



TRI ALPHA ENERGY
THE POWER OF INGENUITY

PIC Simulation of Thermal Distribution Driven Non-Maxwellian by Neutral-Beam Injection in a High Beta Plasma

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and the Entire TAE Team**

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92688, USA***

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Outline

■ Experimental Motivation:

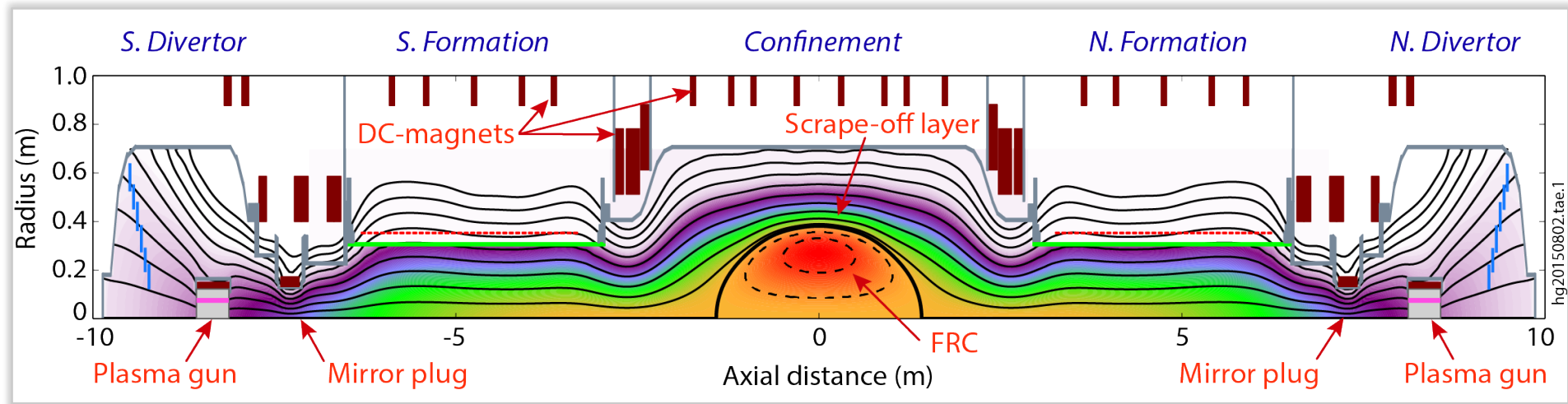
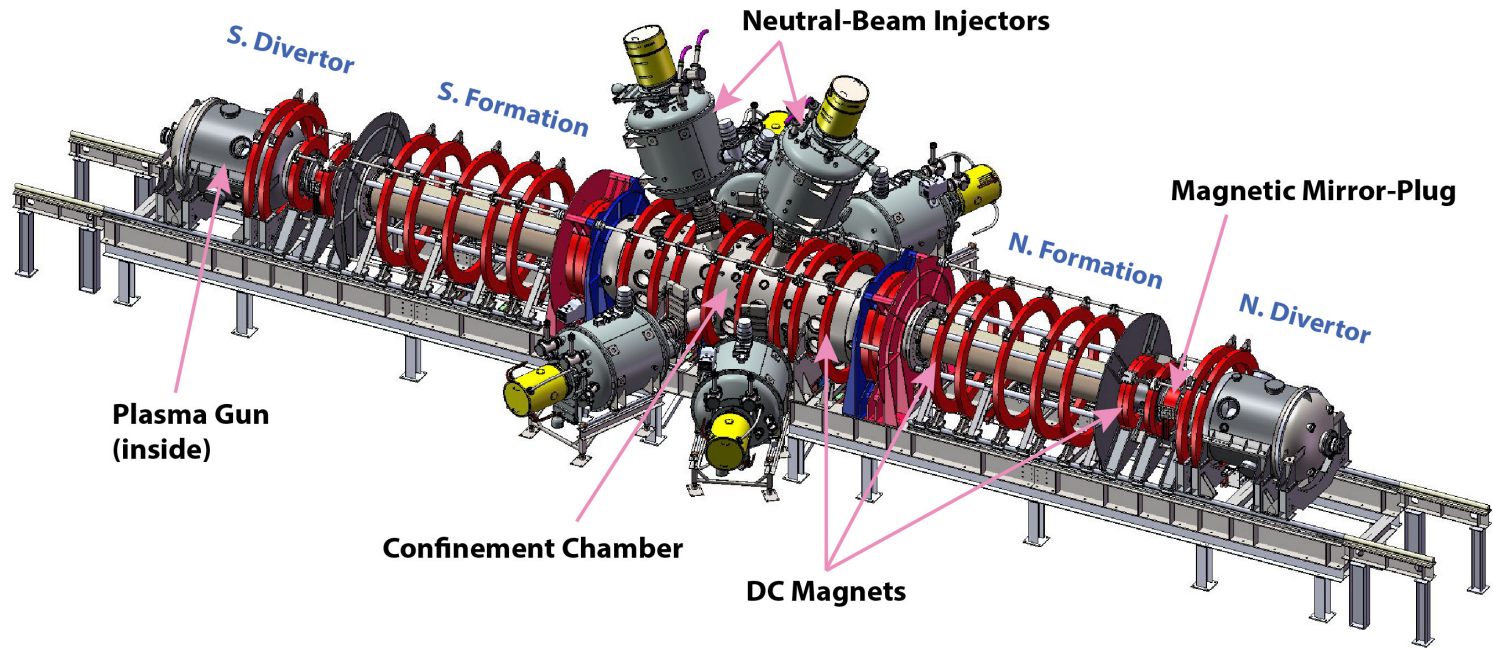
- Enhanced neutron yield

■ Theory and computation:

