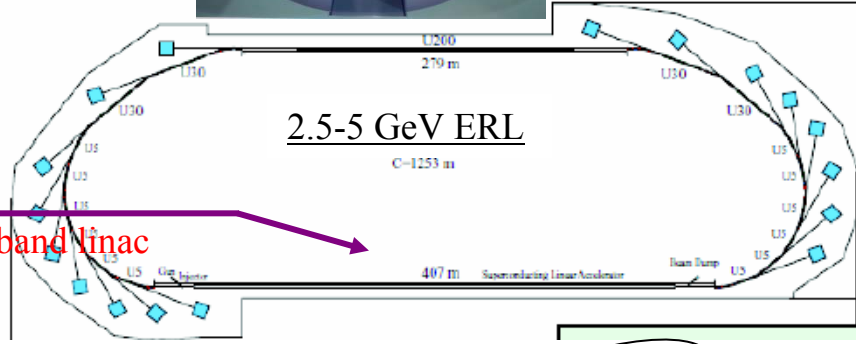


Table-top high energy accelerator

Meeting Suzuki's Challenge:



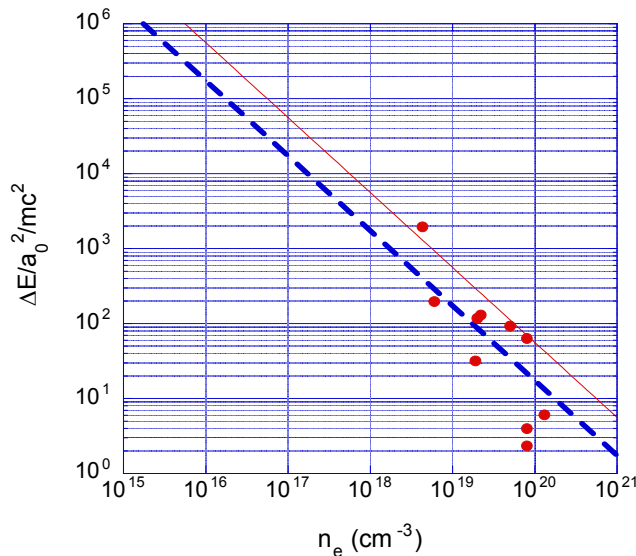
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Laser acceleration toward ultrahigh energies

$$\Delta E \approx 2m_0c^2 a_0^2 \gamma_{ph}^2 = 2m_0c^2 a_0^2 \left(\frac{n_{cr}}{n_e} \right), \quad (\text{when 1D theory applies})$$

$$L_d = \frac{2}{\pi} \lambda_p a_0^2 \left(\frac{n_{cr}}{n_e} \right), \quad L_p = \frac{1}{3\pi} \lambda_p a_0 \left(\frac{n_{cr}}{n_e} \right),$$

