Resolving the Physics of Relaxation: Flows, Reconnection, Heating, EUV bursts, and Waves



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Relaxation as a dynamic cascade



Double Solar Loop Experiment





3D B-field Measurements

Vector B-field measured at 2700 positions with ~2cm spatial resolution

15 θ -positions, 10 z-positions 5 shots per position Set of 750 shots



3D Measurement Volume



3D Laboratory Vector Measurements of \mathbf{B} and \mathbf{J}



Validation of Probe Current Density



3D JxB Force Vectors

JxB Vectors plotted at central plane

Axial JxB forces near footpoint

Axial cross-sections



• To compare with theory, we extract cross-sections along the axis 8

Force cross-sections





- Out of plane forces shown with colormap
- In-plane forces shown with black vectors
- 50% current radius traced with dotted line

Theory for MHD Axial Force

Any flared current channel (finite J_r) will have an axial JxB force in the direction of larger minor radius

$$(\vec{J} \times \vec{B}) \cdot \hat{z} = J_r B_\phi$$

This is equivalent to the component of magnetic pressure in the axial direction:





Pressure Estimate from Perpendicular Forces

Assuming force balance along the minor radius, we have:

Force balance in minor radius
$$P(r) = \int_{\infty}^{r} -\mathbf{J} \times \mathbf{B} \cdot \hat{r} |dr'|$$

Due to asymmetry of measurement data, we integrate along multiple different rays to bound the on-axis pressure. Rays parallel to hoop force will overestimate pressure, rays anti-parallel to hoop force will underestimate pressure.



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Number Density



Assuming relatively constant temperature (1-3 eV), we can calculate the number density:

$$\frac{N}{V} = \frac{P}{RT}$$

This gives reasonable bounds on the number density present in the flux tube.

- 300x increase in density
- Background gas cannot account for more than first factor of 10.
- Most density must originate from footpoints

Strong mass flux from boundary!

Inferred Flow Speeds



Since we have estimates of density as a function of space and time, we can use the continuity equation to estimate the velocities needed to produce the observed mass flux.

$$\frac{\partial \rho}{\partial t} = -\nabla \cdot (\rho \mathbf{u}) \approx -\left(\frac{\partial \rho}{da}u_a + \frac{\partial u_a}{da}\rho\right)$$

Assuming the compressible term is small near the electrodes:

$$u_0 \approx -\left(\frac{\partial \rho}{\partial t} \left(\frac{\partial \rho}{\partial a}\right)^{-1}\right)\bigg|_{a=0}$$

Axial flows observed in similar solar experiment for same currents have the same magnitude (~ 5 cm/us).

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Stenson, Eve Virginia, and P. M. Bellan. "Magnetically driven flows in arched plasma structures." PRL 109.7 (2012): 075001.

Axial Forces Consistent with Observed Flows



- Measurements of axial and perpendicular forces give two self-consistent estimates of axial velocity
- Inferred velocity from imaging gives same order of magnitude

 Alfvenic axial flows can be driven by axial MHD forces

Axial Field Collimation



- Axial field increases by factor of 10 during lifetime
- Factor of 3 radial compression (self-similar)

Summary

(1) Laboratory plasma experiment (20 kA, 3kV, 10us, 1m)

(2) 3D MHD force measurements demonstrate simple self-consistent mechanism for mass flux from flared current channel.

(3) Mass flux relevant to dynamics of coronal loops, astrophysical jets, and compact tori.

Relaxation as a dynamic cascade



Formation of plasma jet



Ideal MHD kink instability



- As jet length exceeds a critical value plasma becomes kink unstable
- Can be explained by Kruskal-Shafranov criterion:
 - $q = rB_z/RB_{\phi} = 2\pi rB_z/LB_{\phi}$
 - *L* increases \rightarrow *q* decreases: kink occurs when *q* < 1

Rayleigh-Taylor (RT) instability



- Kink instability induces effective gravity in the r direction
- RT instability occurs at interface between heavy plasma and light vacuum
- Finger-like structures: RT ripples.

Kink-induced Rayleigh-Taylor (RT)

- Kink-induced Rayleigh-Taylor instability severs plasma jet.
- We simultaneously observe:
 - 1) Strong, localized 25-40 nm EUV bursts
 - 2) <u>Doppler and Stark broadenings</u> in plasma emission spectra
 - 3) 2-20 MHz polarized B-field fluctuations
 - 4) Strong and transient voltage spikes
 - 5) Sudden changes in B-field profile

 Conclusion: Hall-MHD magnetic reconnection occurs and causes
 Ohmic electron heating, stochastic ion heating, whistler wave emission κ.-в





Not the same shot

K.-B. Chai, X. Zhai, P. Bellan, Phys. Plasmas 23, 132122 (2016).