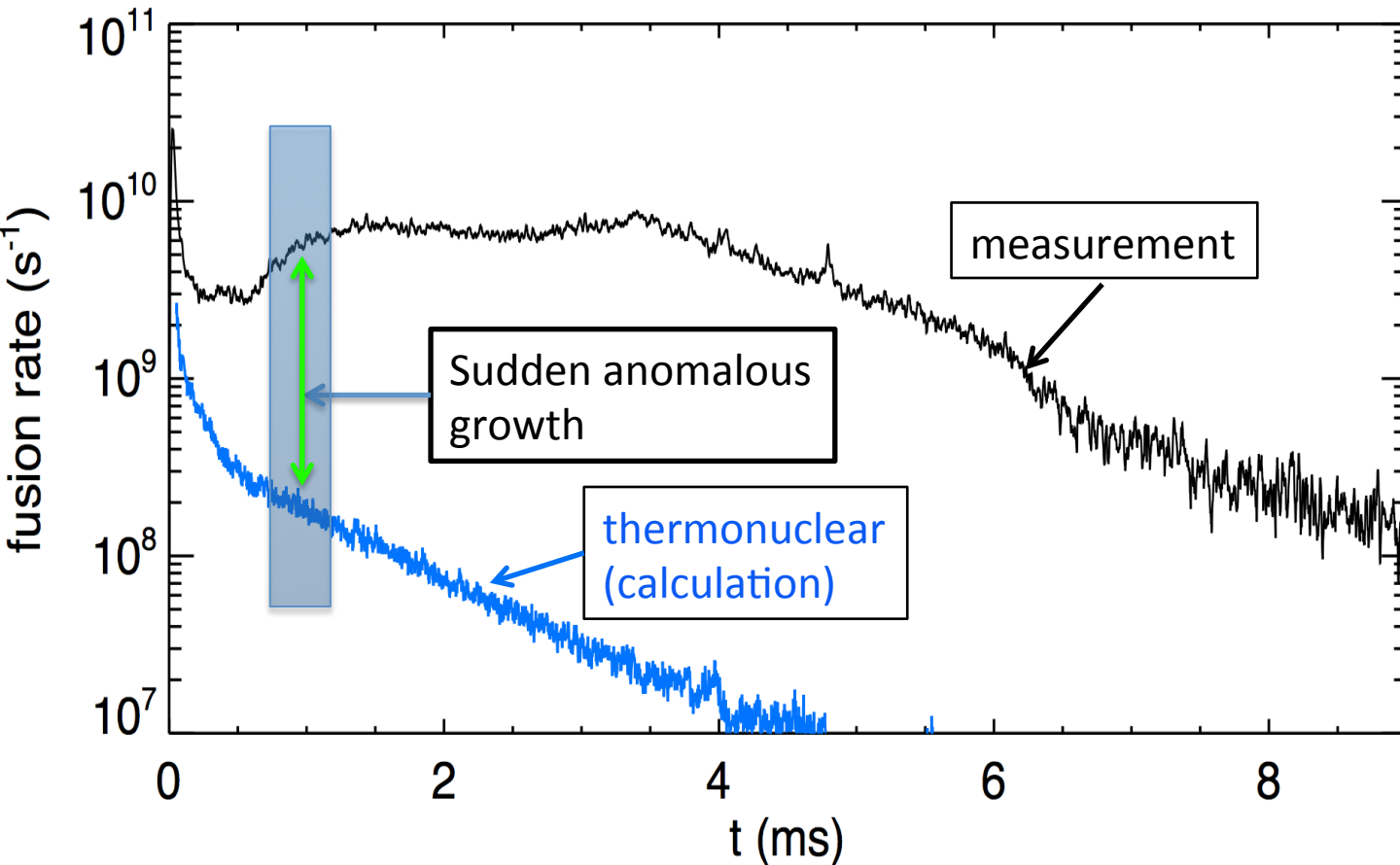
- Hypothesis
- Simulation problem set up
- ES benchmark
- Hypothesis verification
- Klimontovich accounting – neutron generation
- Comparison between high and low(er) beta – Mode discussion
- Comparison to Neutral Particle Analyzer data

■ Summary

C-2U Machine Configuration



Motivation - search for beam-driven micro-instabilities suggested by enhanced reactivity



- Hydrogen beam injection enhances D-D fusion rate.
- ~ 10 times increase in neutron

Courtesy: Richard Magee

Outline

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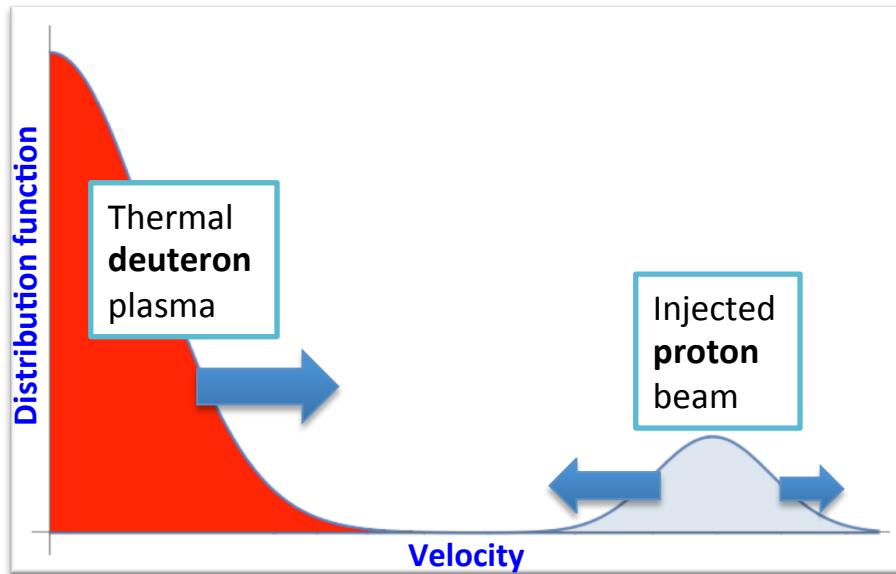
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■ Summary

Hypothesis



Presence of proton beam is source of “free energy”