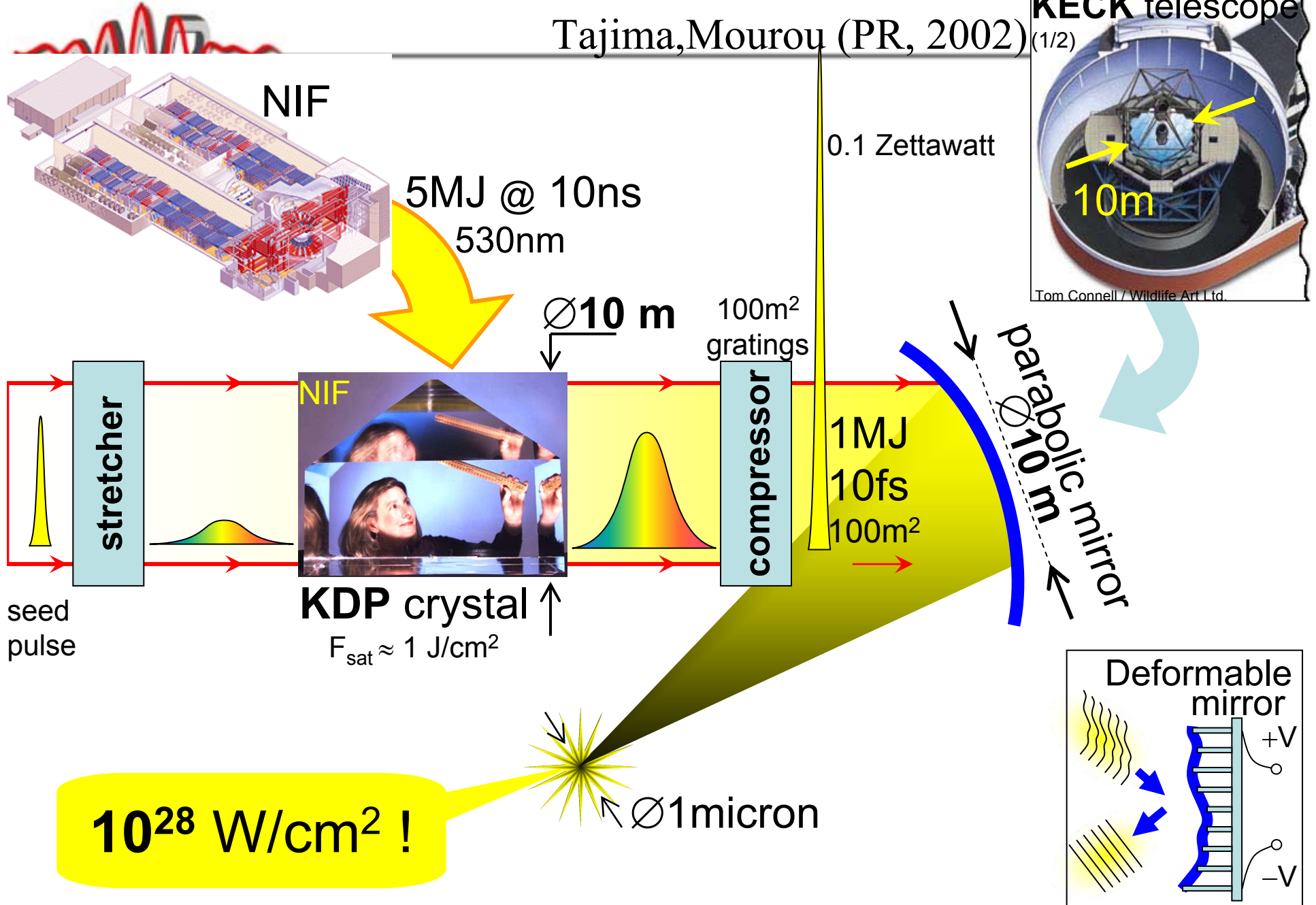


| | | case I | case II | case III |
|---------------------|------------------|----------------------|----------------------|----------------------|
| a_0 | | 10 | 3.2 | 1 |
| energy gain | GeV | 1000 | 1000 | 1000 |
| plasma density | cm ⁻³ | 5.7x10 ¹⁶ | 5.7x10 ¹⁵ | 5.7x10 ¹⁴ |
| acceleration length | m | 2.9 | 29 | 290 |
| spot radius | μm | 32 | 100 | 320 |
| peak power | PW | 2.2 | 2.2 | 2.2 |
| pulse duration | ps | 0.23 | 0.74 | 2.3 |
| laser pulse energy | kJ | 0.5 | 1.6 | 5 |

Even 1PeV electrons (and gammas) are possible, albeit with lesser amount
 → exploration of new physics such as the reach of relativity and beyond?
 (laser energy of 10MJ, plasma density of 10¹⁶/cc)

