

HOUSEKEEPING: • SEE SCHEDULE ONLINE, LET ME KNOW IF YOU SNAG SLOTS
 • THIS SEMESTER: JOINT MTGS ONCE A MONTH w/ hep-ex
 → TO COMPLEMENT FRIDAY JOINT MEETINGS

TODAY: 30 MIN THEORY SUMMARY by HLP
 30 MIN EXP. SUMMARY by WALTER VIA SKYPE

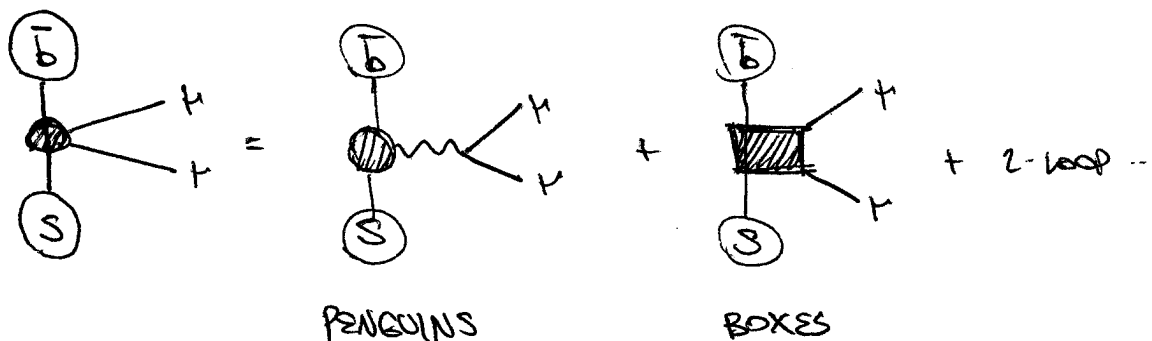
→ THIS IS A DISCUSSION, PLEASE ASK (ANSWER) QUESTIONS

THE BASICS: def (PDG): $\begin{matrix} B_d^0 \equiv \bar{b}d \\ B_s^0 \equiv \bar{b}s \end{matrix}$ ("B meson")
 ("B_s meson")

- PSEUDOSCALAR MESONS
- $m_b \gg m_d, m_s \dots \rightarrow$ (HRET) ← but irrelevant for this decay

FLAVORFUL: CARRIES -1 bottom charge +1 strange charge } but B^0 is charge neutral

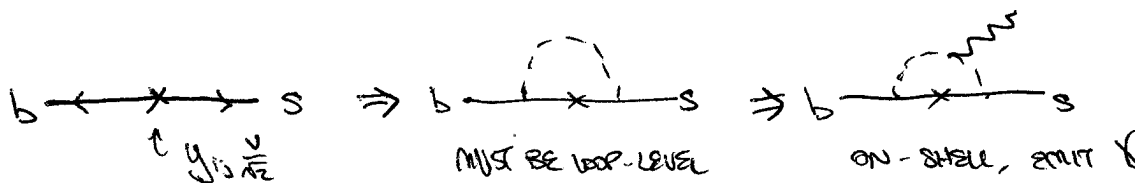
Decays will ~~involve FCNC~~ INVOLVE FLAVOR-CHANGING NEUTRAL CURRENT
 ↳ but we know FCNC NOT ALLOWED IN SM @ tree level



A POINT ABOUT PENGUINS:

NOTE: NP IN LOOPS APPEARING AS A DEVIATION FROM SM.

WANT: FCNC. DRAW THE EXT FERMION LINE:



IF YOU WANT TO BE ARGUMENTATIVE, YOU'D ARGUE THAT THIS PENGUIN NEEDS 'LEGS' (as H does for $B_s \rightarrow \mu\mu$). BUT YOU CAN SEE THAT THE PHYSICS IS ALREADY IN THE χ EMISSION.

BIG PIC: the $\mu\mu$ is there to make up for the $\Delta p \sim m_b - m_s$.