Transfer energy from beam to deuteron



Energetic super-thermal tail in the thermal **deuteron** plasma



Accelerated deuterons collide with thermal (or super-thermal) deuterons



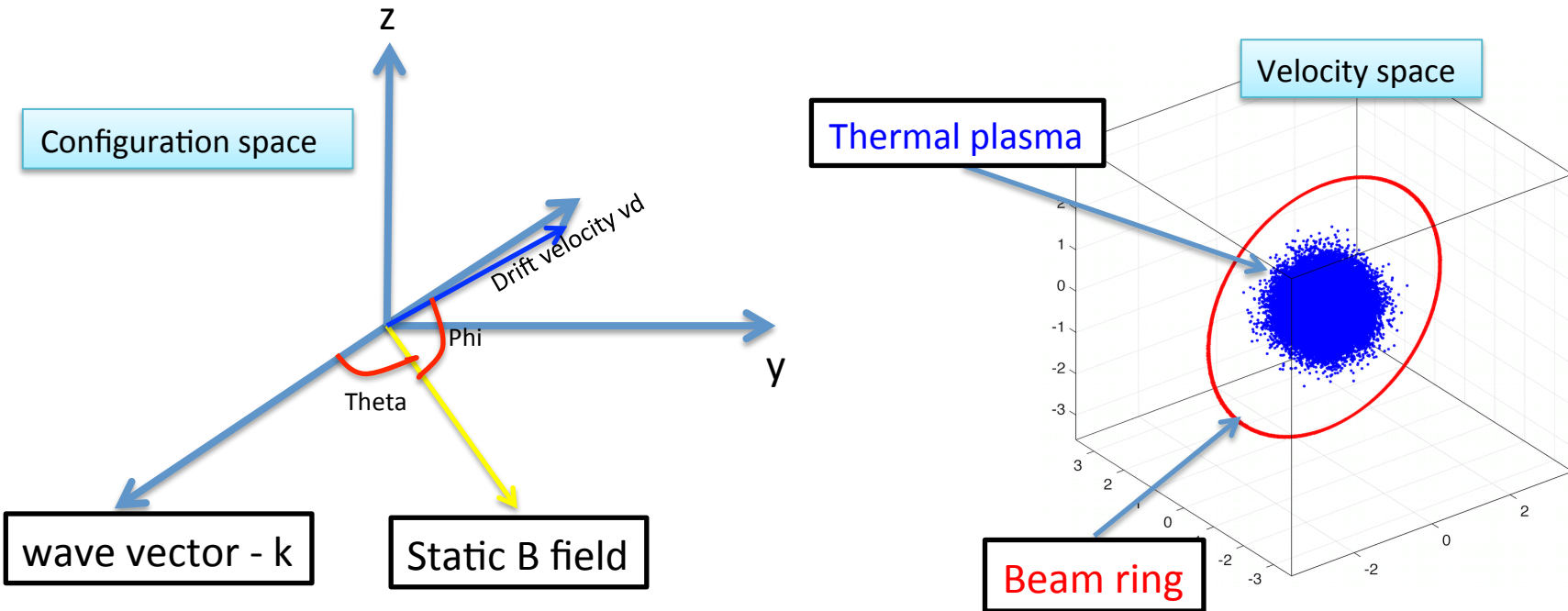
Enhanced increase in neutron production

Testing the hypothesis



Use a 1D PIC code (EPOCH)

Numerical Problem Set Up



- Fully electromagnetic **1D3V** PIC code: Maxwell eqns + Lorentz force
- **Slab model with periodic BC**
- **Initial value problem (introduce beam at t=0)**
- **Thermal plasma** is Maxwellian, optional: anisotropy, drift
- Beam is set up as a Maxwellian distribution, **ring** or slowing down distribution – sampling problem. Anisotropy optional.
- Beam set up with velocity parallel and perp to B field
- Particle interaction via collective effects **ONLY**

Electrostatic Case – Perpendicular Propagation I

Dispersion relation

$$\frac{\omega_{pi}^2}{\lambda_i} \sum_{n=1}^{\infty} \exp[-\lambda_i] I_n(\lambda_i) \frac{2n^2}{\omega^2 - n^2 \Omega_i^2} +$$

Thermal ion

$$\frac{\omega_{pb}^2}{k v_b} \sum_{n=1}^{\infty} \frac{2n^2 \omega_{cb}}{\omega^2 - n^2 \omega_{cb}^2} J_n(\nu) \left[J_{n-1}(\nu) - J_{n+1}(\nu) \right] =$$

Beam

$$\lambda_i = 0.5(k\rho_i)^2$$

$$\nu = \frac{k v_b}{\omega_{cb}}$$

$$1 + \left(\frac{\omega_{pe}}{\Omega_{ce}} \right)^2 \left(1 + \frac{\omega_{pe}^2}{c^2 k^2} \right)$$

Thermal electron

- Beta = 1e-4
- Let $\omega = \omega + i\gamma$

Electrostatic Case – Perpendicular Propagation II

Dispersion relation

$$\frac{\omega_{pi}^2}{\lambda_i} \sum_{n=1}^{\infty} \exp[-\lambda_i] I_n(\lambda_i) \frac{2n^2}{\omega^2 - n^2 \Omega_i^2} +$$

Thermal ion

$$\frac{\omega_{pb}^2}{k v_b} \sum_{n=1}^{\infty} \frac{2n^2 \omega_{cb}}{\omega^2 - n^2 \omega_{cb}^2} J_n(\nu) \left[J_{n-1}(\nu) - J_{n+1}(\nu) \right] =$$

Beam

$$\lambda_i = 0.5(k\rho_i)^2$$

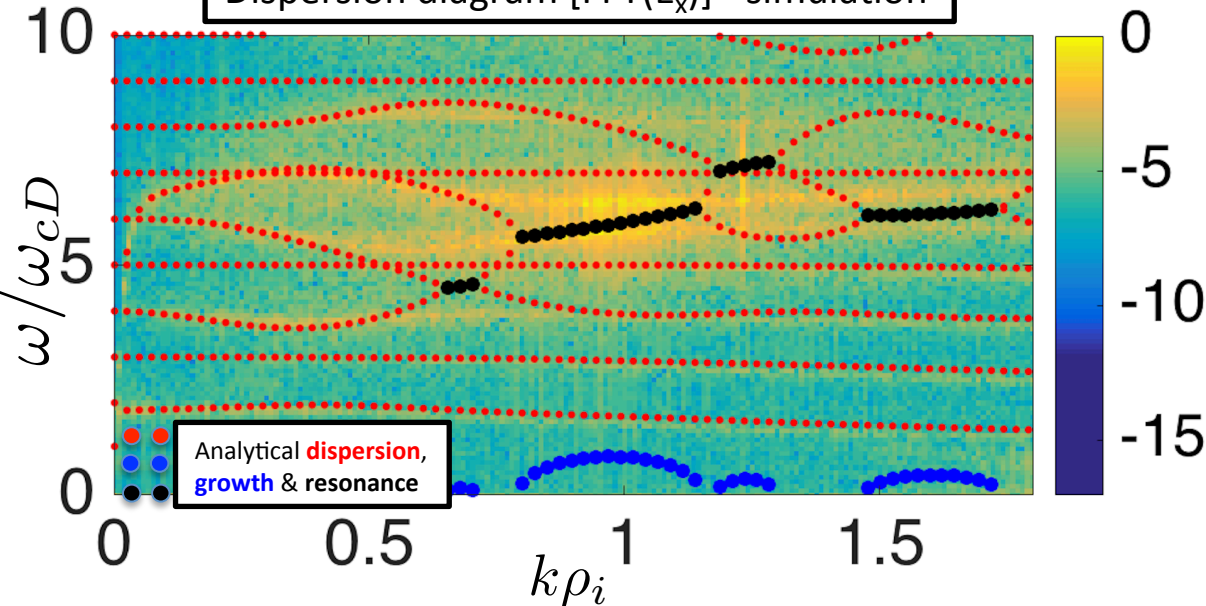
$$\nu = \frac{k v_b}{\omega_{cb}}$$

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Thermal electron

- Beta = 1e-4
- Let $\omega = \omega + i\gamma$
- Coupling of plasma lower hybrid wave and n=6 ion Bernstein harmonic
- Growth rate = 0.4 ω_{ci}

Dispersion diagram [FFT(E_x)] - simulation



Electrostatic Case – Perpendicular Propagation III

Dispersion relation

$$\frac{\omega_{pi}^2}{\lambda_i} \sum_{n=1}^{\infty} \exp[-\lambda_i] I_n(\lambda_i) \frac{2n^2}{\omega^2 - n^2 \Omega_i^2} +$$

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Beam

$$1 + \left(\frac{\omega_{pe}}{\Omega_{ce}} \right)^2 \left(1 + \frac{\omega_{pe}^2}{c^2 k^2} \right)$$

