



# 2016 CTW

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## Multitple plasmoid formation and flux closure during transient-CHI start-up process on HIST

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### Outline

- 1) Introduction
- 2) HIST device and diagnostics
- 3) Experimental topics
  - a) Kink instability and relaxation in T-CHI generated plasmas
  - b) Flux closure and plasmoid reconnection
- 4) Summary

# Progress on Helicity Injection Experiments

The **helicity injection** is a promising candidate method for the non-inductive steady-state current drive and plasma start-up.

## Electrostatic HICD

Start-up

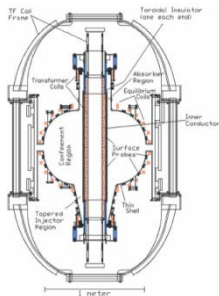
Steady-state

**Transient CHI**

**Muti-plusing CHI**

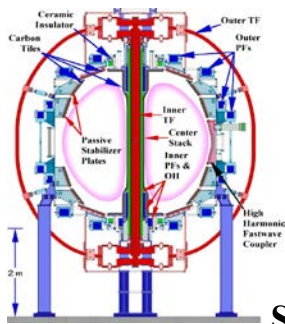
ST (HIT II, NSTX, HIST, SPECTOR, QUEST)

Spheromak(SSPX) ST(HIST)

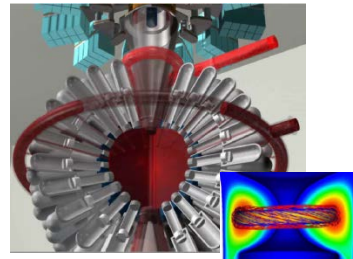


HIT-II

(U. Washington)



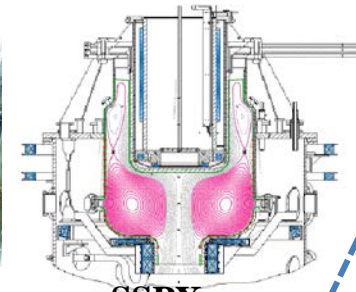
NSTX  
(PPPL)



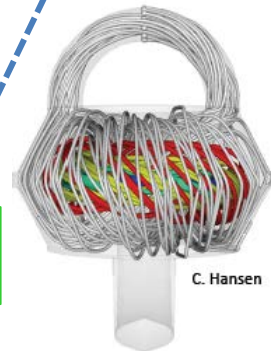
SPECTOR (General Fusion)



HIST  
(U. Hyogo)



SSPX  
(LLNL)

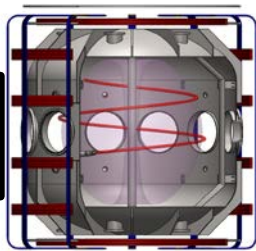


HIT-SI  
(U. Washington)

**LHI**

ST (Pegasus)

Pegasus  
(U. Wisconsin)



QUEST (Kyusyu U.)

**Inductive HICD**

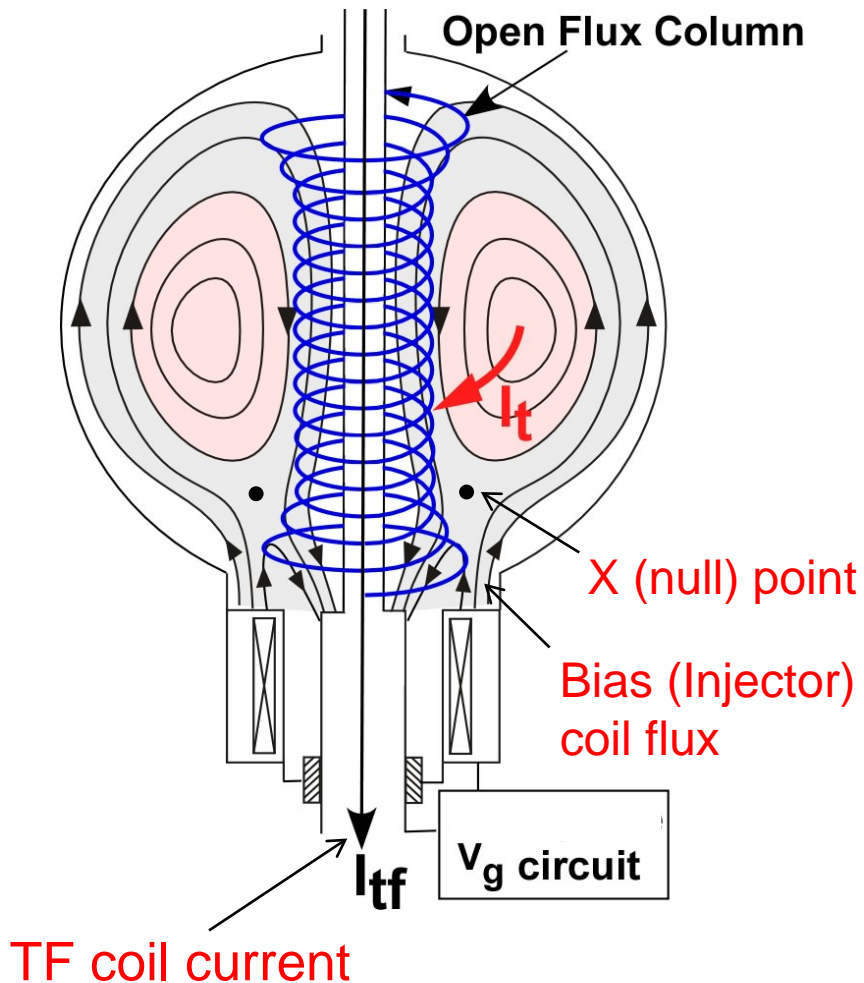
Spheromak (HIT-SI)

C. Hansen



# Key features of T-CHI method

## Gun-spheromak with Center conductor and External toroidal field; Spherical Torus



- T-CHI plasmas are generated by nearly an axisymmetric process.
- Any dynamo mechanism such as non-axisymmetric relaxation is not necessary, but only fast magnetic reconnection is required for the formation of the X-point during the short time scale of the start-up.
- Injection current in the OFC region ✓ the Kruskal-Shafranov limit
- T-CHI experiments provides a good platform for studying the MHD relaxation and reconnection physics.

# What are key issues in T-CHI experiments ?

## ■ Formation of closed flux surfaces (flux closure)

- The magnetic reconnection for the flux closure should be completed as soon as possible. How fast the magnetic reconnection occurs during the start-up time scale ?
- The reconnection rate slows down at a **large toroidal guide field and/or a high S number**. What is the fast reconnection mechanism ?

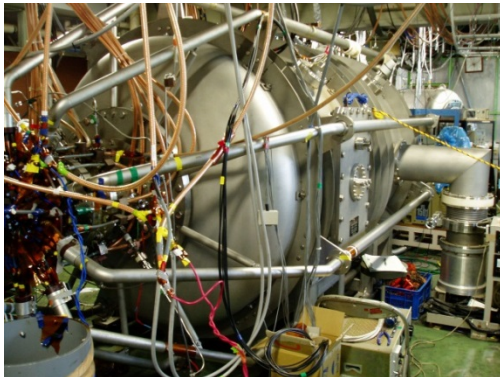
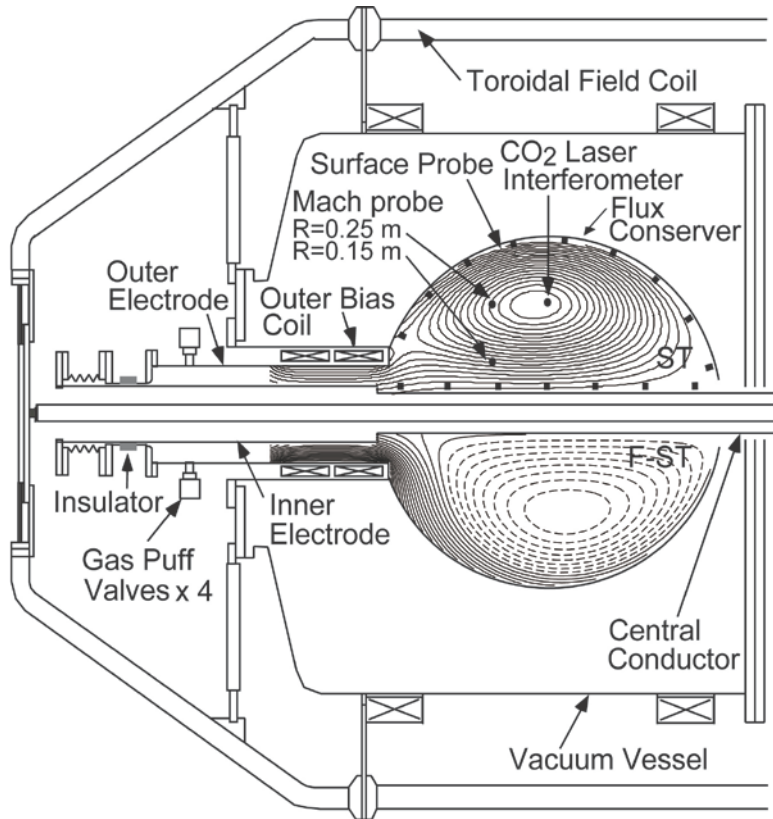
## ■ Optimization of current density profile

- The **excess edge injection current** causes the **n=1 kink instability** on the OFC. **Absorber arc** sometimes occurs at the rear gap of the FC. How to control the current profiles ?

## ■ Verification of CHI and LHI scaling $\frac{dK}{dt} = -K/\tau_K + 2V_{gun}\Psi_{bias}$

- **Helicity balance** based on the **helicity conservation law** in the presence of high TF should be experimentally verified.