^{VERY}
A RARE DECAY: $B_{s,d} \rightarrow \mu\mu$ IS ONE OF THE REMAINING STRONGHOUDS OF THE SM. LEPTON SECTOR.
 IT IS BOTH THEORETICALLY INTERESTING (BSM) & HAS YET TO BE DISCOVERED (... CDF?)
 will discuss shortly \rightarrow who nice: CLEAN SIG, THY UNCERTAINTY ENCODED IN FB

THE SM BRANCHING RATIO FOR $B_{s,d} \rightarrow \mu\mu$ IS SUPPRESSED BY

- LOOP $\sim (4\pi)^{-2}$
 - GIM $\sim V_{id}^* V_{is} \Delta m_{ui}/M_W$
 - MASS INS. $\sim m_\mu^2$
- $\rightarrow B^0$ IS A PSEUDOSCALAR
 \rightarrow LEPTONS MUST HAVE SAME HELICITY
 \rightarrow BUT IN SM, WEAK CURRENT ONLY COUPLES TO LEFT CHIRALITY

$Br(B_s \rightarrow \mu\mu)_{SM} \approx 4 \times 10^{-9}$ \leftarrow I DUNNO, MAYBE IT'S 3×10^{-9}
 $Br(B_d \rightarrow \mu\mu)_{SM} \approx 2 \times 10^{-10}$... I DON'T REALLY CARE. DO YOU?

EFFECTIVE THEORY REMINDER

$M = \langle H_{eff} \rangle = \sum_i c_i \langle \mathcal{O}_i \rangle$

WILSON COEFFICIENTS \rightarrow ENCODES UV THEORY (DOF W/ $M > M_W$) INCLUDES RUNNING FROM HI SCALE TO M_W

EFFECTIVE OPERATOR ENCODES IR DOF (OPERATOR ASSUMED TO BE @ LOW SCALE; IN GENERAL ONE MUST INCLUDE \oplus MIXING EFFECTS INTO THE c_i)

ASSUMPTION: OPERATORS FACTORIZE INTO ~~SWARK~~ SWARK & LEPTON PIECES. OTHERWISE BARYON / LEPTON # VIOLATION (SUCH \mathcal{O} 'S SHOULD HAVE NEGLECTIBLE WILSON COEFFICIENTS)

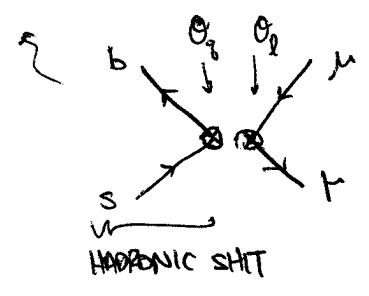
$\mathcal{O}_{XY}^V = \bar{b} \gamma^\mu P_X s \quad \otimes \quad F \gamma_\mu P_Y \mu$ VECTOR

$\mathcal{O}_{XY}^S = \bar{b} P_X s \quad \otimes \quad F P_Y \mu$ SCALAR

$\mathcal{O}_X^T = \bar{b} \sigma^{\mu\nu} P_X s \quad \otimes \quad F \sigma^{\mu\nu} P_Y \mu$ TENSOR

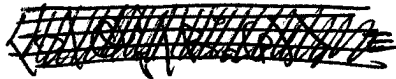
note: OTHER PEOPLE PARAMETERIZE IN TERMS OF V, A, S, P, T.

$P_{L,R} = \frac{1}{2}(1 \mp \gamma^5)$



$$\langle H\bar{F} | H_{\text{eff}} | B_s^0(P) \rangle = \sum_i C_i \underbrace{\langle H\bar{F} | \mathcal{O}_i^1 | 0 \rangle}_{\text{EASY}} \underbrace{\langle 0 | \mathcal{O}_i^1 | B_s^0(P) \rangle}_{\text{CONTAINS QCD EFFECTS!}}$$

THE $\langle 0 | \mathcal{O}_i^1 | B \rangle$ MATRIX ELEMENT IS NON-PERTURBATIVE \rightarrow IS PARAMETERIZED BY A DECAY CONSTANT



$$\langle 0 | \bar{b} \gamma^\mu P_{L,R} s | B_s^0(P) \rangle = F_B \frac{i}{2} P^\mu f_{B_s}$$

\uparrow BY LORENTZ INVARIANCE
 \uparrow ALL HADRONIC THY UNCEB
 \uparrow \neq CORRESPONDS TO $P_{L,R} = \frac{1}{2}(1 \mp \gamma^5)$
 \uparrow ONLY γ^5 CONTRIBUTES \leftrightarrow PSEUDOSCALAR



from this, contract w/ P_μ to obtain scalar matrix element

$$\langle 0 | \bar{b} P_{L,R} s | B_s^0(P) \rangle = \pm \frac{i}{2} \frac{M_{B_s} f_{B_s}}{M_b + M_s} \quad (1-P_L)$$