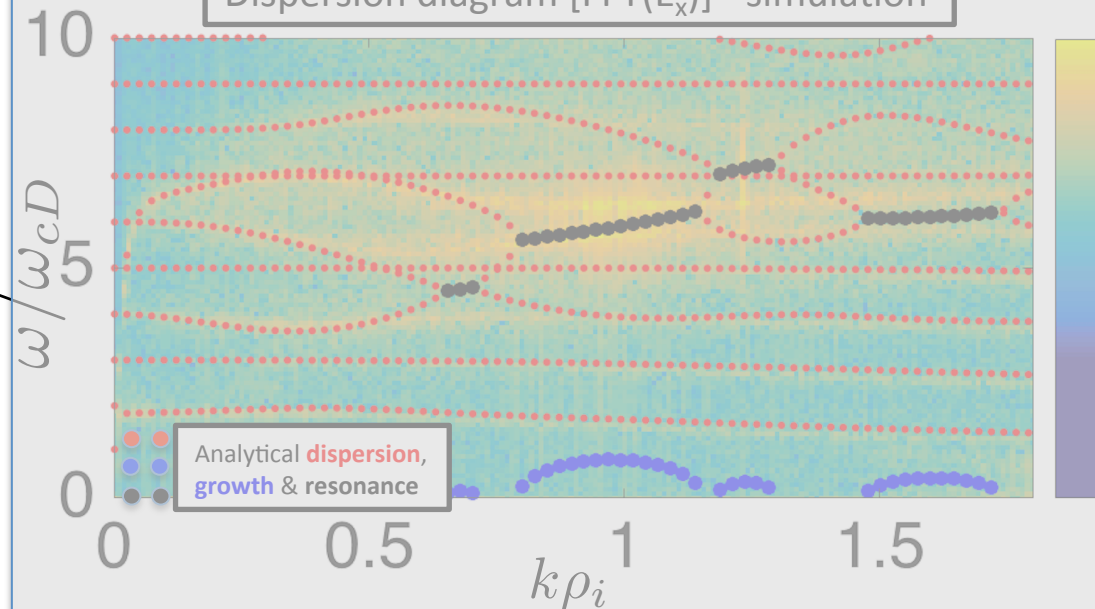
Thermal electron

$$\lambda_i = 0.5(k\rho_i)^2$$

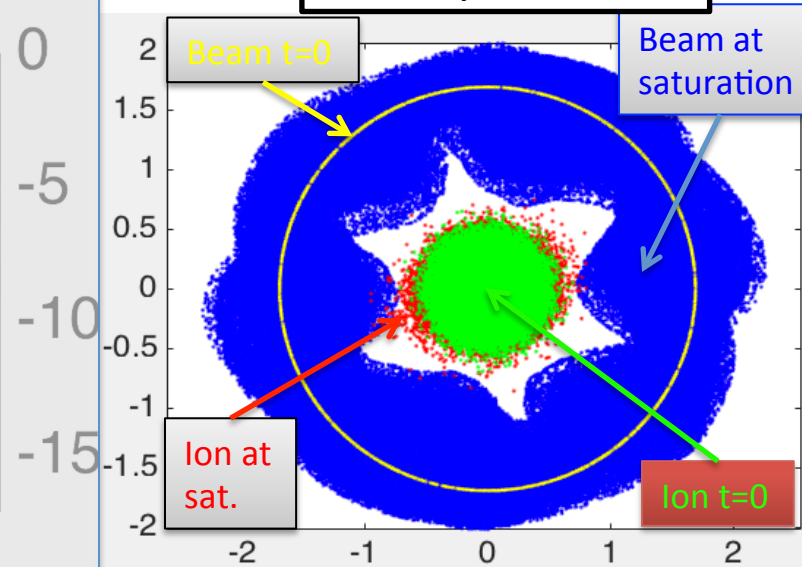
$$\nu = \frac{k v_b}{\omega_{cb}}$$

- Beta = 1e-4
- Let $\omega = \omega + i\gamma$
- Coupling of plasma lower hybrid wave and n=6 ion Bernstein harmonic
- Growth rate = 0.4 ω_{ci}
- **Good agreement between analytics and simulation**

Dispersion diagram [FFT(E_x)] - simulation



Velocity [x10⁶ m/s]



Electrostatic Case – Perpendicular Propagation

Dispersion relation

$$\frac{\omega_{pi}^2}{\lambda_i} \sum_{n=1}^{\infty} \exp[-\lambda_i] I_n(\lambda_i) \frac{2n^2}{\omega^2 - n^2 \Omega_i^2} + \frac{\omega_{pb}^2}{k v_b} \sum_{n=1}^{\infty} \frac{2n^2 \omega_{cb}}{\omega^2 - n^2 \omega_{cb}^2} J_n(\nu) \left[J_{n-1}(\nu) - J_{n+1}(\nu) \right] = 1 + \left(\frac{\omega_{pe}}{\Omega_{ce}} \right)^2 \left(1 + \frac{\omega_{pe}^2}{c^2 k^2} \right)$$