# HIST device



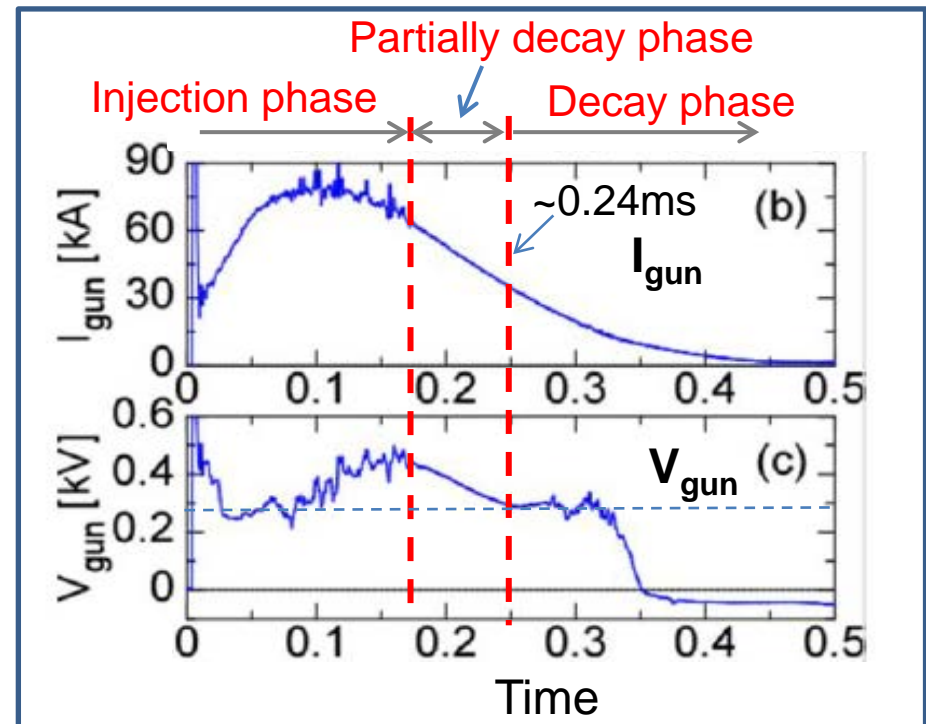
- **HIST parameters**

$R=0.3\text{ m}$ ,  $a=0.24\text{ m}$ ,  $A=1.25$

TF coil current  $I_{\text{tf}}=150\text{-}250\text{ kA}$

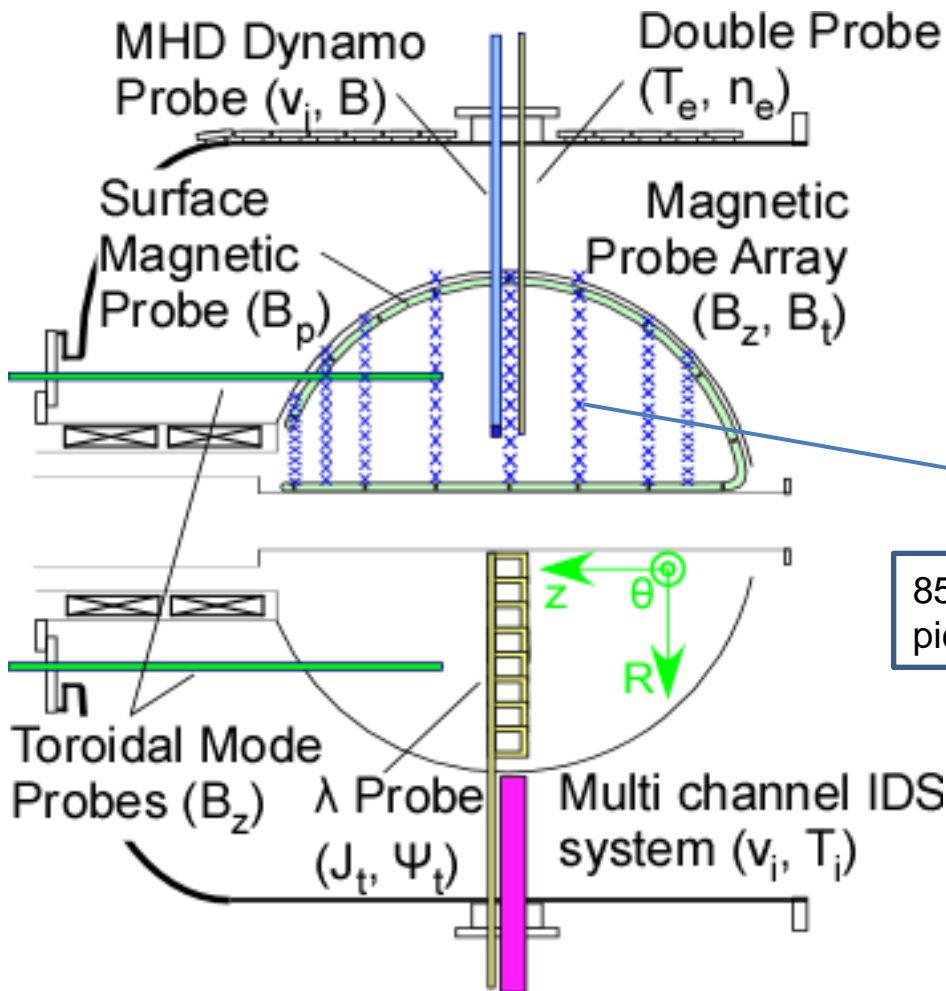
- **T-CHI capacitor banks**

Voltage :  $V = 5\text{ kV}$ ,  $C = 2.9\text{ mF}$



# Diagnostics

## Diagnostics



- ❑ Surface magnetic probes:  $I_{tor}$
- ❑ Internal magnetic probes:  $B_\theta, B_z, (B_r)$
- ❑  $\lambda$  probe:  $J_t, \Psi_t$
- ❑ Dynamo probes:  $V_\theta, V_z, V_r, B_\theta, B_z, B_r$
- ❑ Hall probe:  $J_\theta, J_z, J_r$
- ❑ Double electrostatic probe:  $T_e, n_e$
- ❑ Ion Doppler Spectroscopy:  $T_{D,i}$
- ❑ Electric field probe:  $E_r$
- ❑ CO<sub>2</sub> laser interferometer  $\langle n_e \rangle$

85 x  $B_z$  and 85 x  $B_t$   
pick-up coils

### Channel distance

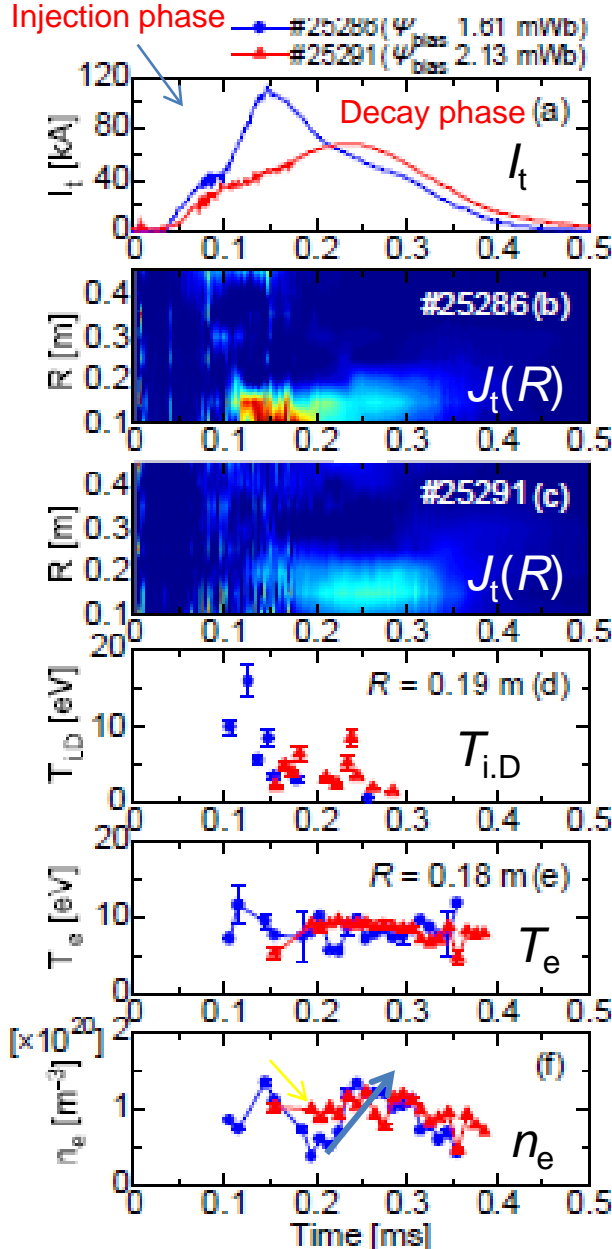
$\Delta R$	20 - 50 mm
$\Delta Z$	74 or 148 mm

Poloidal flux

$$\Psi_p = \int_0^R 2\pi r B_z(t) dr$$

by assuming axisymmetry

# Characteristics of T-CHI generated ST plasmas



$I_{TF} = 135$  kA

**Toroidal plasma current**

**Low bias flux**

Blue line : Low bias flux

Red line : High bias flux

**Toroidal current density**

**High bias flux**

**Doppler ion temperature**

**Electron temperature**

**Electron density**

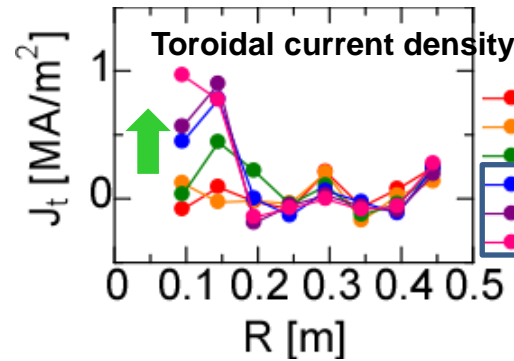
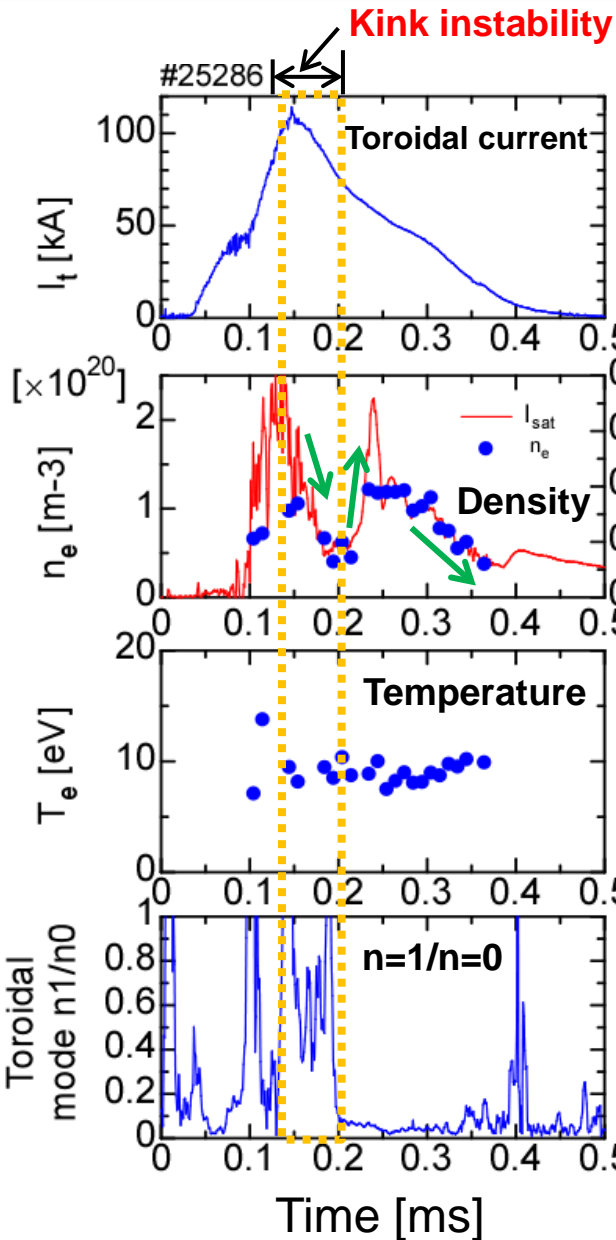
□ Typical T-CHI discharge characterized with varying the amount of the bias (injector) flux.