Zettawatt Laser

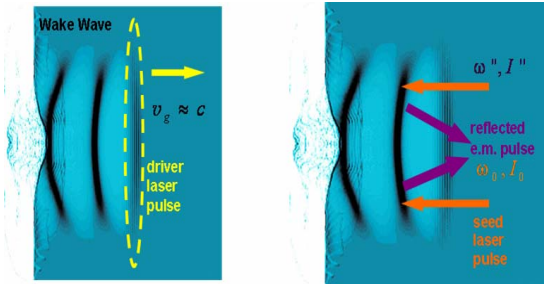
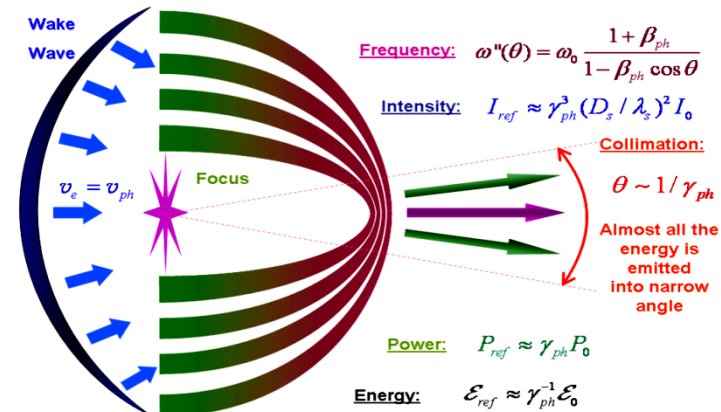
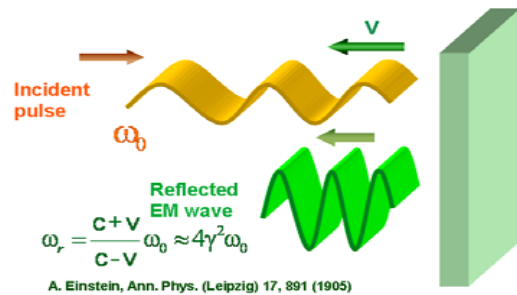
Tajima, Mourou (PR, 2002) ^(1/2)



Relativistic Engineering: an example

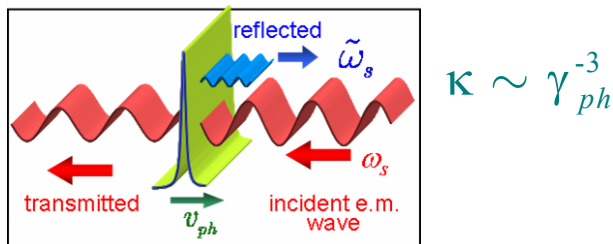
EM Pulse Intensification and Shortening by the Flying Mirror

(Bulanov, Esirkepov, Tajima, 2003)

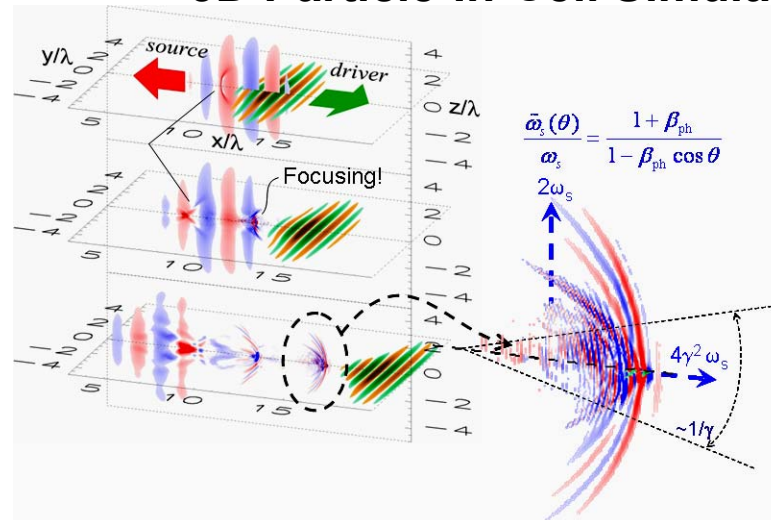


$$\omega'' = \frac{c + v_{ph}}{c - v_{ph}} \omega \approx 4\gamma_{ph}^2 \omega_0$$

$$\frac{I''_{max}}{I_0} \approx \kappa \gamma_{ph}^6 \left(\frac{D}{\lambda} \right)^2$$