Thermal ion

Beam

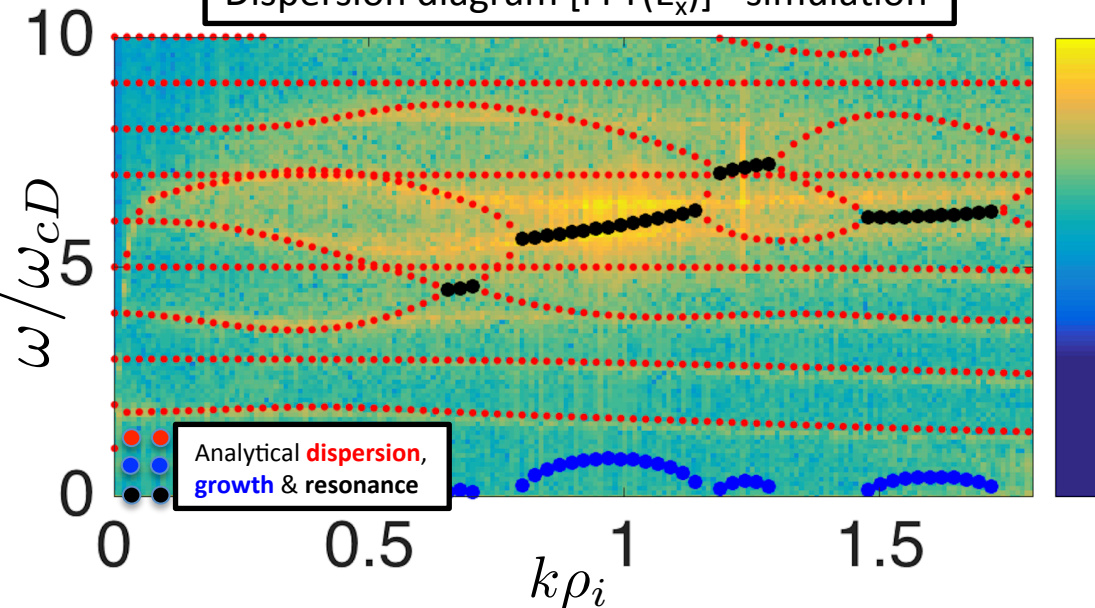
Thermal electron

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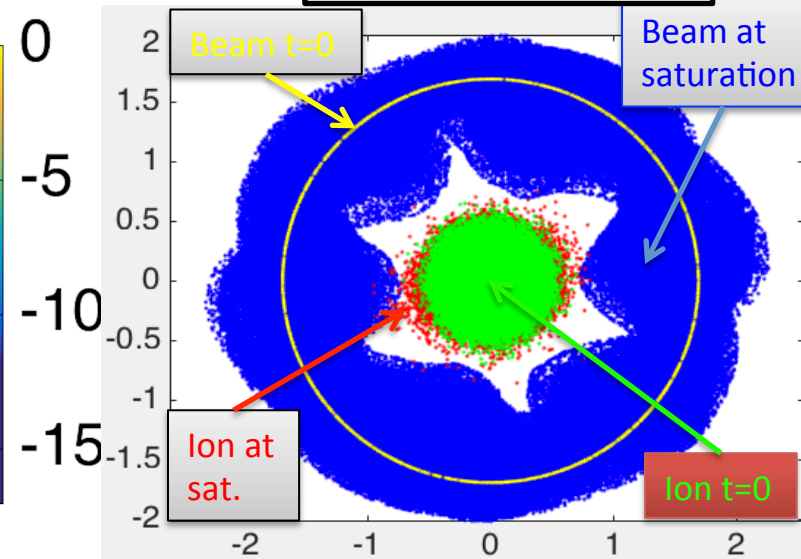
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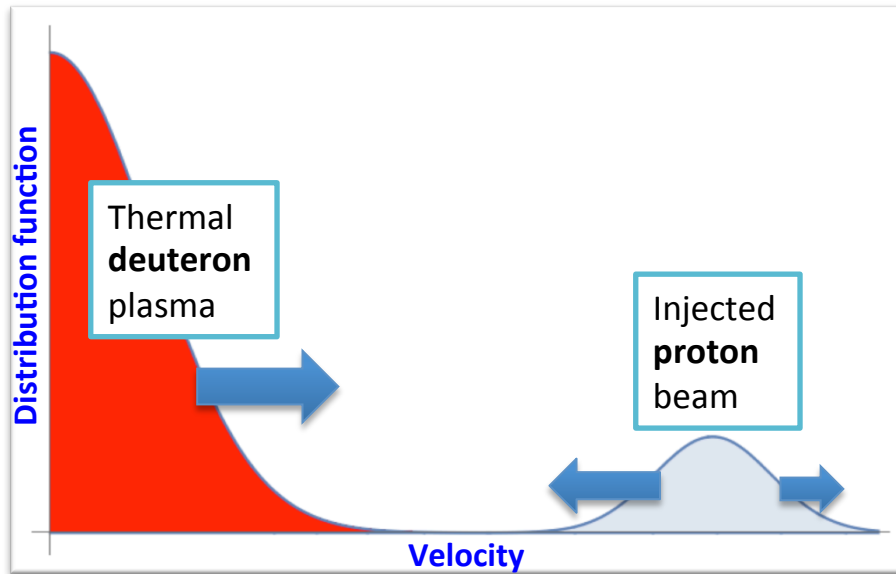
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Velocity [$\times 10^6$ m/s]