Diagnostics for investigating RT signatures



Simultaneous EUV and RF bursts

- Observe transient EUV bursts & high frequency fluctuations when RT occurs:
 - EUV bursts: could be associated with electron heating
 - High frequency fluctuation: could be associated with whistler waves
- Goal: study details of magnetic reconnection occurring in our jet by using comprehensive diagnostics



K.-B. Chai & P. Bellan, Rev. Sci. Instrum. 84, 123504 (2013)





K.-B. Chai and P. M. Bellan, Rev. Sci. Instrum. 84, 123504 (2013).

EUV (25-40 nm) vs visible light

End-on view

red: EUV; blue: visible light



- As RT instability grows:
 - EUV (red) becomes extremely bright in localized area
 - Visible light (blue) becomes dark where EUV gets bright
- \rightarrow Ratio of EUV to visible light becomes extremely high
- \rightarrow Higher ionization state when RT occurs (electron heating)

Electron heating: Spectroscopic line ratios



- As RT occurs:
 - Ar II (Ar+) lines disappear and Ar III (Ar2+) lines dominate over Ar II lines
 - Ar IV (Ar³⁺) lines are also observed
- \rightarrow Indicates higher ionization state and electron heating

Ion heating: Doppler & Stark broadening



- Voigt fitting gives both Doppler (T_i) and Stark (n_e) broadenings:
- As RT occurs
 - $T_i: 2.6 \pm 0.4 \text{ eV} 15.8 \pm 2.3 \text{ eV}$
 - $n_{\rm e}$: (1.6 ± 0.3)×10²² m⁻³ (5.1 ± 2.1)×10²² m⁻³

High frequency B-dot probe

• B-dot probe with excellent shield



Whistler waves: Magnetic field fluctuations



- As RT occurs, broadband (2-20 MHz) high frequency magnetic fluctuations measured by whistler probe.
 - Ion cyclotron freq. < observed wave freq. < electron cyclotron freq.
 - Have power-law dependence on freq. (~ f^{-1}) but not turbulence because of coherence between different frequencies

Whistler waves: Circular polarization



- Hodograms of magnetic vector show circular polarization: confirms whistler wave character
 - Angle between waves and background magnetic field: < 60°

Observation of whistler waves suggests Hall-MHD reconnection

Voltage spikes



- When RT breaks the jet:
 - Reproducible, >500 V voltage spikes lasting ~1 µs appear
- Interpretation: sudden break of jet current induces strong emf

Conclusion

- Hall-MHD reconnection is believed to occur because
 1) jet diameter is observed to be similar to ion skin depth
 - 2) whistler waves are observed

3) Hall term and resistive term are calculated to have same order of magnitude

- Electron heating is likely caused by Ohmic dissipation
- Ion heating plausibly results from stochastic heating
- Voltage spikes are likely caused by sudden change in magnetic flux linking the electrode circuit

Relaxation as a dynamic cascade



Questions?

Experimental Parameters



Scintillating material: YAG:Ce Crystal vs powder





powder

	Crystal	powder
Output wavelength	550 nm	
Decay time	70 ns	
Thickness	100 µm	15 μm
Al coating thickness	200 nm	50 nm
Converting efficiency	1	3
X-ray resistance	strong	weak

 \rightarrow Powder type has better efficiency because it is thin

Mo:Si multilayer mirror



- Forms EUV image onto scintillator
 - Concave shape with f=50 mm; diameter 50.4 mm
 - Multilayered to obtain high reflectivity for 25-41.3 nm EUV
 - Maximum reflectivity 13% @ 36.5 nm (34 eV) with normal incident
 - Custom manufactured by NTT-AT (Japan)
- Principal: stack Mo/Si periodically to obtain constructive interference
 - Periodicity: satisfy Bragg constructive interference condition
 - $ightarrow d\sin\theta = \lambda / 2$ (d: periodicity, θ : angle to mirror surface, λ : wavelength)

Efficiency of EUV optics





Photon energy (eV)

- Efficiency = mirror reflectivity × AI transmittivity × scintillator efficiency
- Maximum efficiency: ~ 0.25 % @ 36.5 nm (34 eV)
- Average efficiency for 24.8-41.3 nm (30-50 eV) EUV: 0.15 %

→ Efficiency looks low but enough for us because our plasma is bright

High frequency B-dot probe

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Probe pair facing opposite direction



- Differential component $V_{d} = V_{1} V_{2} = 2A \frac{dB}{dt}$: magnetic signal
- Common component $V_c = V_1 + V_2 = 2V_{common}$: unwanted common pickup, e.g., capacitively coupled between probe and plasma

RF ground loop diverting



R. J. Perkins and P. M. Bellan, AIP Conf. Proc. 1406, 531 (2011).



Examples of EUV images



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