□ In the low bias flux case, the rise time of  $I_t$  is faster and its peak value is higher compared to the high bias case. Time evolution of  $I_t$  is divided by the injection phase and the decay phase.

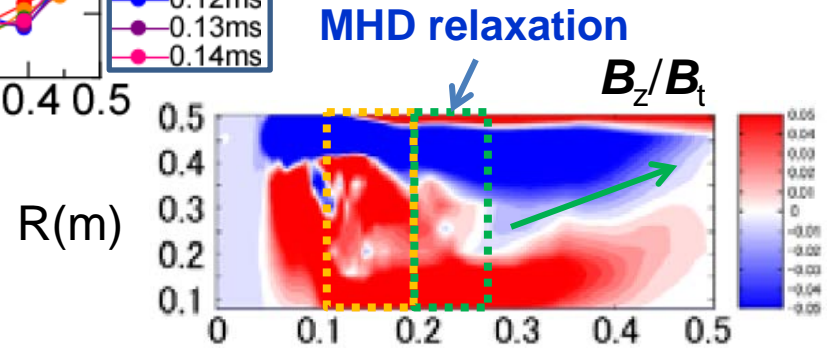
□ We found that the radial profile of the toroidal current density  $J_t(R)$  depends on the amount of injector flux.

□ The current density  $J_t$  is concentrated mainly in the open central column (OFC). Therefore,  $J_t(R)$  in the low bias flux shows a kink instability occurring at  $t = 0.14$  ms.

# Kink instability and MHD relaxation



**Low bias case**



A) During the injection phase → **Kink instability**

- ✓ Inner edge current increases and becomes unstable.
- ✓ The density decreases rapidly, but temperature does not so change.
- ✓ Correspondingly, the n=1 toroidal mode grows rapidly at t=0.12 ms

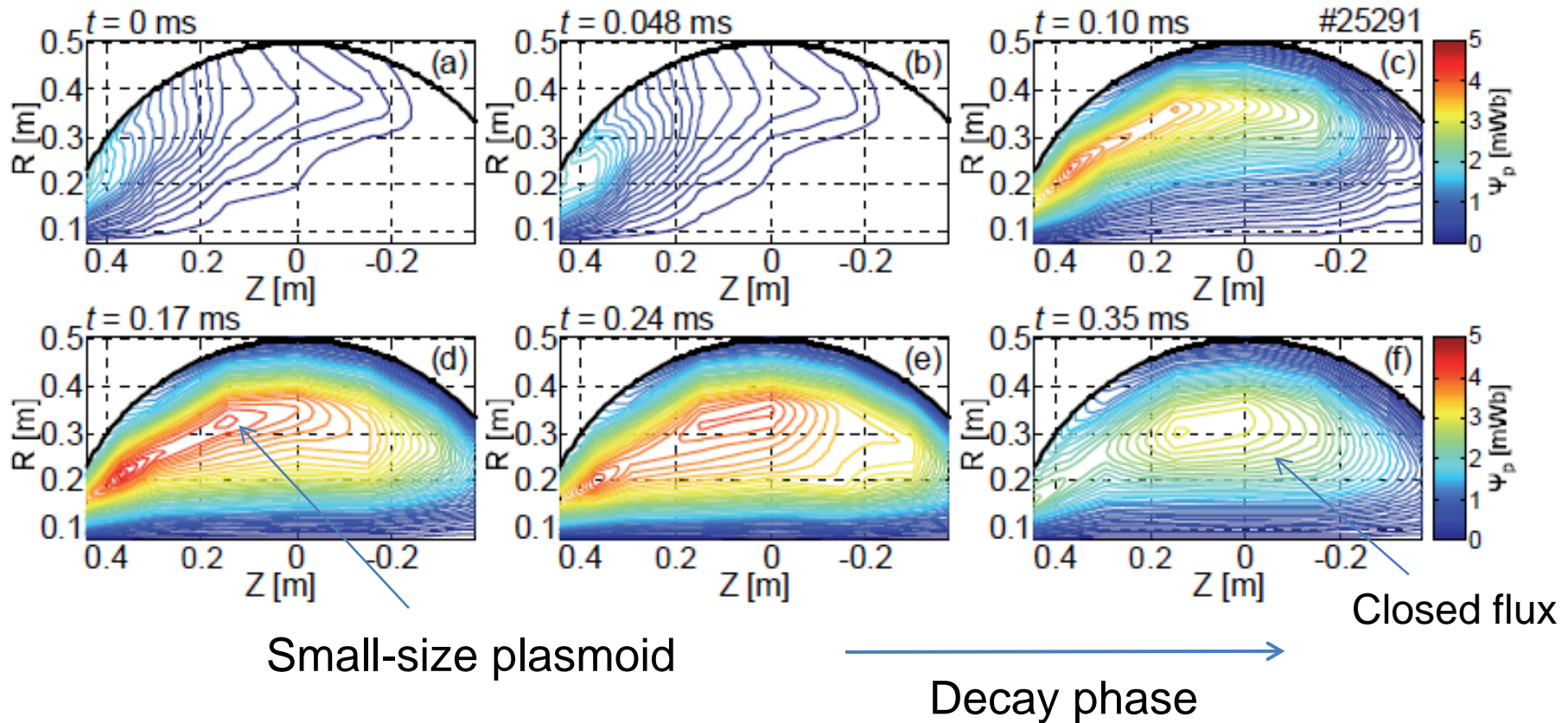
B) During the decay phase → **MHD relaxation**

- ✓ Density recovers shortly and the n=1 mode goes down
- ✓ The plasma relaxes into a stable state.