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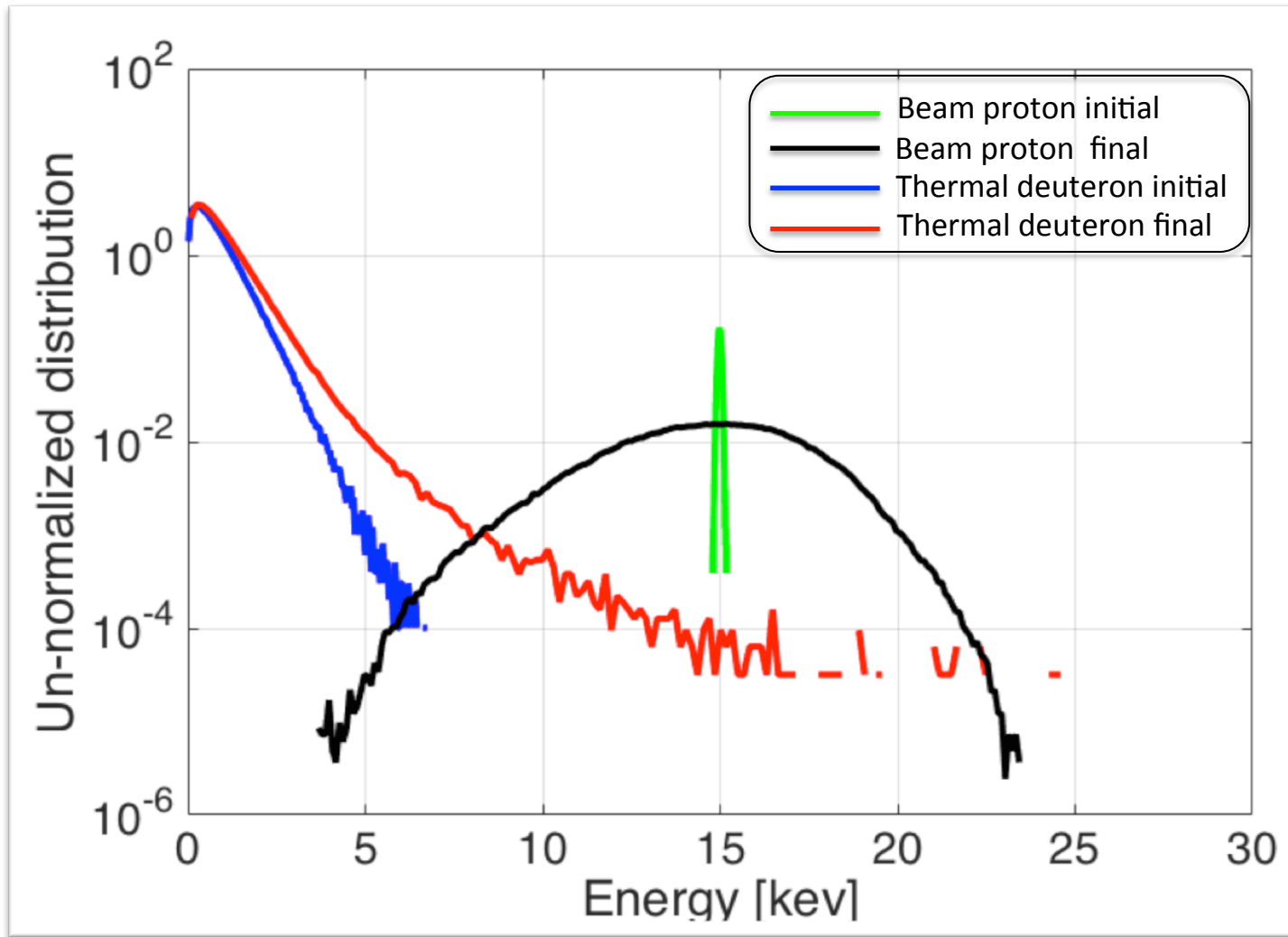
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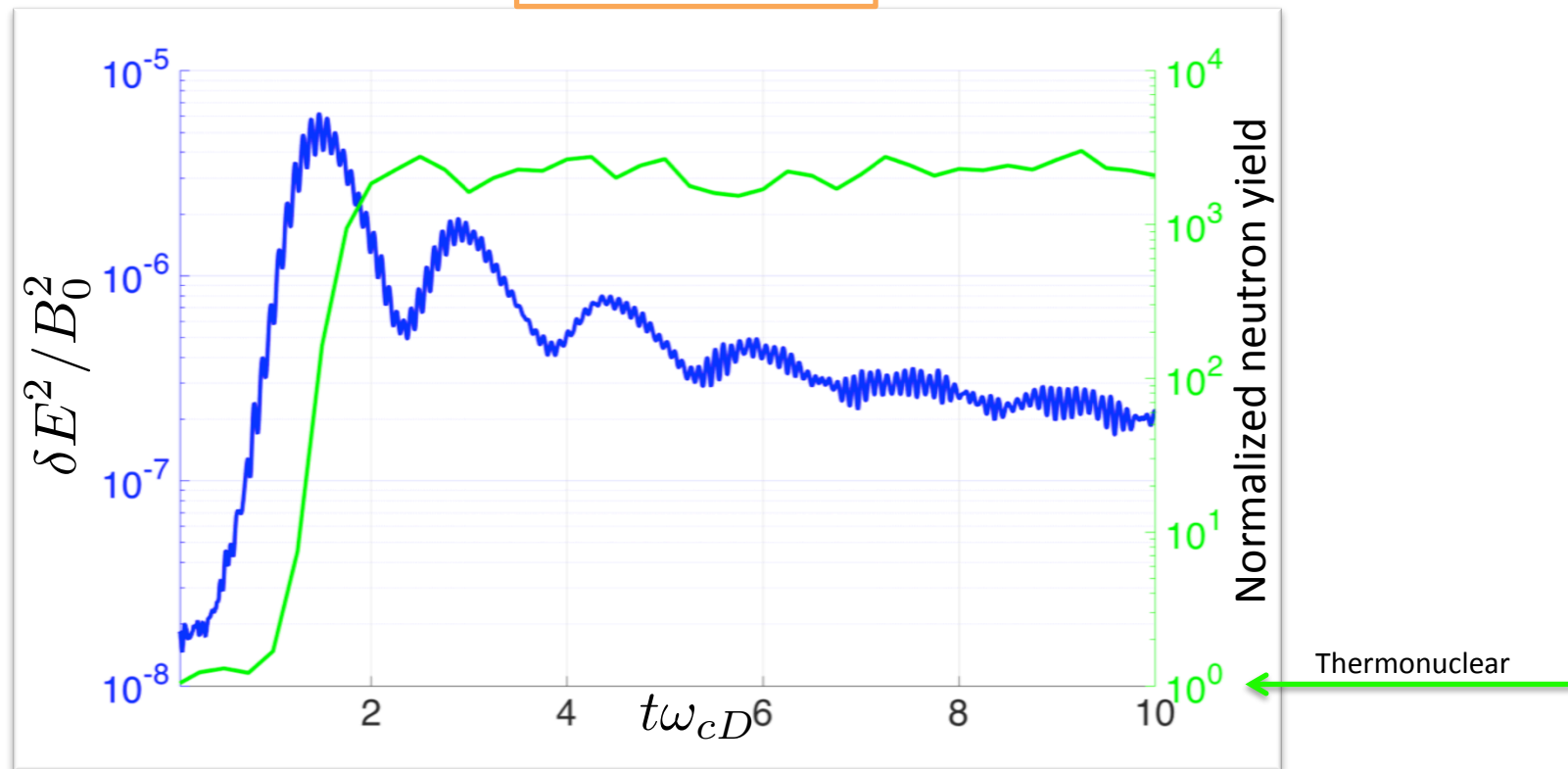
Test our Hypothesis



Klimontovich Reactivity for ES Case

$$\langle \sigma v \rangle = \sum_{i \neq j} \sum_j \sigma(\|\mathbf{v}_i - \mathbf{v}_j\|) \|\mathbf{v}_i - \mathbf{v}_j\|$$

DD-cross section



Enormous increase in reactivity due to the beam driven collective effects

Outline

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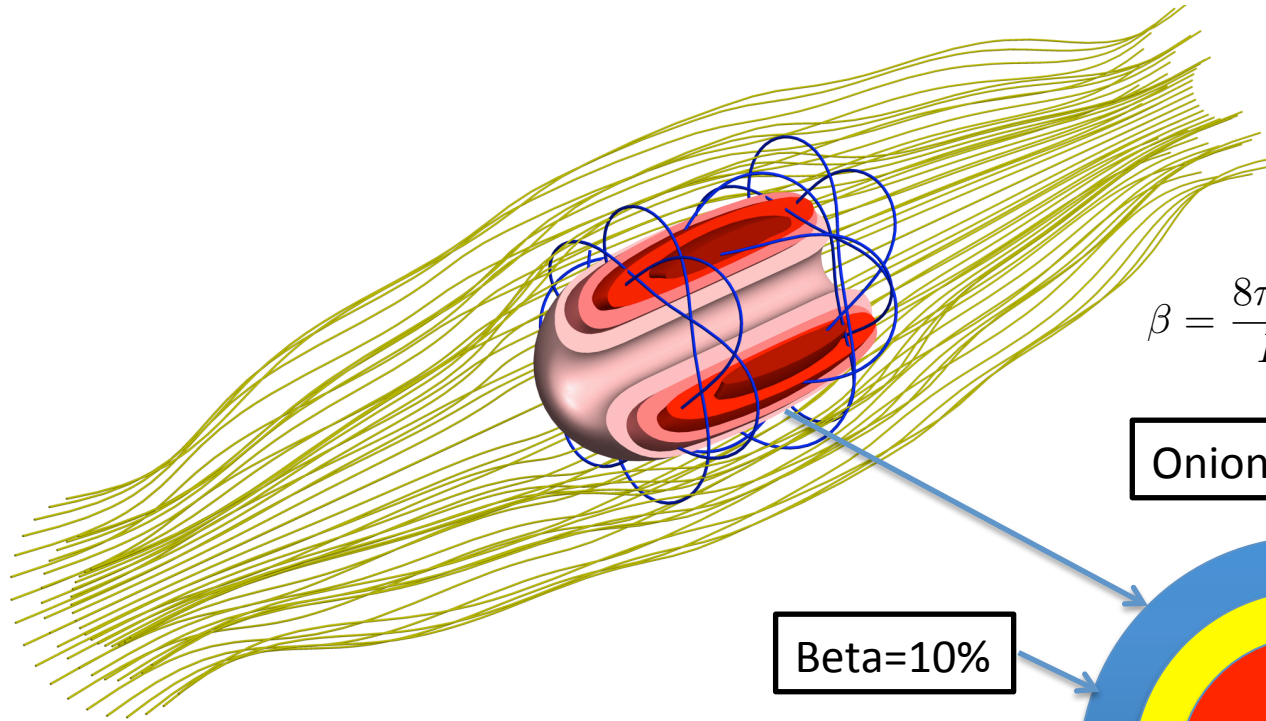
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Beam Orbit Samples FRC Volume

