SK FORCE M EUROPEAN FUSION DEVELOPMENT AGREEME



Modelling of Frequency Sweeping with the HAGIS code

S.D.Pinches¹

H.L.Berk², S.E.Sharapov³, M.Gryaznavich³

¹Max-Planck-Institut für Plasmaphysik, EURATOM Assoziation, Garching, Germany ²Institute for Fusion Studies, University of Texas at Austin, Austin, Texas, USA ³EURATOM/UKAEA Fusion Association, Culham Science Centre, Abingdon, UK

UKAEA Fusion



Structure of Talk

- What is frequency sweeping?
 - Experimental evidence
 - Theoretical understanding
- Numerical modelling
 - Description of the HAGIS code
 - Simulations of frequency sweeping
- Summary

Experimental Observations

Frequency sweeping in MAST #5568







Simon Pinches, 8th IAEA Technical Meeting on Energetic Particles, San Diego



 Universality in nonlinear response of resonant particles to low amplitude wave

[Berk, Breizman, Pekker (1997)]

- Particle distribution satisfies a 1-dimensional equation (two phase-space coordinates)
- Constants of motion for wave $E(r,t) = C(t) E(r,\theta,n\phi - \omega_0 t)$
 - Magnetic moment, μ (if $\omega_0 \ll \omega_c$ and $L_{\omega} > \rho_i$)
 - Energy in rotating frame, H' = H (ω_0/n) P_z (if 1/C dC/dt « ω_0)

Wave-Particle Interaction

Define:
$$\Omega_l(P_{\phi}, H', \mu) = n \langle \omega_{\phi} \rangle - l \langle \omega_{\theta} \rangle$$

As P_{ϕ} changes due to interaction at fixed H', μ
 $P_{\phi} - P_{\phi,l} = \frac{\partial P_{\phi}}{\partial \Omega_l}\Big|_{H',\mu} \left[\Omega_l(P_{\phi}) - \omega(t)\right]$
Equations of particle motion for fixed H', μ
 $\frac{d\xi}{dt} = \Omega_l - \omega_0, \quad \frac{d\Omega_l}{dt} = -\omega_{bl}^2(t) \sin\xi$
Hence, $\frac{d^2\xi}{dt^2} + \omega_{bl}^2(t) \sin\xi = 0$
"Pendulum equation"
F is a phase space dependent form factor
Trapping frequency, $\omega_{bl}(t) \propto |E|^{1/2}F(H', \mu)$





• Frequency sweep is related to trapping frequency [Berk et al., (1997)]

$$\delta\omega\propto\omega_b^{3/2}t^{1/2}$$

Amplitude related to frequency sweep

$$\Rightarrow \quad \delta B \propto \left(\frac{\delta \omega^2}{t}\right)^{2/3}$$

Analytic estimates give correct order of magnitude. Numerical simulation required for more accurate estimate.





- Use experimentally observed rate of frequency sweeping to determine wave amplitude
- In general, numerical modelling is needed to establish the form factor that relates $\delta\omega$ and δB
- Validate HAGIS for model case
- Employ HAGIS to establish δB in general case
 - General geometry (including tight-aspect ratio)
 - Mode structure: global mode analysis

The HAGIS Code







- Straight field-line equilibrium
 - Boozer coordinates
- Hamiltonian description of particle motion [White & Chance 1984]
- Fast ion distribution function
 - δf method
- Evolution of waves
 - Wave eigenfunctions computed by CASTOR

Equilibrium Representation

 Coordinates ψ_p, θ, ζ chosen to produce straight field lines



 $\mathbf{B} = \delta(\psi_p, \theta) \nabla \psi_p + I(\psi_p) \nabla \theta + g(\psi_p) \nabla \zeta,$ $\mathbf{B} = \nabla \psi \wedge \nabla \theta - \nabla \psi_p \wedge \nabla \zeta,$

 $\Rightarrow \mathbf{A} = \psi \nabla \theta - \psi_p \nabla \zeta.$

Particle Description

Exact particle Lagrangian, $\mathcal{L}_{exact} = \sum_{ep} \frac{1}{2}mV^2 + e\mathbf{V} \cdot \mathbf{A} - e\phi$ is gyro-averaged and written in the form,

$$\mathcal{L}_{ep} = \sum_{j=1}^{n_p} P_{\theta j} \dot{\theta}_j + P_{\zeta j} \dot{\zeta}_j - \mathcal{H}_j$$

with

$$\mathcal{H}_j = \frac{1}{2}m_j v_{\parallel j}^2 + \mu_j B_j + e_j \phi_j$$





Fast Particle Orbits

- ICRH ions in JET deep shear reversal
 - On axis heating[†]:
 - $\Lambda = \mu B_0 / E = 1$
 - -E = 500 keV
- Produces predominately potato orbits