# Poloidal flux contour plots from 2D magnetic probe measurements

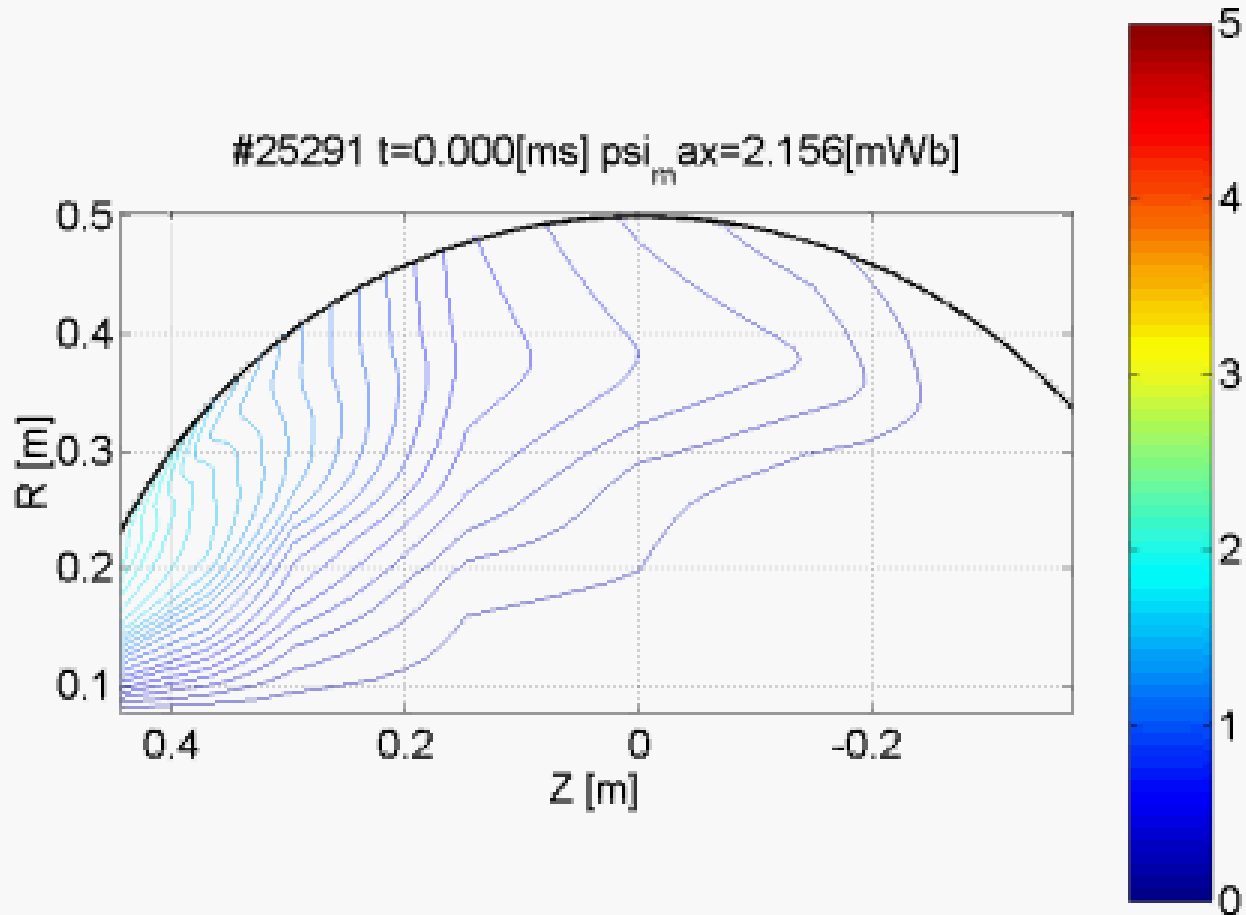
## High bias flux case



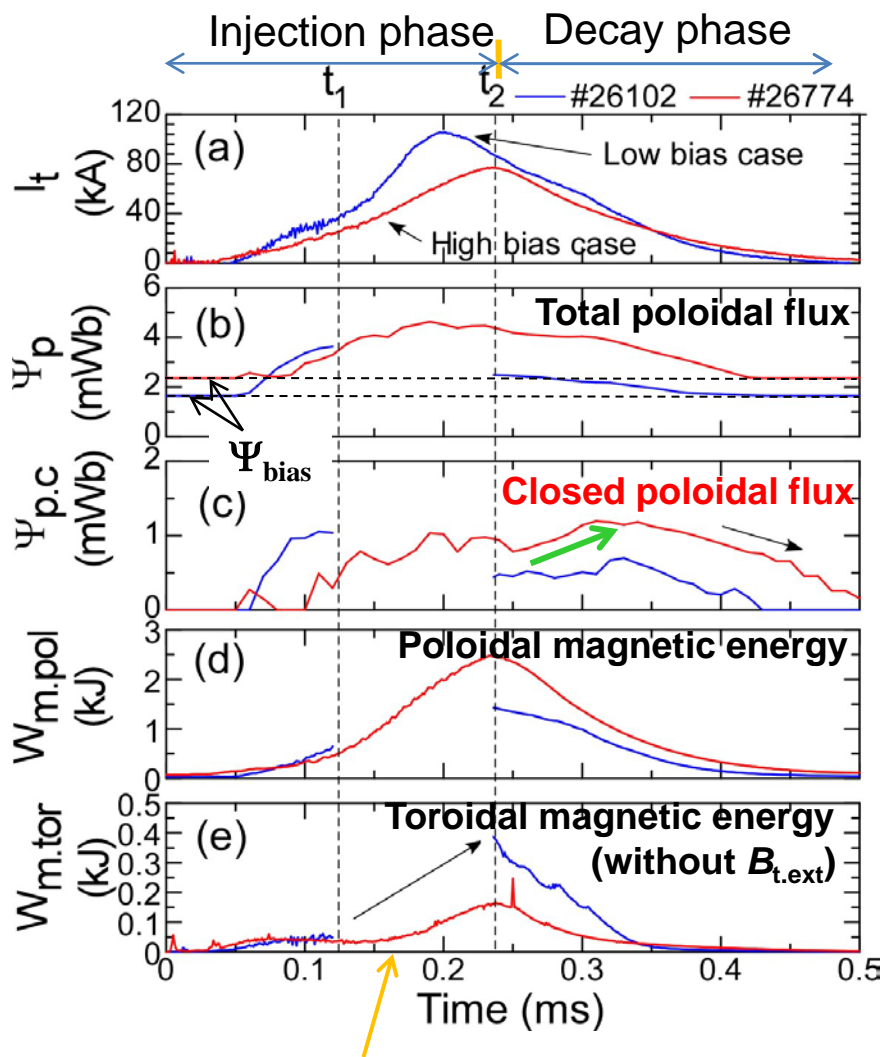
# Time evolution of poloidal flux contour plots

$\Psi_p$

High bias flux case



# Closed flux formation and flux conversion



Kink relaxation period in the low bias case

- The **stable closed configuration** can be obtained under the high bias flux case.
- The ratio of poloidal flux/the bias poloidal flux in the both cases is about 2. Flux amplification  $(\Psi_p/\Psi_{bias}) = 2$
- In the high bias case, the ratio of closed poloidal flux /the total poloidal flux increases from 25 % to 35 %.

$$\Psi_{p.closed}/\Psi_p = 25-35 \%$$

Poloidal magnetic field energy  $W_{m.pol}$



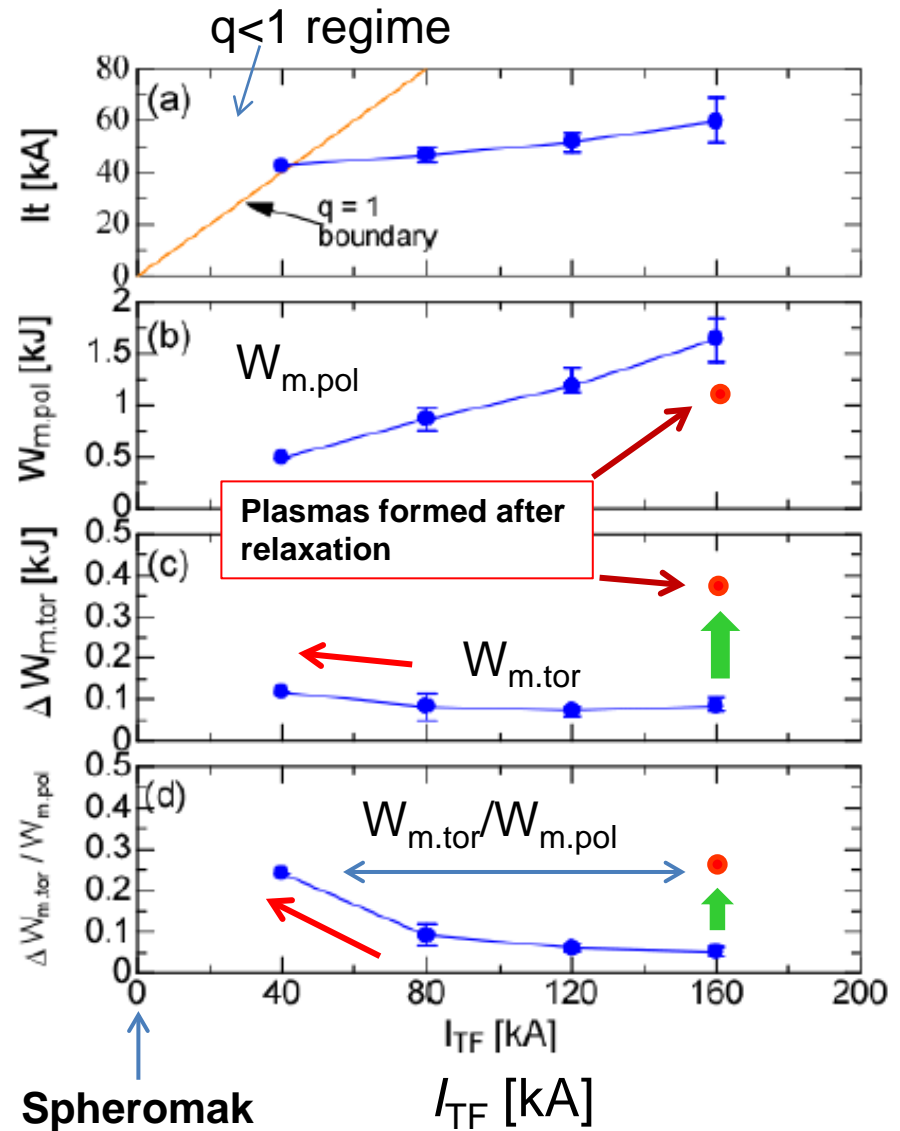
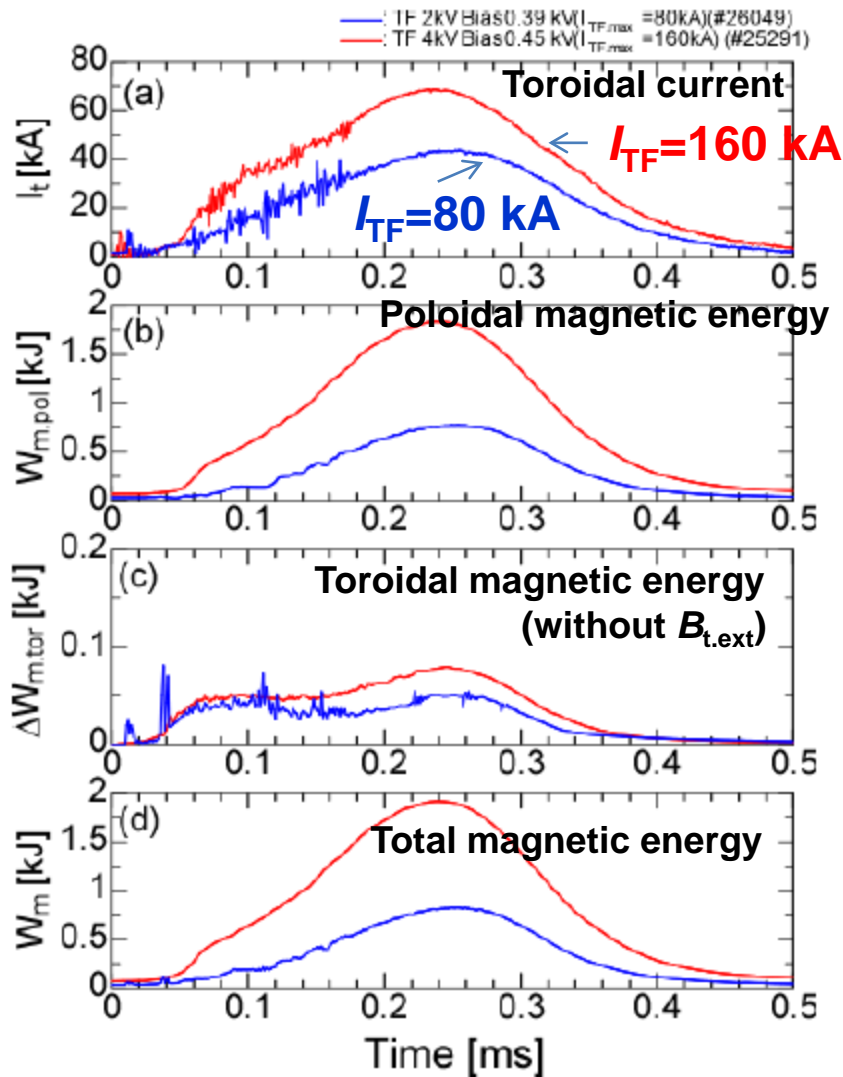
Kink relaxation

Toroidal magnetic field energy  $W_{m.tor}$

- ✓ The toroidal (injection) current in the OFC is converted to the poloidal current.

◆ This should be utilized for the steady-state CHI dynamo current drive.

# Dependence of TF coil current $I_{TF}$



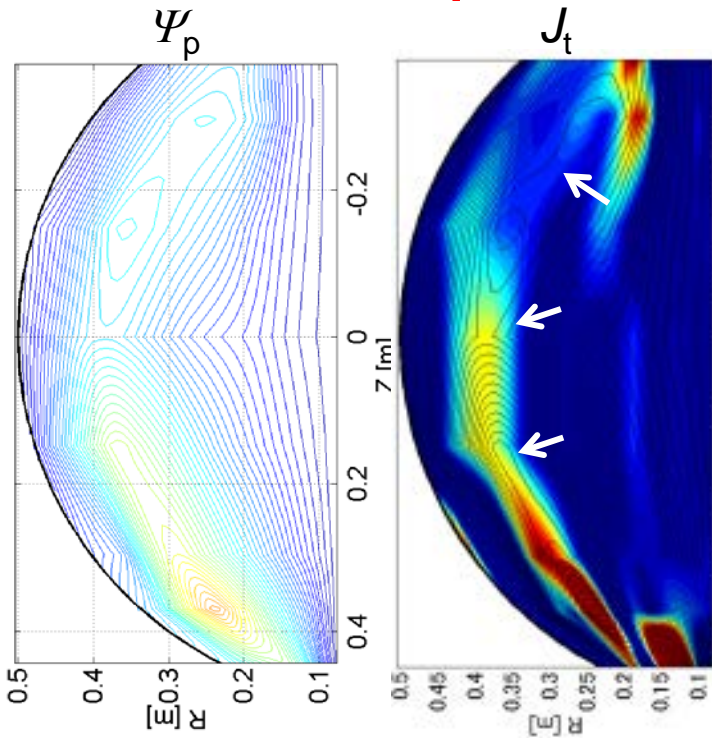
✓ Paramagnetic toroidal flux in the core region increases.