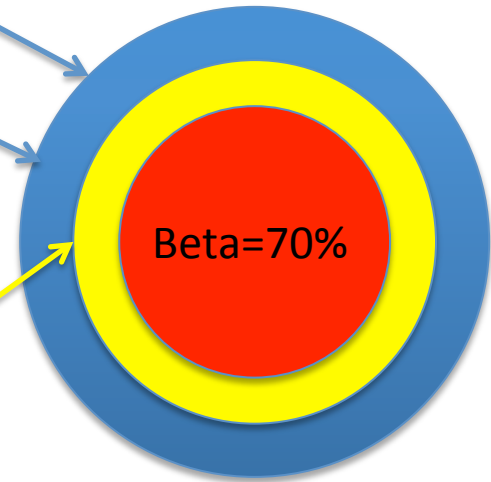


$$\beta = \frac{8\pi nT}{B^2} = \frac{\text{Plasma pressure}}{\text{Magnetic pressure}}$$

Onion Skin Model

Beta=10%

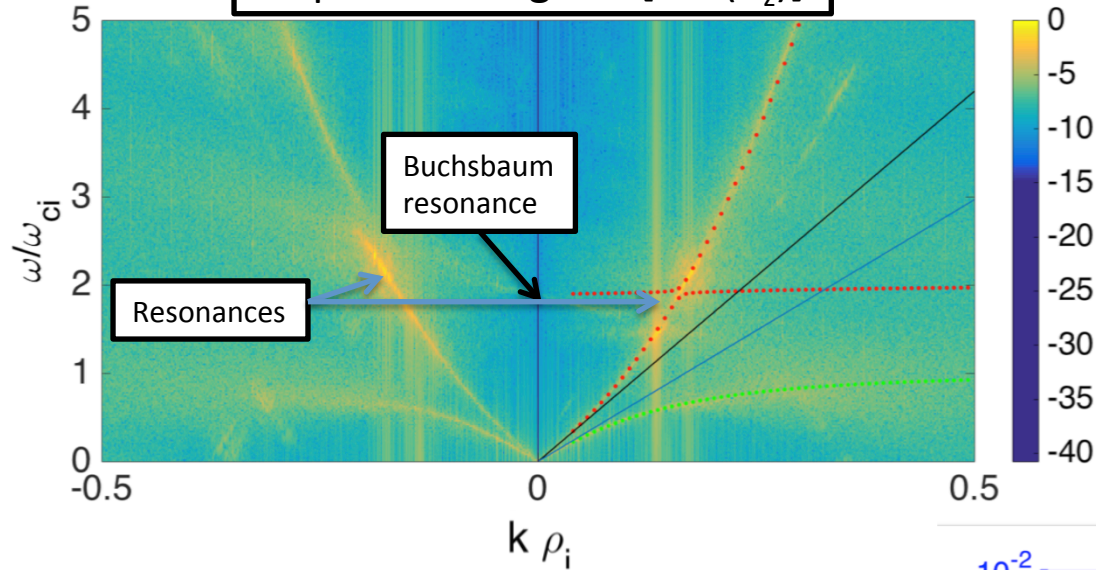
Beta=30%



- Study each onion layer separately
- Moving radially inward, increases beta

Beta=10 % -- Open field line Case

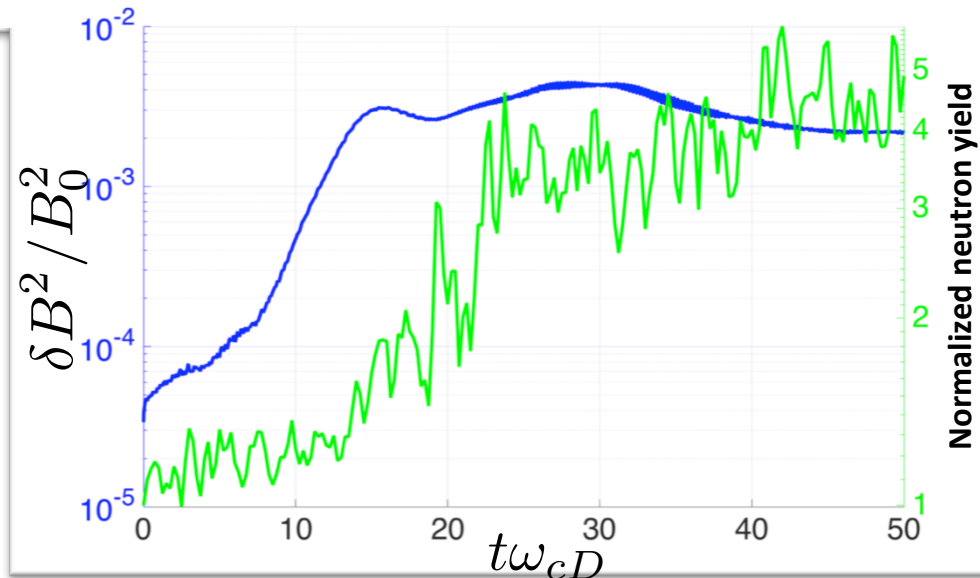
Dispersion diagram [FFT(B_z)]