\leftarrow can also note that the scalar makes no contribution

check: $P_\mu \langle 0 | \bar{b} \gamma^\mu P_{L,R} s | B \rangle = \langle 0 | \bar{b} (\cancel{P}_b + \cancel{P}_s) P_{L,R} s | B \rangle$

$$= -M_b \langle 0 | \bar{b} P_{L,R} s | B \rangle + M_s \langle 0 | \bar{b} P_{L,R} s | B \rangle$$

$$= -(M_b + M_s) \langle 0 | \bar{b} P_{L,R} s | B \rangle + M_s \langle 0 | \bar{b} s | B \rangle$$

PSEUDOSCALAR

finally: $\langle 0 | \bar{b} \sigma^{\mu\nu} P_{L,R} s | B_s^0(P) \rangle = 0$

\uparrow by LORENTZ: CANT FORM ANTSYM. TENSOR OUT OF ONLY P^μ .

IT IS NOT ESPECIALLY ILLUMINATING TO GO OVER PARTICULAR DIAGRAMS, BUT IT IS WORTH MENTIONING THAT THE PHOTON PENGUIN VANISHES BY THE WARD IDENTITY:

$$\langle H\bar{F} | F \gamma^\mu P_{L,R} H | 0 \rangle \langle 0 | \bar{b} \gamma_\mu P_{L,R} s | B \rangle \sim P_\mu \langle H\bar{F} | F \gamma^\mu P_{L,R} H | 0 \rangle = 0$$

NOW YOU CAN JUST PLUG IN CHUG

$$Br(B_s^0 \rightarrow \mu\mu) = \frac{\Gamma_B}{16\pi} \left(\frac{|M|^2}{M_B} \right) \sqrt{1 - \left(\frac{M_\mu + M_\mu}{M_B} \right)^2} \sqrt{1 - \left(\frac{M_\mu - M_\mu}{M_B} \right)^2}$$

CONTAINS WILSON COEFFICIENTS

IF YOU WANT TO GEN. TO B → ll'

PARAMETRICALLY:

$$Br(B_s^0 \rightarrow \mu\mu) = 3.5 \times 10^{-9} \left[\frac{\Gamma_{B_s}}{1.6 \text{ ps}} \right] \left[\frac{f_{B_s}}{210 \text{ MeV}} \right]^2 \left[\frac{|V_{ts}|}{0.040} \right]^2 \left[\frac{M_\mu(M_\mu)}{170 \text{ GeV}} \right]^{3/2}$$

POSSIBLY FOR MASS

BUCHANAN & BURAS NUCL PHYS B400 225 (1993)

That's all I want to say about the SM.

BSM: 2HDM (TWO HIGGS DOUBLET MODEL) eg. SUSY (MSSM)

GENERIC 2HDM ARE VERY CONSTRAINED BY FCNC @ TREE LEVEL

IF H_u COUPLES TO BOTH U_L & D_R THEN CANNOT DIAGONALIZE

WHAT WE REALLY CARE ABOUT

SO: TYPICALLY HAVE TO IMPOSE SYMMETRIES TO KILL FCNC

TYPE II 2HDM: H_u COUPLES ONLY TO U_R, H_d COUPLES ONLY TO D_R

$$\mathcal{L} = y_u \bar{Q} H_u U_R + y_d \bar{Q} H_d D_R$$

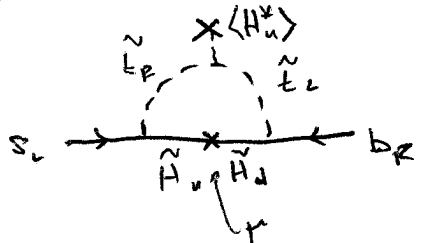
Remark: 2HDM PREDICTED BY MSSM, this is the natural set up.

↳ ANOMALY CANCELLATION & HOMOMORPHY → U(1)

√ g H_uH_d → HAS TO BE THERE BE EWSD

BUT: ~~SUSY~~ TERMS BREAK THIS STRUCTURE (EVEN HOMOMORPHIC PRET.)

eg:



COUPLES H_u TO b_R!