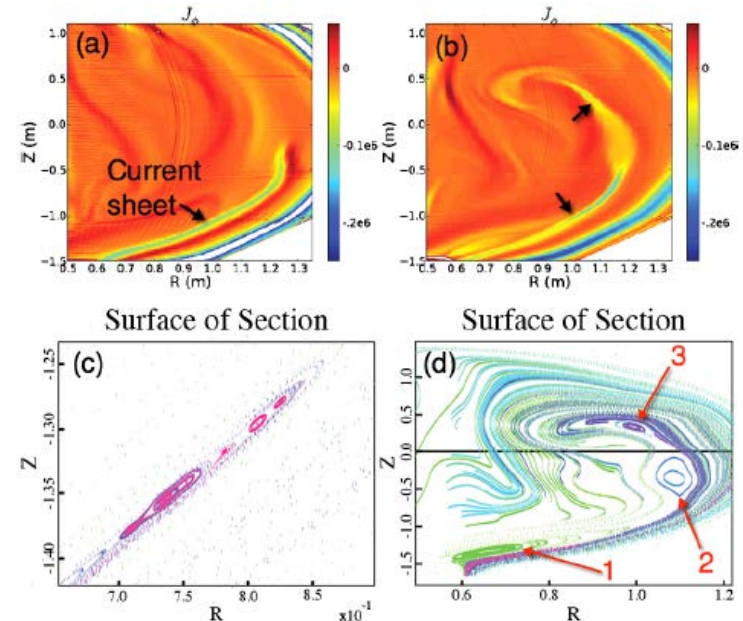
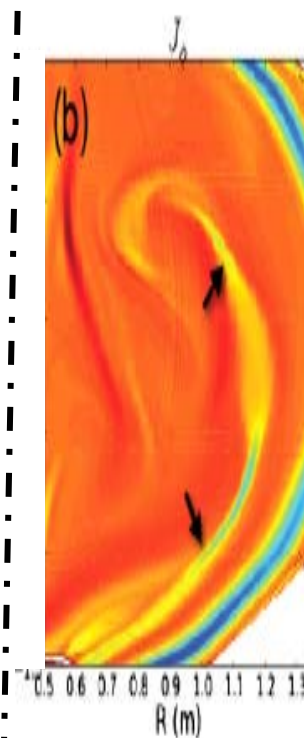
# Why is the current density localized and extended so long ?

- Ebrahimi's MHD simulation predicted that at high  $S > 3000$ , the elongated current sheet becomes tearing unstable, then the transition to plasmoid instability occurs during the injection phase, leading to the fast flux closure.
- In the HIST experiments, formation of the elongated current sheet and two or three plasmoids have been observed during T-CHI.

**HIST Exp.**



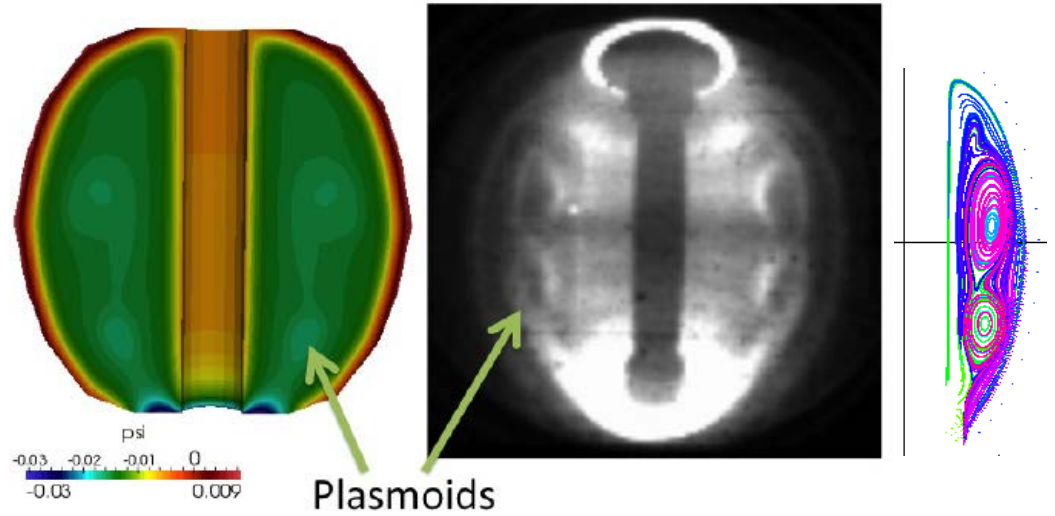
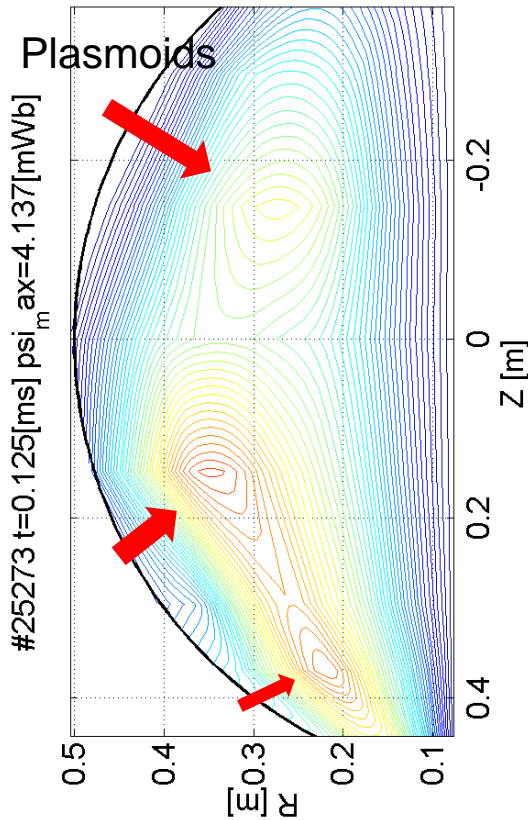
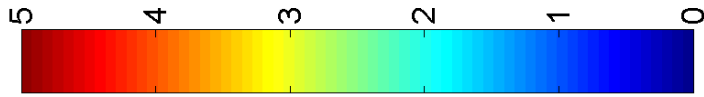
**MHD Simulation (NSTX)**



**F. Ebrahimi and R. Raman, PRL, 114, 205003 (2015).**

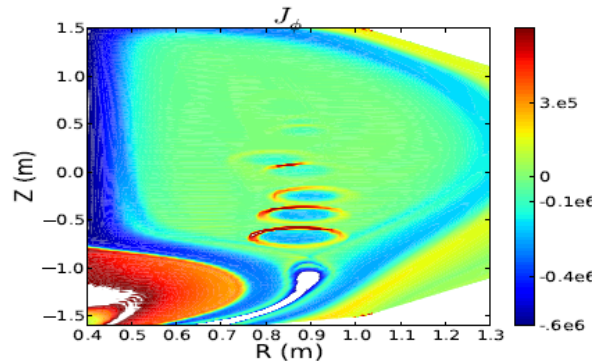
# Doublet type equilibrium configuration formed by two discrete plasmoids

HIST exp.



$T_e \sim 5$  eV,  $n \sim 4-5 \times 10^{18} \text{ m}^{-3}$ ,  $B_{\text{rec}} \sim 0.01$  T,  $L \sim 1$  m,  $S = 2000-4000$ .

✓ Camera image during CHI discharge in NSTX shows the formation of two discrete plasmoids.



✓ Plasmoid instability with continued injection of plasmoid is observed during the injection phase. ( $S \sim 29000$ )

# Spontaneous magnetic reconnection driven by plasmoid instability

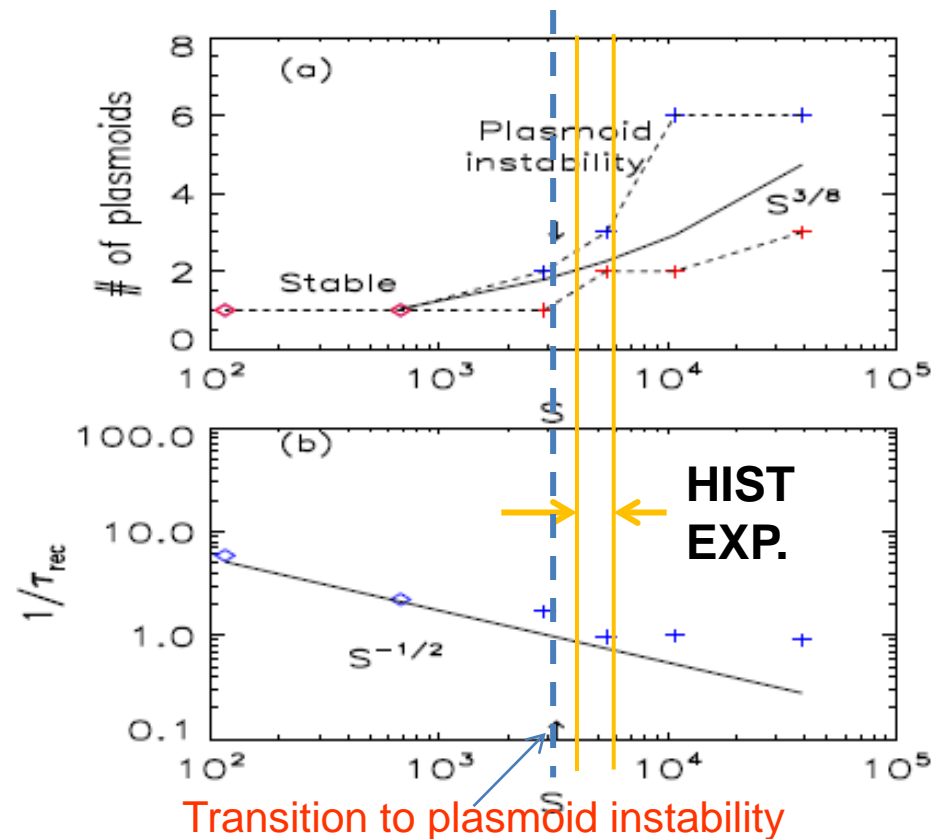
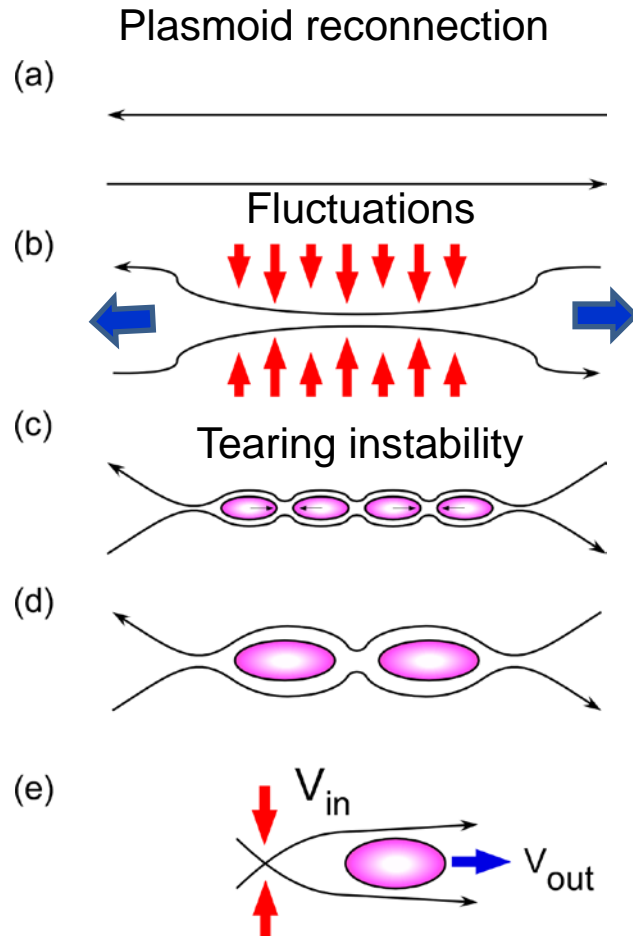