3D Particle-In-Cell Simulation



A lot ideas for new attosecond pulses

Relativity helps acceleration

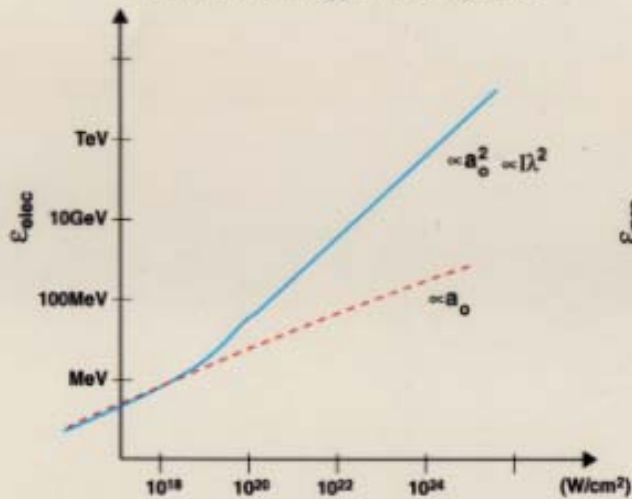
Extreme Field Science



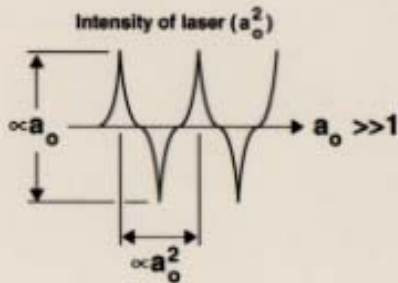
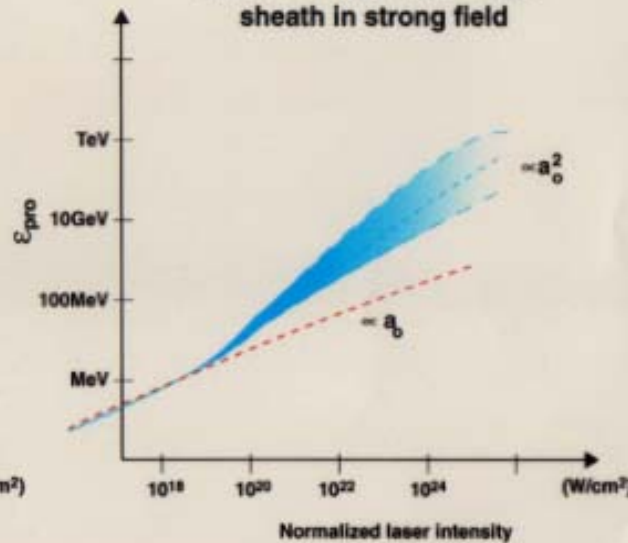
The National Ignition Facility

Ultra-relativistic Regime:
charged particles move with photons

Electron Energy in strong field



Proton Energy from Debye sheath in strong field



$$a_0 \sim L.5 \left(\frac{\lambda}{1 \mu\text{m}} \right) \left(\frac{I}{10^{20} \text{ W/cm}^2} \right)$$

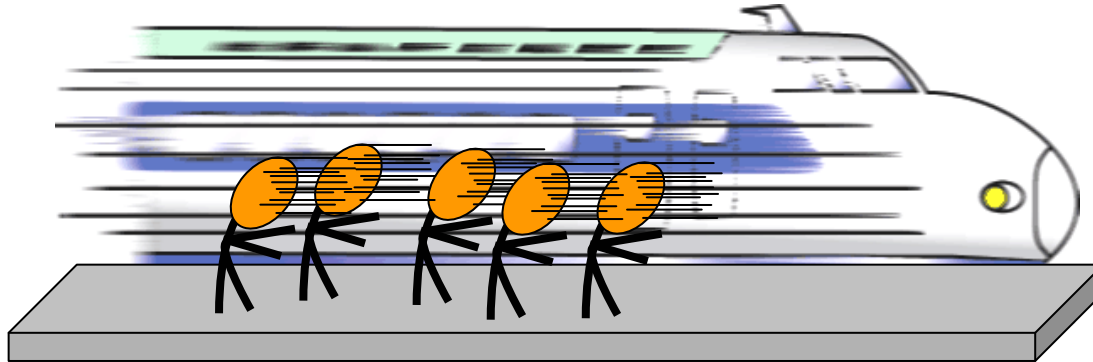
Strong fields:
rectifies laser
to longitudinal
fields

In relativistic regime,
photon x electrons
and even protons
couple **stronger**.