[SEE KANE, KOLDA, LENDNER 103]

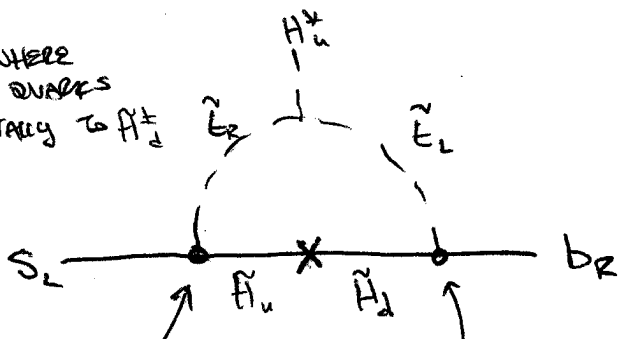
SO, AFTER EWSS :

(IGNORE 1ST GENERATION FOR SIMPLICITY)

$$Y_{1-loop}^{mass} = \begin{pmatrix} \bar{S}_R & \bar{b}_R \end{pmatrix} \begin{pmatrix} M_s & 0 \\ y_b E V_u & M_b \end{pmatrix} \begin{pmatrix} S_L \\ b_L \end{pmatrix}$$

INDUCED BY LOOP DIAGRAM
E INCLUDES $y_t V_s$ & FINEMATIC (LOOP) FACTORS

IN BASIS WHERE
DOWN-TYPE QUARKS
COUPE DIAGONALLY TO H_d^\pm



LOOP LEVEL
FLAVOR CHANGES
& WIRING-LIBBS
COUPLING. THE
SIGNIFICANCE OF
THE LATTER IS
THIS

$y_t V_s = \epsilon$
OFF DIAG.

y_b COUPLING, DIAGONAL
[ALSO: \tilde{g} DIAGRAMS w/ S^3 INS.]

YUHSIN IS AN EXPERT

DIAGONLIZE LOOP MASS MATRIX
ROTATE BY AN ANGLE

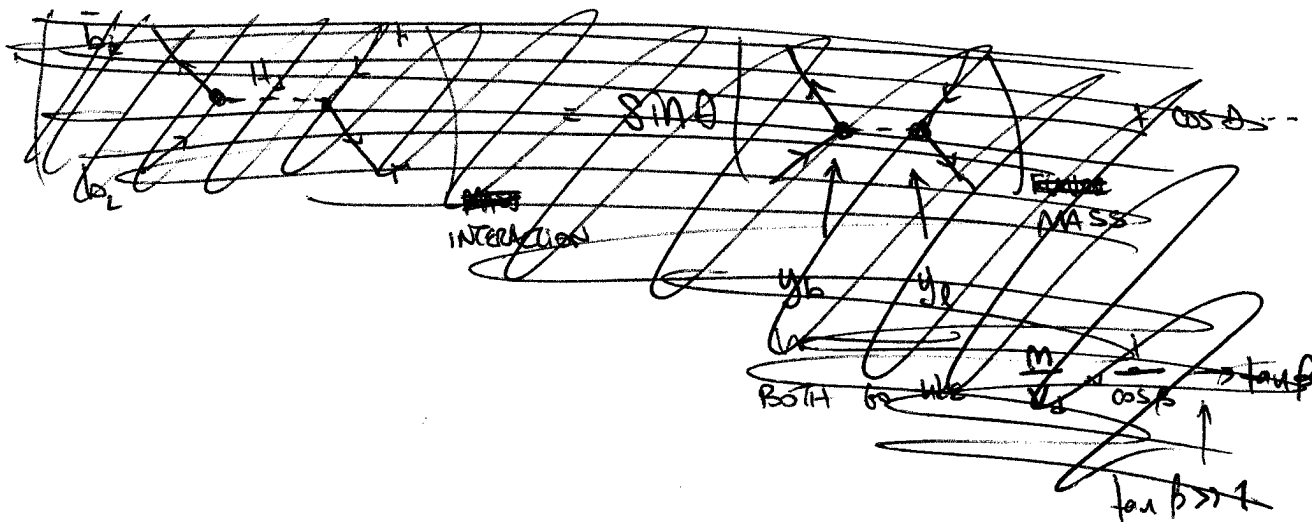
$$\sin \theta \approx y_b E V_u / M_b = \epsilon \frac{V_u}{V_d} = \epsilon \tan \beta$$