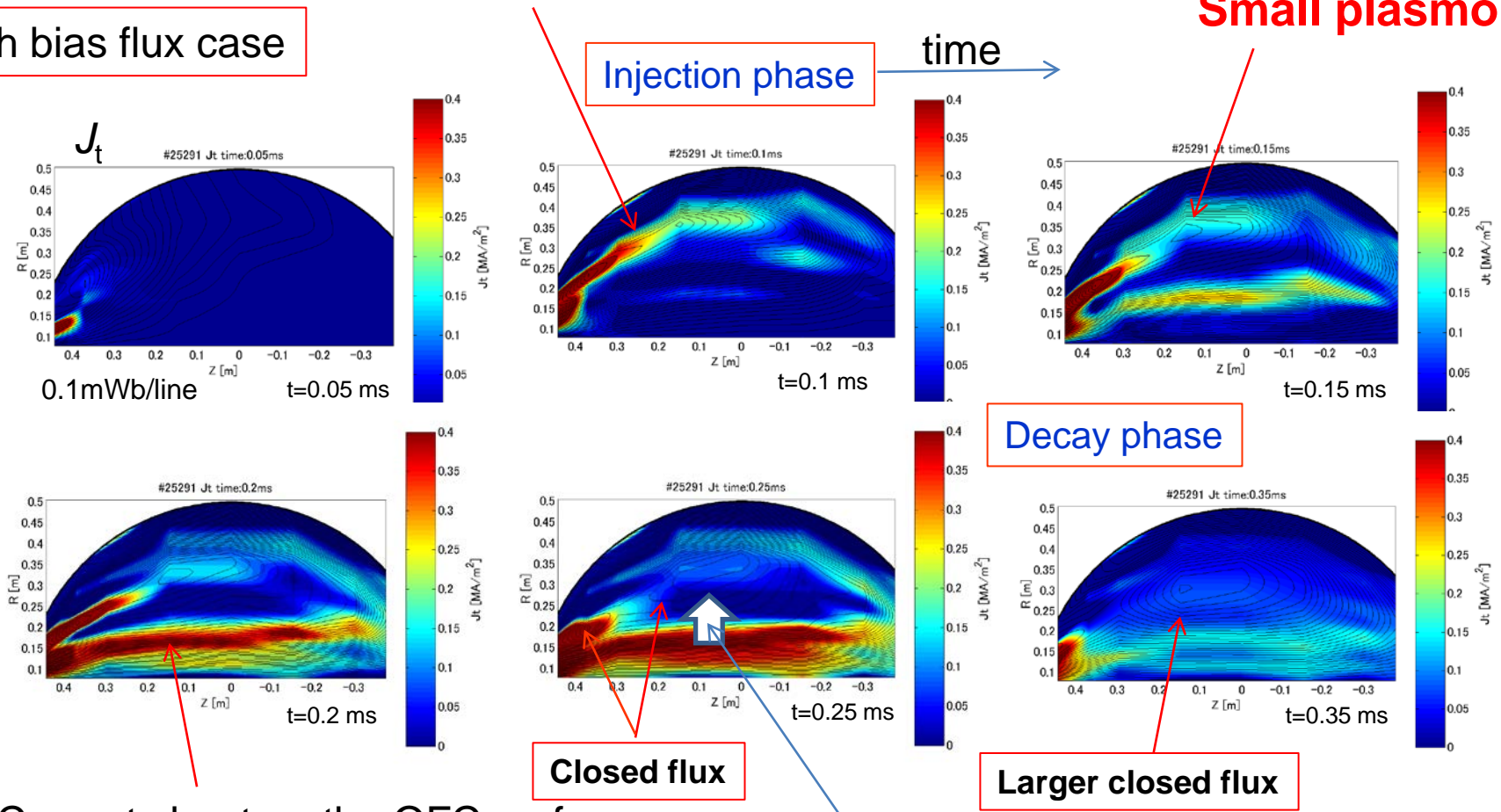
FIG. 3 (color online). (a) Number of plasmoids vs  $S$ . Blue: small sized transient plasmoids during the early phase of discharge (Fig. 2). Red: large scale and persistent plasmoids during the later phase of discharge. The solid line is the linear theoretical  $S$  scaling. (b) The reconnection rate,  $1/\tau_{rec}$  vs  $S$ . The transition to plasmoid instability is shown at  $S \sim 3000$ . The solid line is the S-P scaling.

- ✓ Plasmoid formation allows the fast reconnection rate, nearly independent of Lundquist number  $S$  at the high  $S$  regime.

# Current sheet elongation, magnetic reconnection and small-size plasmoids

## Current sheet elongation

High bias flux case



Current sheet on the OFC surface



Diamagnetic  $B_t$  in the OFC

✓ Inner edge current diffuses into the small plasmoid, leading to a large-scale closed flux.



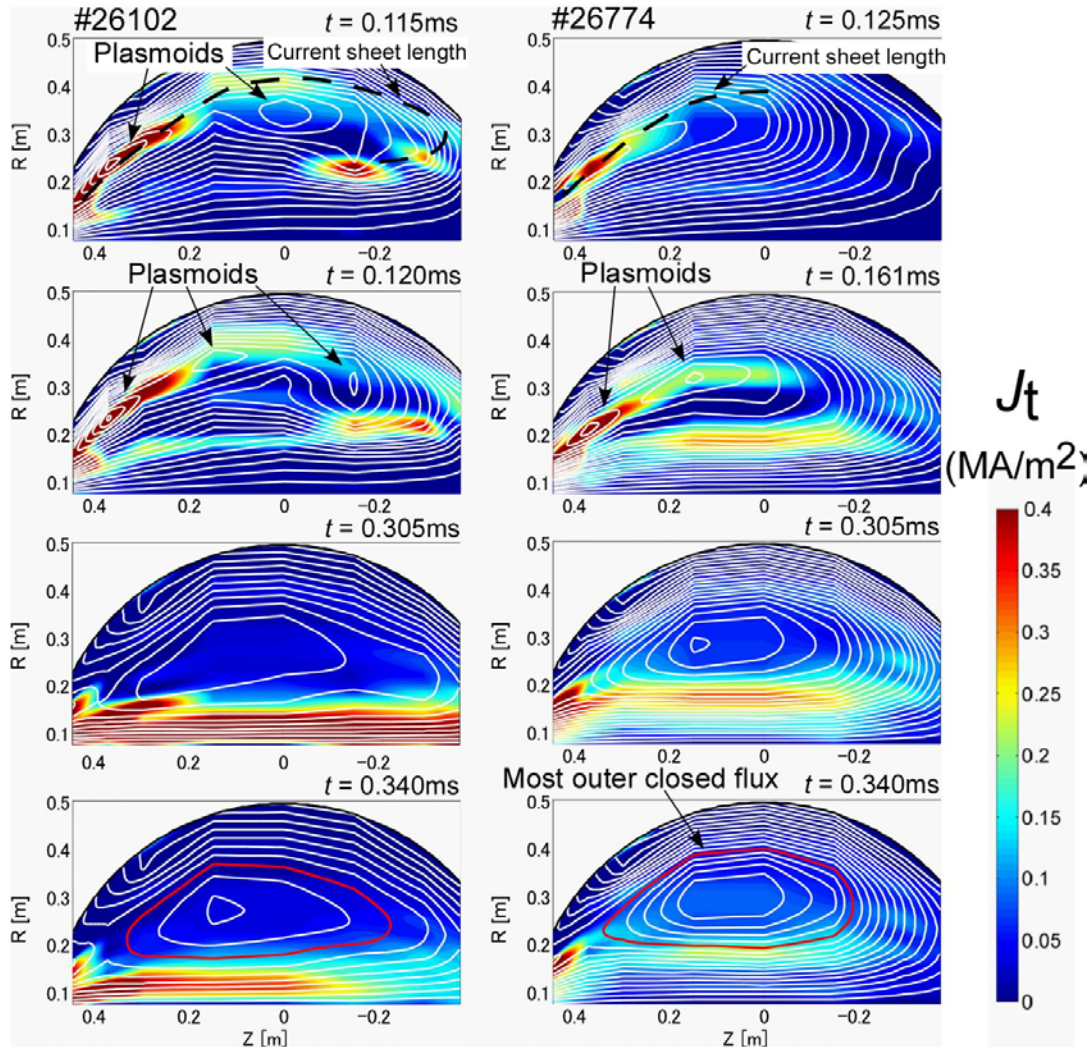
# Comparison of $J_t$ and $\Psi_p$ contours between in both cases