[†]J. Hedin, Thesis 1999

Distribution Function

- Represented by a finite number of *markers*
- Markers represent deviation from initial distribution function - so-called δf method
 - Dramatically reduces numerical noise

$$f = \underbrace{f_0(\mathcal{E}, P_{\zeta}; \mu)}_{\text{analytic}} + \underbrace{\delta f(\Gamma^{(p)}, t)}_{\text{markers}}$$
$$\frac{df}{dt} = 0 \Rightarrow \dot{\delta f} = -\dot{P}_{\zeta} \frac{\partial f_0}{\partial P_{\zeta}} - \dot{\mathcal{E}} \frac{\partial f_0}{\partial \mathcal{E}}$$
$$\int f g \, d\Gamma^{(p)} \longleftrightarrow \int f_0 g \, d\Gamma^{(p)} + \sum_{j=1}^{n_p} \delta n_j g_j$$
where
$$\delta n_j(t) \equiv \delta f_j(t) \, \Delta \Gamma_j^{(p)}(t) \qquad \text{Parker \& Lee, 1993}$$

Wave Equations

- Linear eigenstructure assumed invariant
- Introduce slowly varying amplitude and phase:
- $\tilde{\Phi}_{k} = A_{k}(t) \sum_{m} \tilde{\phi}_{km}(\psi) e^{i(n_{k}\zeta m\theta \omega_{k}t \alpha_{k}(t))}$ Gives wave equations as:

$$\begin{aligned} \dot{\mathcal{X}}_k &= \frac{1}{2E_k} \sum_{j=1}^{n_p} \delta f_j \Delta \Gamma_j^{(p)} \sum_m (k_{||m} v_{||j} - \omega_k) S_{jkm} + \mathcal{X}_k \gamma_d, \\ \dot{\mathcal{Y}}_k &= -\frac{1}{2E_k} \sum_{j=1}^{n_p} \delta f_j \Delta \Gamma_j^{(p)} \sum_m (k_{||m} v_{||j} - \omega_k) C_{jkm} + \mathcal{Y}_k \gamma_d, \end{aligned}$$

where

$$\begin{array}{lll} \mathcal{X}_k &\equiv & A_k \cos(\alpha_k), \\ \mathcal{Y}_k &\equiv & A_k \sin(\alpha_k), \end{array} & \begin{array}{lll} C_{jkm} &\equiv & \Re e[\tilde{\phi}_{km}(\psi_j)e^{i\Theta_{jkm}}] \\ &S_{jkm} &\equiv & \Im m[\tilde{\phi}_{km}(\psi_j)e^{i\Theta_{jkm}}] \\ &\Theta_{jkm} &\equiv & n_k \zeta_j - m\theta_j - \omega_k t \end{array}$$

Simon Pinches, 8th IAEA Technical Meeting on Energetic Particles, San Diego

Additional mode damping rate, γ_d

HAGIS Code Performance



- HAGIS code parallelises very well
 - relatively low level of inter-processor communication traffic

Self-Consistent Frequency

Sweeping



Linear Growthrate



...with additional damping





Fourier spectrum of evolving mode



Linear Growthrate









Fourier spectrum of evolving mode



Simon Pinches, 8th IAEA Technical Meeting on Energetic Particles, San Diego

23

тby

Fast Ion Redistribution



MAST #5568



• Obtain factor relating ω_{b} and δB



Particle Trapping in MAST

- Particles trapped in TAE wave
 - All particles have same
 - $H' = E \omega/n P_{\zeta}^{\circ}$ = 20 keV
 - TAE amplitude: $\delta B/B = 10^{-3}$



Scaling of Nonlinear Bounce

Frequency



Scaling of Nonlinear Bounce

Frequency





For monotonic q-profiles we now know:

 $\omega_b = C_1 \left(\frac{\delta B}{B}\right)^{1/2}$ where $C_1 = 1.156 \times 10^6$

• For a single resonance, $\delta \omega = C_2 \omega_b^{3/2} t^{1/2}$

where
$$C_2 = \frac{\pi}{2\sqrt{2}} \simeq 1$$

• Therefore, $\frac{\delta B}{B} = \frac{1}{C_1^2} \left(\frac{\delta \omega^2}{C_2^2 t}\right)^{2/3}$

TAE Amplitude in MAST



Conclusions



 Frequency sweeping has been modelled using the HAGIS code

- Benchmarked against analytic theory

• The amplitude of a frequency sweeping mode in MAST has been calculated to be $\delta B/B = 4 \times 10^{-4}$