(Tajima, 1999
@LLNL;
Esirkepov et al.,
PRL,2004)

Adiabatic (gradual) acceleration

↓ Accelerating structure

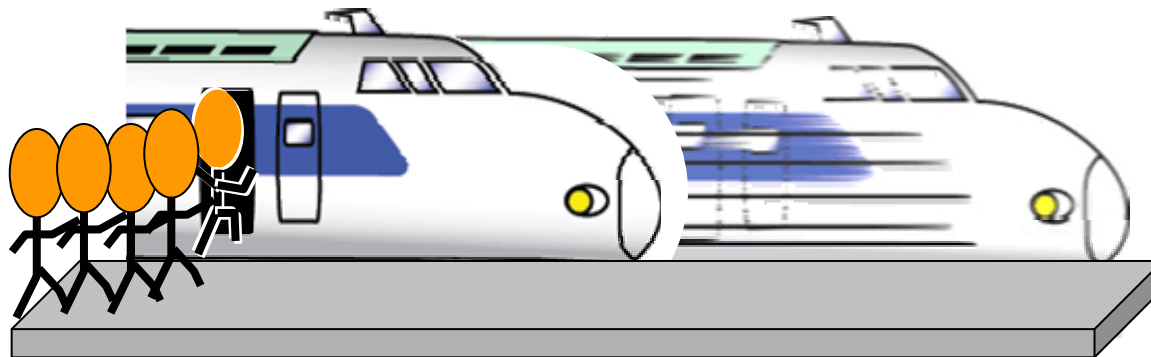


**Inefficient if
suddenly
accelerated**

protons ↑



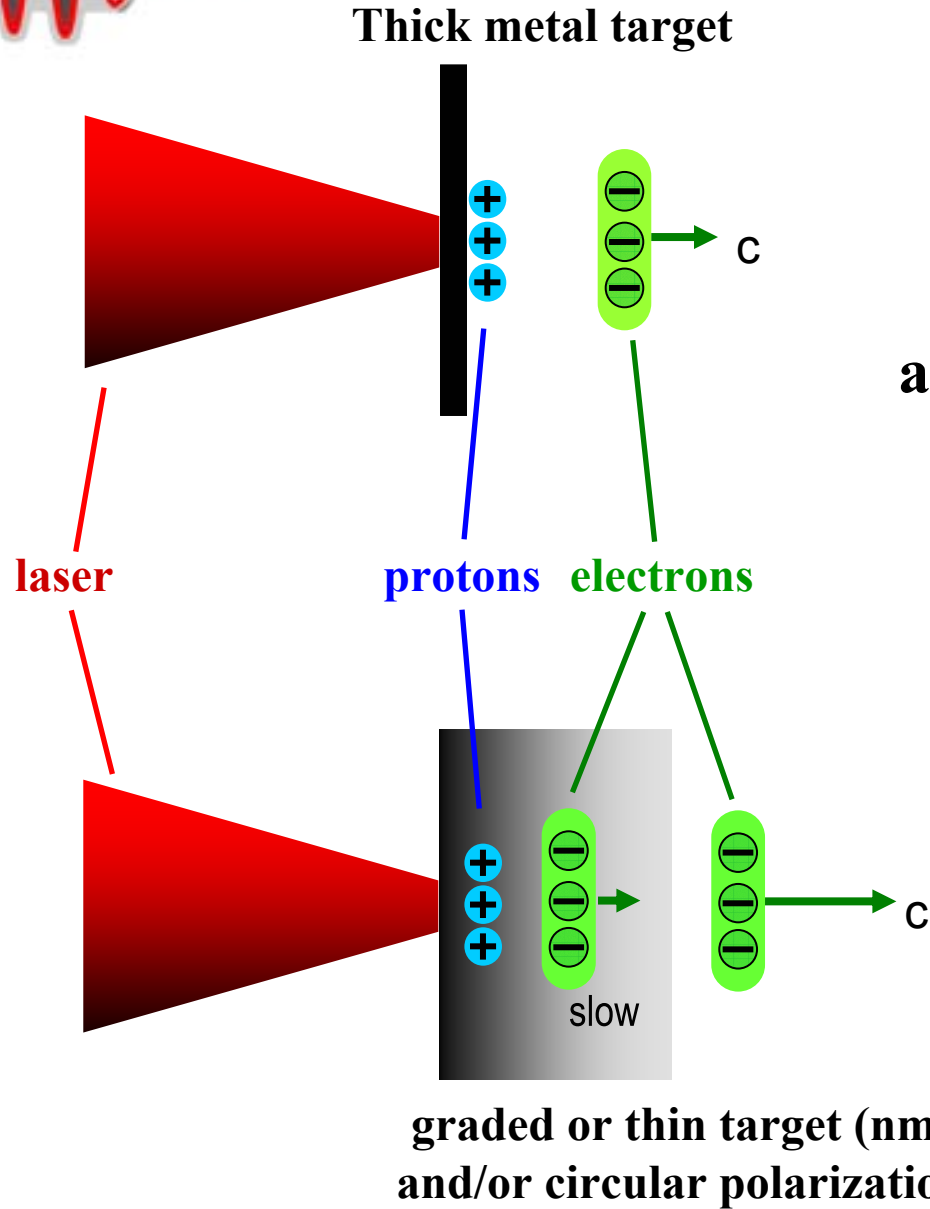
↓ Accelerating structure



**Efficient
when
gradually
accelerated**

#1 → Relevant to ion acceleration

Adiabatic acceleration(2)



Most experimental configurations of proton acceleration(2000-2009)