Low bias flux case

High bias flux case

(a) Low bias case

(b) High bias case



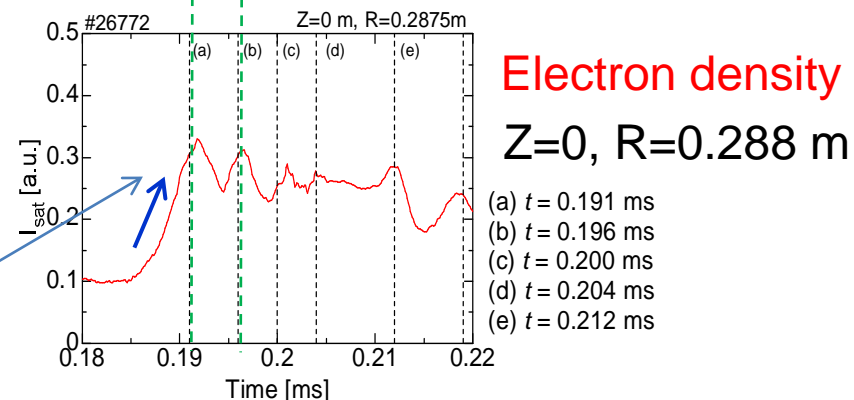
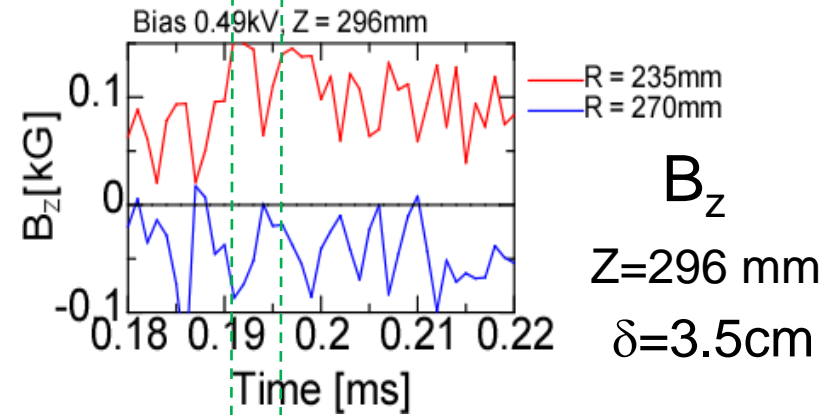
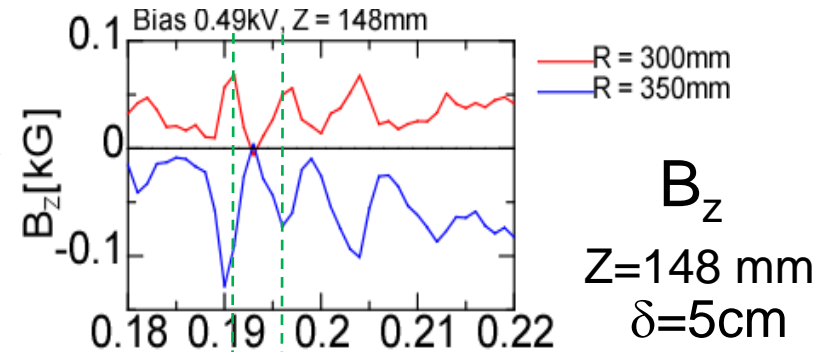
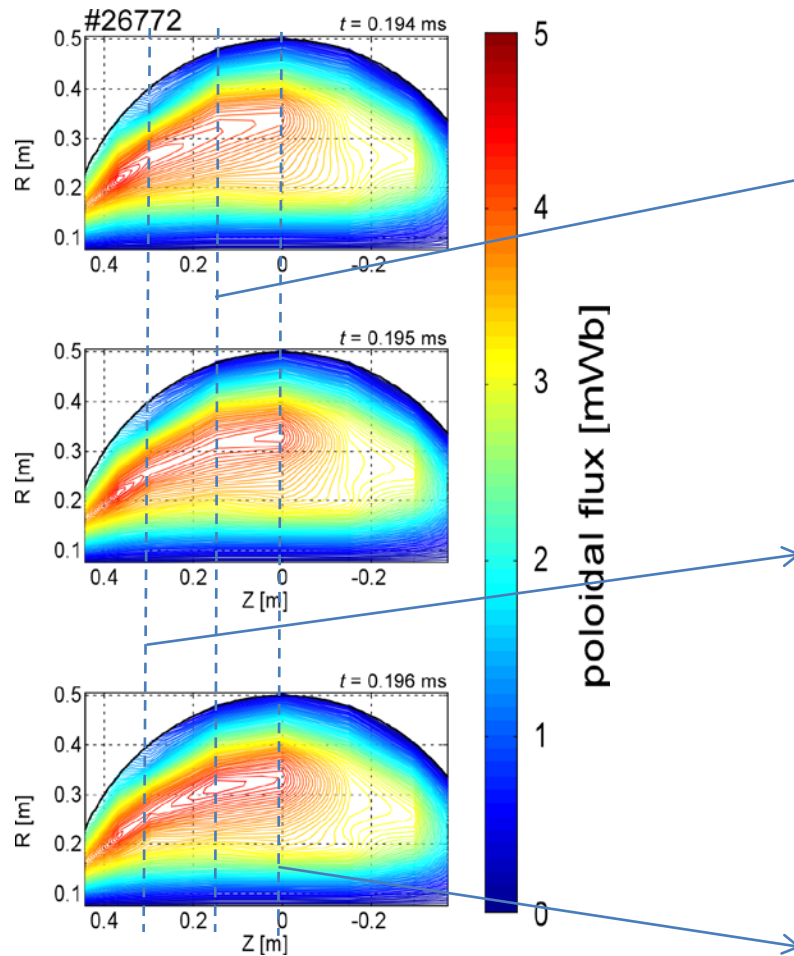
□ The current sheet is elongated during the current rise phase in both cases.

□ The reconnection occurs as the **current sheet length  $L$**  increases and so the **current sheet width  $\delta$**  becomes enough thin.

□ The **plasmoids** are generated due to the reconnection in the **elongated current sheet**.

□ The **number of plasmoids** in the low bias case with a longer  $L$  is larger by one or two than that in the high bias case.

# Behavior of plasmoids and reconnected field lines

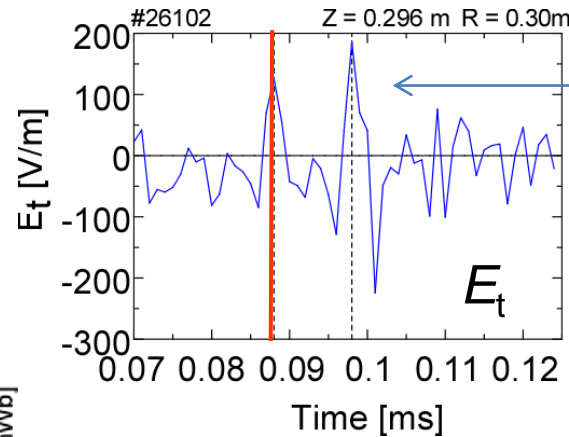
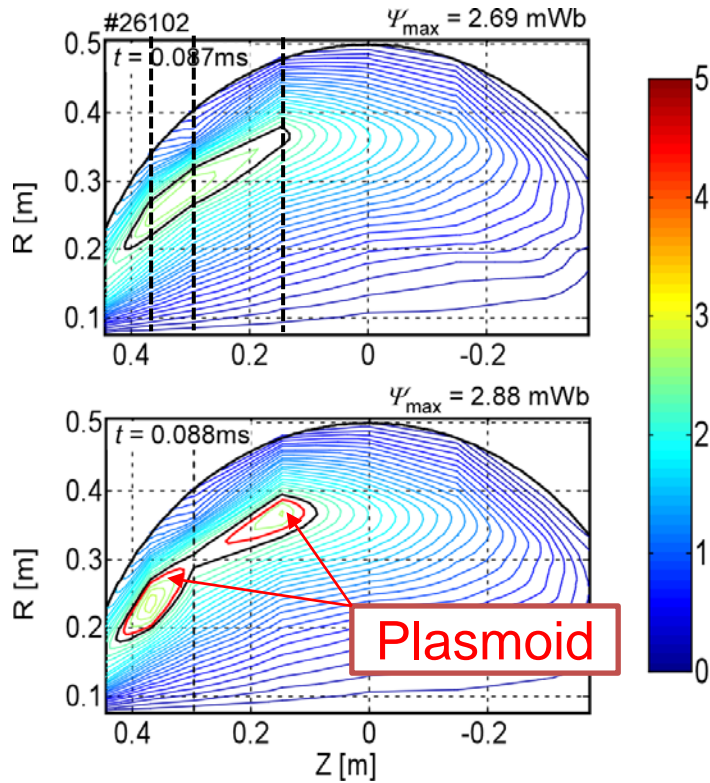


## Plasmoid dynamics

- ✓ Oscillation of the oppositely directed poloidal fields  $B_z$ , which together to be reconnected, synchronizes with that of the electron density.

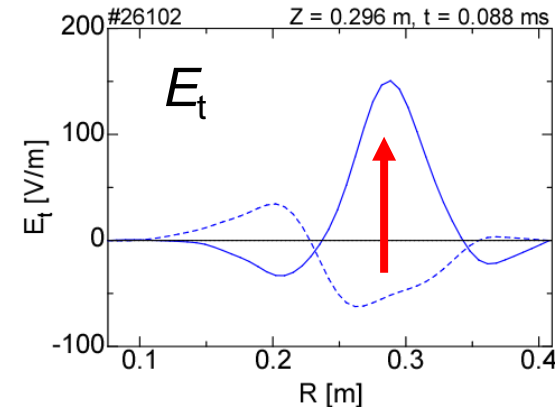
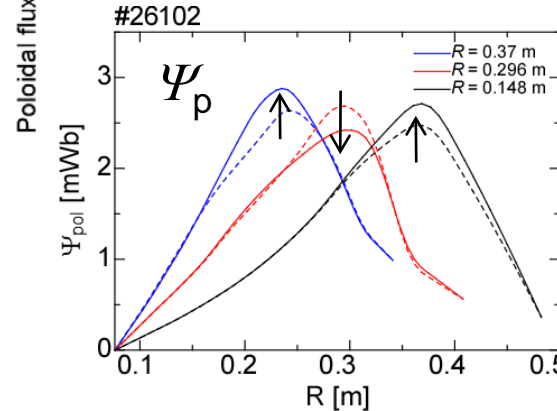
$I_{\text{sat}}$  (a.u.)

# Time evolution of reconnection $E_t$ and $\Psi_p$ during plasmoid's separation



The reconnection  $E_t$  is enhanced.