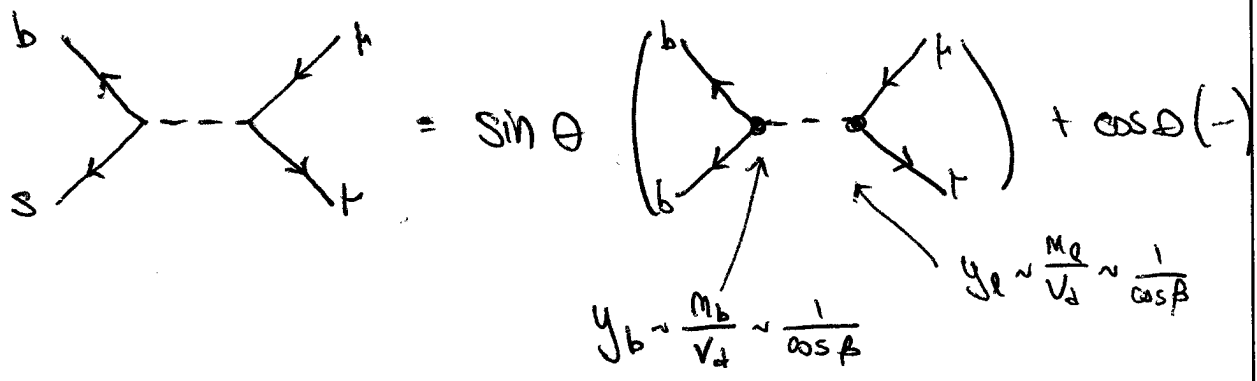
the phase \rightarrow MASS
ROTATION IN QR
TOY 2 FLAVOR QVE

$$M_b = y_b V_d$$

PARAMETER OF ZHM
CAN OVERCOME
LOOP SUPPRESSION

tan β DEPENDENCE :





$$\sim \sin \theta \frac{1}{\cos^2 \beta} \xrightarrow{\tan \beta \gg 1} \boxed{\tan^3 \beta}$$

LARGE ENHANCEMENT

Remark: large $\tan \beta$ favored by soft GUT
 ↳ want to make $y_t \approx y_b$ @ GUT scale

CONNECTIONS TO OTHER OBSERVABLES

3 HINTS FOR MODEL BUILDING

• FLEISCHER, SERRA, TUNING (2010)

MAIN SOURCE OF SYSTEMATIC UNCERTAINTY

WANT TO BEAT THIS DOWN.

$$\frac{Br(B_s \rightarrow \mu\mu)}{Br(B_c \rightarrow X)} = \frac{f_b}{f_s} \frac{\epsilon_x}{\epsilon_{\mu\mu}} \frac{N_{\mu\mu}}{N_X}$$

↑ DECAY CONST ↑ DET. EFF. ↑ # EVENTS

UCLB EXTRACTS $B_s \rightarrow \mu\mu$ BY NORMALIZING w/rt

$$B_{q_1} \rightarrow X = \begin{matrix} B_{q_1}^+ \rightarrow (3/4) K^+ \\ B_{q_1}^0 \rightarrow K^+ \pi^- \\ B_{q_1}^- \rightarrow (3/4) K^0 \end{matrix}$$

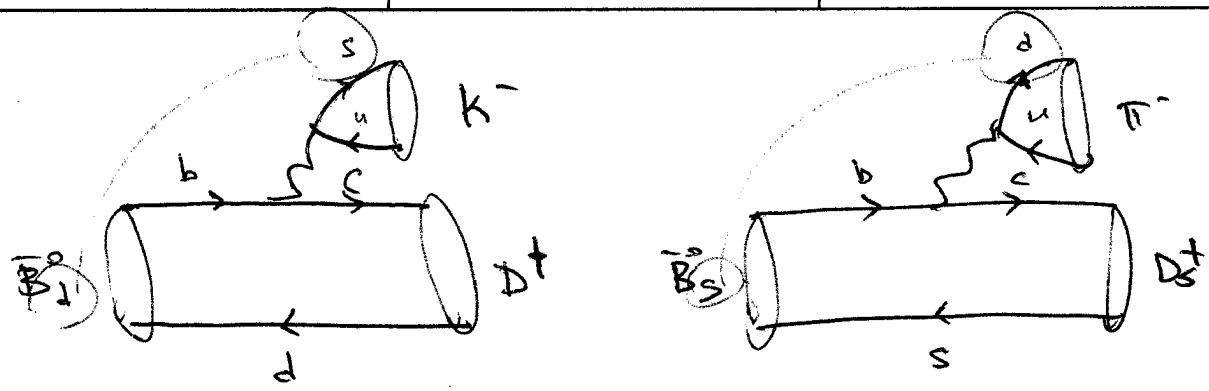
PROPOSE: NEW WAY TO OBTAIN f_d/f_s @ UCLB:

PICK: $B_s \rightarrow X_1$; $B_d \rightarrow X_2$

SUCH THAT

- ① RATIO OF BRANCHING RATIOS EASY TO MEASURE @ UCLB
→ decay into 2 charged particles
- ② DECAYS ROBUST AGAINST NP CONTRIB.
→ NO PENGUINS (WANT: tree-level)
- ③ RATIO OF BR'S THEORETICALLY WELL UNDERSTOOD

$$\Rightarrow \begin{matrix} \bar{B}_s^0 \rightarrow D_s^+ \pi^- \\ \bar{B}_d^0 \rightarrow D^+ K^- \\ \quad \uparrow \text{(CHARGED MESONS)} \end{matrix}$$



RELATED BY U-SPIN \subset SU(3)_F
 NO EXCHANGE TOPOLOGIES \rightarrow B_{\perp} MUST GO TO $(u, s) = K^-$
 B_s \rightarrow $(u, d) = \pi^-$

UP TO U-SPIN BREAKING EFFECTS

CLAIM: CAN INCREASE NP REACH BY A FACTOR OF 2.

RELATION TO $B_s - \bar{B}_s$ MIXING

GOLOWICZ, HEWITT, PAUVASA, PETROV (2009)
 YEGHIYAN (2010)

ALSO: ANUP, SANKAR '07
 BURAS hep-ph/0303067

IDEA: WRITE OUT EFF \mathcal{L} FOR $g \} l$ SECTORS.
 MUST ASSUME A CLASS OF NP MODEL
 (eg INTERMEDIATE VECTOR, FCNC HIGGS, ...)
 COMPARE B_s MIXING CONSTRAINT TO $B_s \rightarrow \mu\mu$

eg. Z' w/ FCNC

$$\Delta M_{B_s}^{(Z')} \sim (\text{factors}) \frac{g_{Z'sb}^2}{M_{Z'}^2} \Rightarrow \frac{g_{Z'sb}^2}{M_{Z'}^2} < 2.5 \times 10^{-11} \frac{1}{\text{GeV}^2}$$

$$\text{Br}(B_s \rightarrow \mu\mu) \sim (\text{factors}) \frac{g_{Z'sb}^2 g_{Z'\mu\mu}^2}{M_{Z'}^4} \Rightarrow \text{Br}(B_s \rightarrow \mu\mu) < 0.25 \times 10^{-9} \left(\frac{\text{TeV}}{M_{Z'}}\right)^2$$

REMARKS: THIS STRATEGY IS MOTIVATED BY A STUDY BY THE AUTHORS CONNECTING D MIXING W/ D → MM.

D SECTOR: IF ALL/MOST OF MIXING FROM NP, CAN OFTEN PREDICT D → MM BECAUSE E AMPLITUDES CAN HAVE SAME PARAMETERS

SM MIXING HAS LARGE THY UNCERTAINTIES & MANY NP MODELS CAN PRODUCE THE OBS. MIXING.

B SECTOR: VERY DIFFERENT! SM MIXING AGREES W/ OBS VALUE

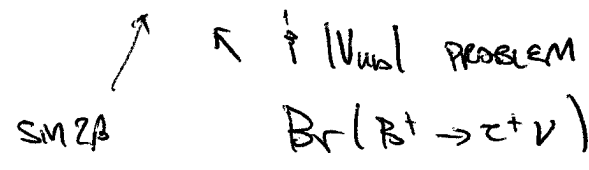
CAN ALSO HAVE INTERFERENCE WHICH MAKES $B_T < B_{SM}$



$$\frac{\Delta M_{B_s}^{NP}}{\Delta M_{B_s}^{SM}} \leq 20\%$$

↳ STRONG CONSTRAINT IN SOME NP MODELS!

OTHER OBSERVABLES: $\epsilon_K - S_{\text{PKS}}$ TENSION



$$S_{\text{PKS}}, \begin{matrix} K^+ \rightarrow \pi^+ \nu \bar{\nu} \\ K_L \rightarrow \pi^0 \nu \bar{\nu} \end{matrix}$$

↳ ALL CORRELATED. SOME HINTS @ DISAGREEMENT W/ EACH OTHER