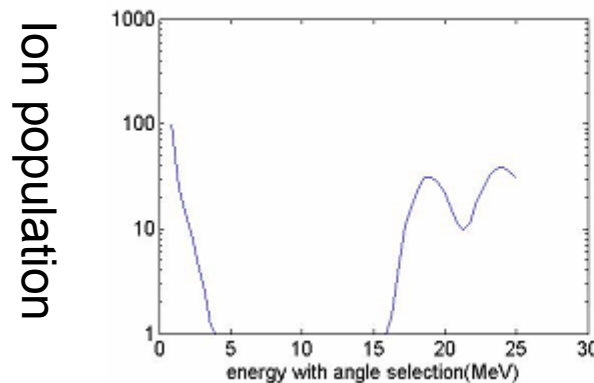
**Innovation (“Adiabatic Acceleration”)
(2009-)**

**graded or thin target (nm)
and/or circular polarization**

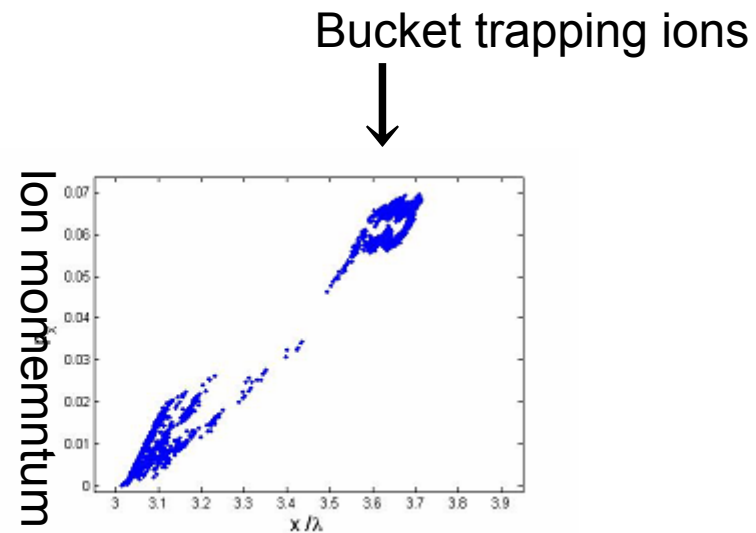
Circularly polarized laser driven



CP laser drives ions out of ultrathin (nm) foil **adiabatically**
Monoenergy peak emerges



laser →



→

Ponderomotive force drives electrons,
Electrostatic force nearly cancels
Slowly accelerating bucket formed

(X. Yan et al: 2009)