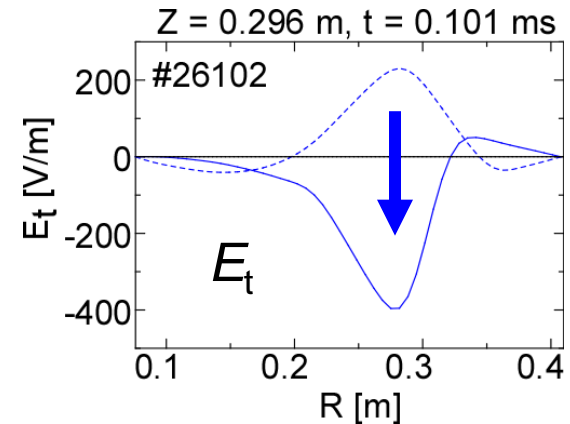
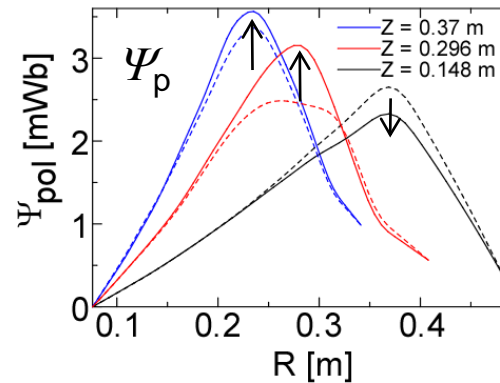
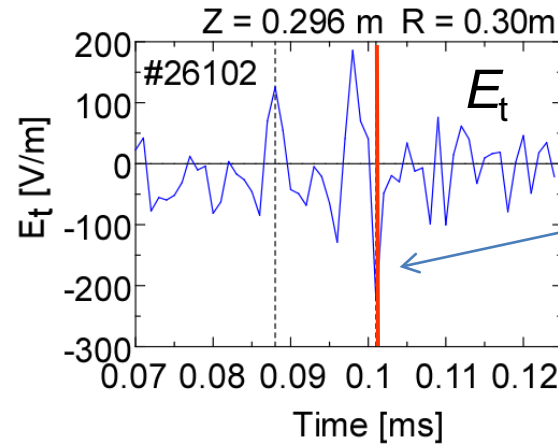
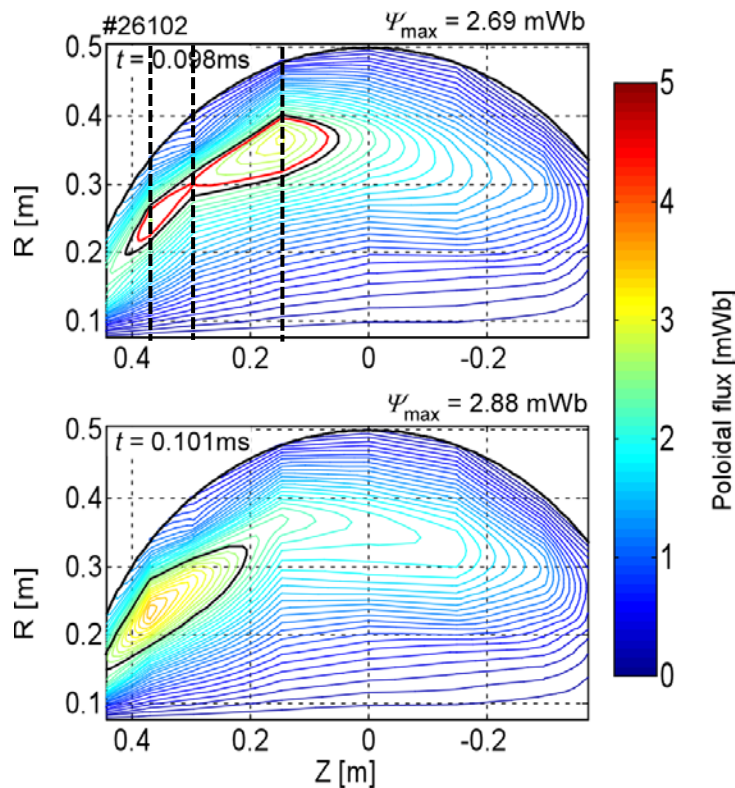
$$E_t = -\frac{1}{2\pi R} \frac{d\psi_p}{dt}$$



✓ One plasmoid separates to two smaller size plasmoids.

✓ At  $t = 0.088$  ms, one plasmoid separates to two small plasmoids and the poloidal flux decreases. At this time,  $E_t$  shows a positive spike signal and the polarity of the profile is reversed during the merging.

# Time evolution of reconnection $E_t$ and $\Psi_p$ during plasmoid's merging



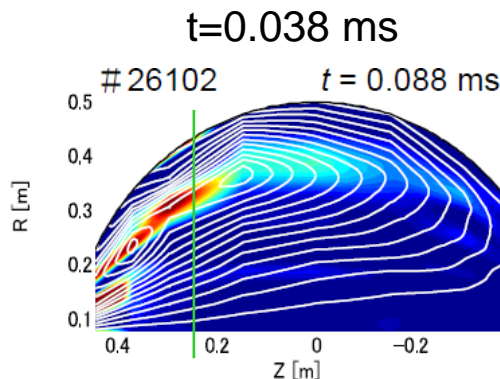
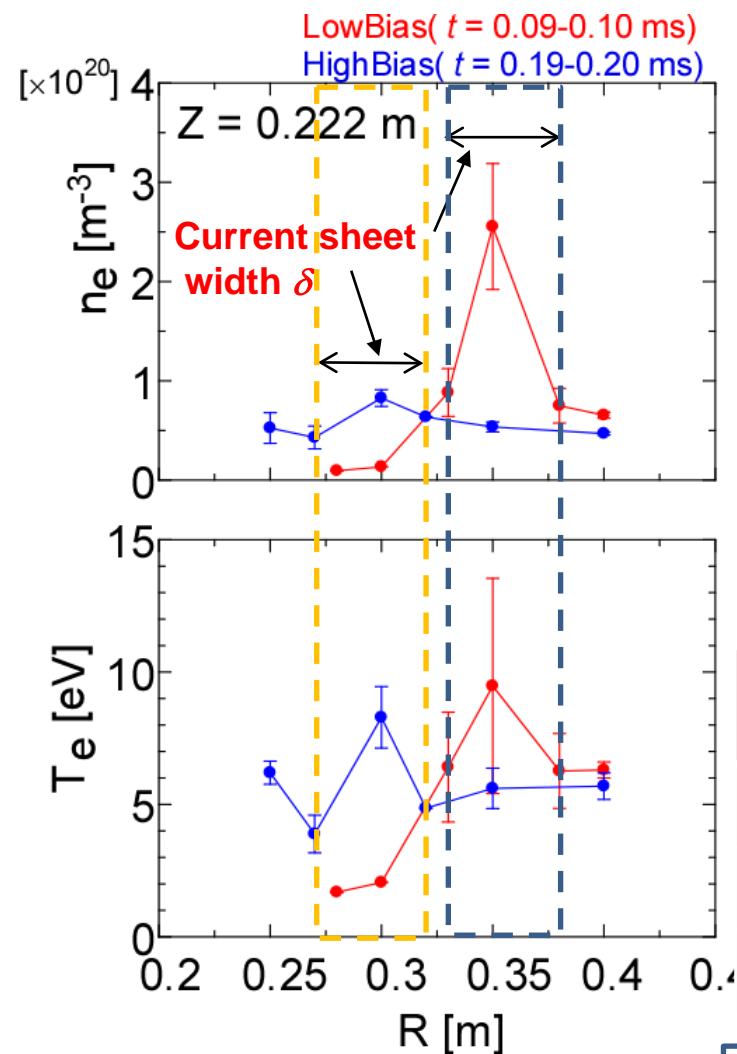
✓ Two plasmoids merge to one plasmoid.

✓ At  $t = 0.101$  ms, two plasmoids merge and plasmoid's poloidal flux increases. At this time,  $E_t$  shows a negative spike signal and the polarity of the profile is reversed during the merging.

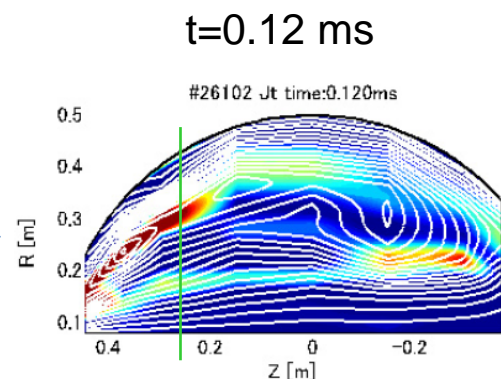


# Estimation of Lundquist number $S$ in the current sheet

## - Correlation with its width $\delta$ and length $L$ , and the number of plasmoids -



The number of plasmoids:  $N=2$   
 $L \sim 0.59$  m,  $n_e \sim 2.5 \times 10^{20} m^{-3}$



The number of plasmoids:  $N=3$   
 $L \sim 1.1$  m,  $n_e \sim 0.9 \times 10^{20} m^{-3}$

$$B = B_p + B_t$$

$$S = \frac{\mu_0 \cdot v_A \cdot L}{\eta}$$

Bias flux	B [T]	L [m]	$\delta$ [cm]	$T_e$ [eV]	$n_e$ [ $m^{-3}$ ]	$S$	The number of Plasmoid
Low	0.08	0.59	$\sim 5$	10	$2.5 \times 10^{20}$	5000	2
Low	0.08	1.1	$\sim 5$	6	$0.9 \times 10^{20}$	7200	3
High	0.08	0.56	$\sim 5$	8	$0.8 \times 10^{20}$	5800	2

S-P scaling:  $\delta_{sp}/L \sim S^{-1/2}$

Ion skin depth

$d_i = 2 \sim 3$  cm

$\delta \geq d_i$

No two-fluid effect ?

# Summary

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- **Flux closure and magnetic reconnection in the T-CHI generated plasmas have been investigated on HIST.**
- It has been found that the toroidal current sheet is elongated in the presence of the high TF.
- Two or three plasmoids in the elongated current sheet have been directly detected by internal 2D magnetic probe measurements. Also, we have examined the dynamics of plasmoids (merging and separation).
- The each parameter ( $\delta$ ,  $L$ ,  $S$ ,  $T_e$ ,  $n_e$ ) of the elongated current sheet has been measured. The  $S$  number is estimated to be in the range of 5000-7000. According to the MHD simulation, the onset of the plasmoid instability may be possible around the measured  $S$  number.
- Plasmoid instability could be one possible mechanism for generation of larger-size closed flux surfaces.
- **Self-organization of T-CHI generated plasmas have been observed.**
- The relaxation process exhibited the current conversion from toroidal to poloidal direction.