Conclusions

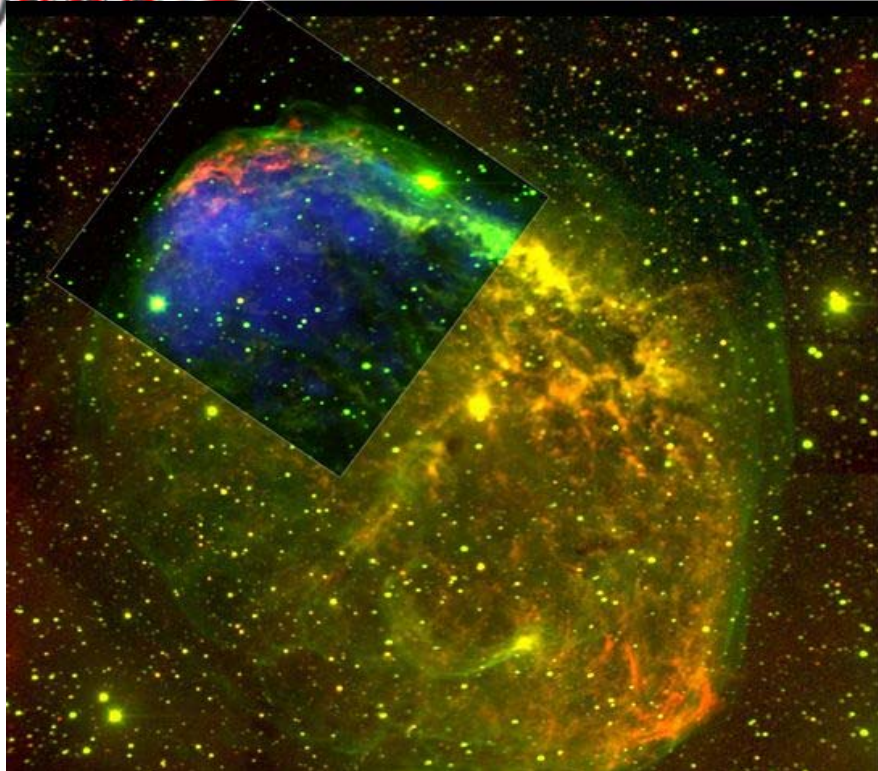


- **Collective acceleration: hard birth / long way and near maturation**(electron→ion; **laser**→electron; **laser**→photon; electron→electron; ion→electron); unexpected ‘homecoming’ (**laser** →ion), too
- Leap by **many orders** (≥ 3) in many respects; equally more demanding by many orders : **N²** vs. N.
- **Laser** has come around to match the condition set 30 years ago; Still some ways to go to realize the dream
- GeV electrons; 10 GeV soon; 100GeV contemplated; **laser** collider contemplated; PeV ?
- **Societal obligations and applications**: already beginning, soon to flourish (e.g., **cancer therapy**, radiolysis, **bunch decelerator**, nuclear detection, compact FEL source, compact radiation sources, ultrafast diagnosis,...)

Cosmic acceleration in the wake of intense radiation and particle flows



LMU
www.attoworld.de



UHECR:
beyond **Fermi** acceleration,
wakefiled acceleration?

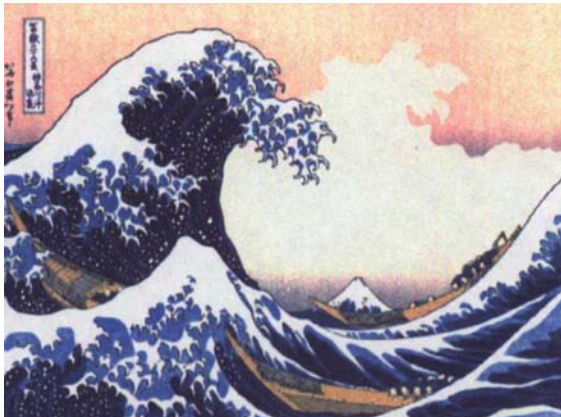
NGC 6888, also known as the Crescent Nebula, is a cosmic bubble of interstellar gas about 25 light-years across. Created by winds from the bright, massive star seen near the center of this composite image, the shocked filaments of gas glowing at optical wavelengths are represented in green and yellowish hues. X-ray image data from a portion of the nebula viewed by the Chandra Observatory is overlaid in blue. Such isolated stellar wind bubbles are not usually seen to produce energetic x-rays, which require heating gas to a million degrees celsius. Still, NGC 6888 seems to have accomplished this as slow moving winds from the central star's initial transition to a red supergiant were overtaken and rammed by faster winds driven by the **intense radiation** from the star's exposed inner layers.



Grazie!



Late-Professor John Dawson



(based on the slide given
by F.Krausz and J. Meyer-te-Vehn
from Dream Beams Symposium, 2007)