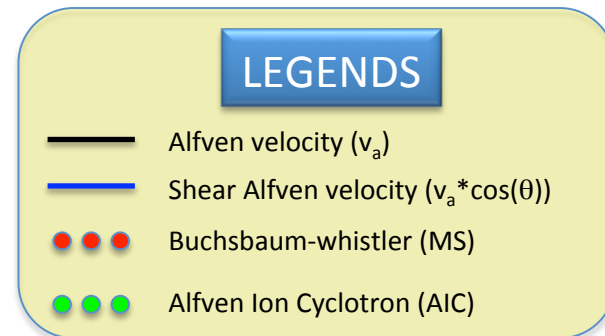
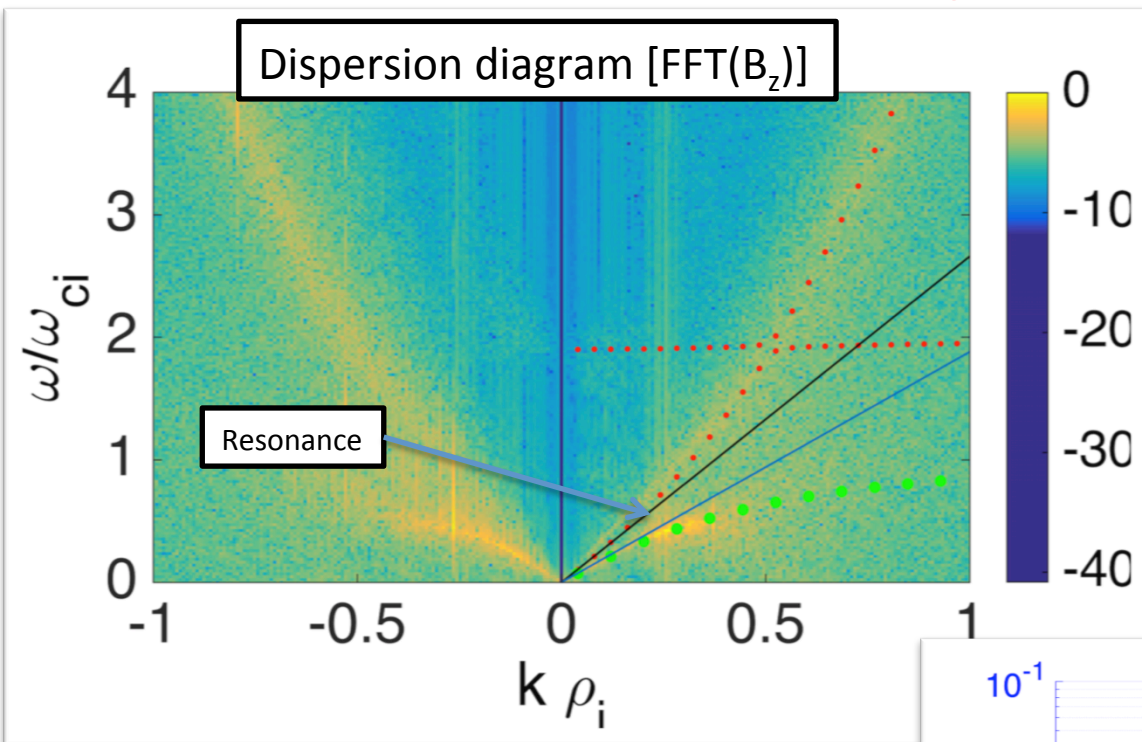
LEGENDS

- Alfvén velocity (v_a)
- Shear Alfvén velocity ($v_a \cdot \cos(\theta)$)
- ⋯ Buchsbaum-whistler (MS)
- ⋯ Alfvén Ion Cyclotron (AIC)

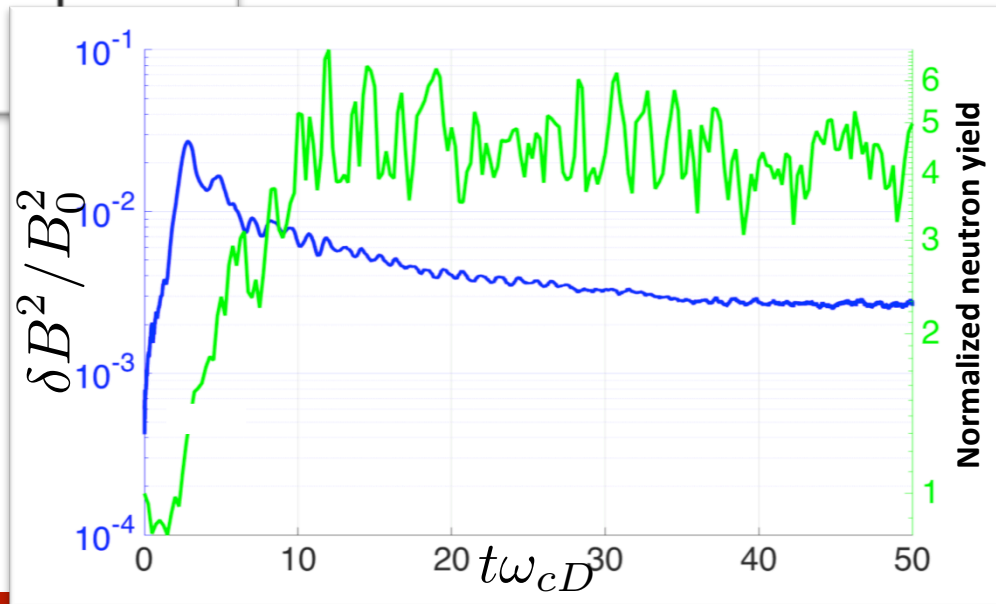
- Buchsbaum-whistler coupling facilitated by beams in presence of two q/m ion species.
- Letting $m_t = m_b$ results in the excitation of AIC mode
- Various resonances/modes are excited as beam slows down
- Slow growth ($\gamma \sim 0.03 \omega_{cD}$)
- Enhanced neutron signal



Beta=70 % -- Core

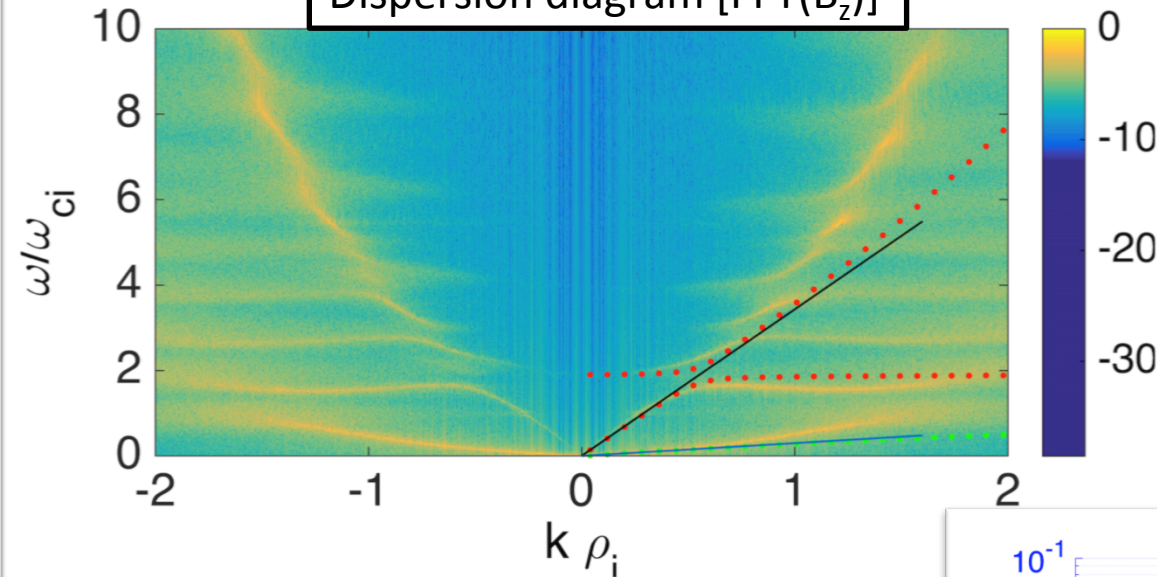


- AIC mode excited due to the presence of anisotropic ($W_{\text{perp}} > W_{\parallel}$) beam population
- Proton beam and deuteron thermal
- Only one mode excited
- Fast growth ($\gamma \sim 0.13 \omega_{cD}$)
- Enhanced neutron signal



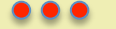



Beta=30 % -- SOL (~impact par. of NPA)

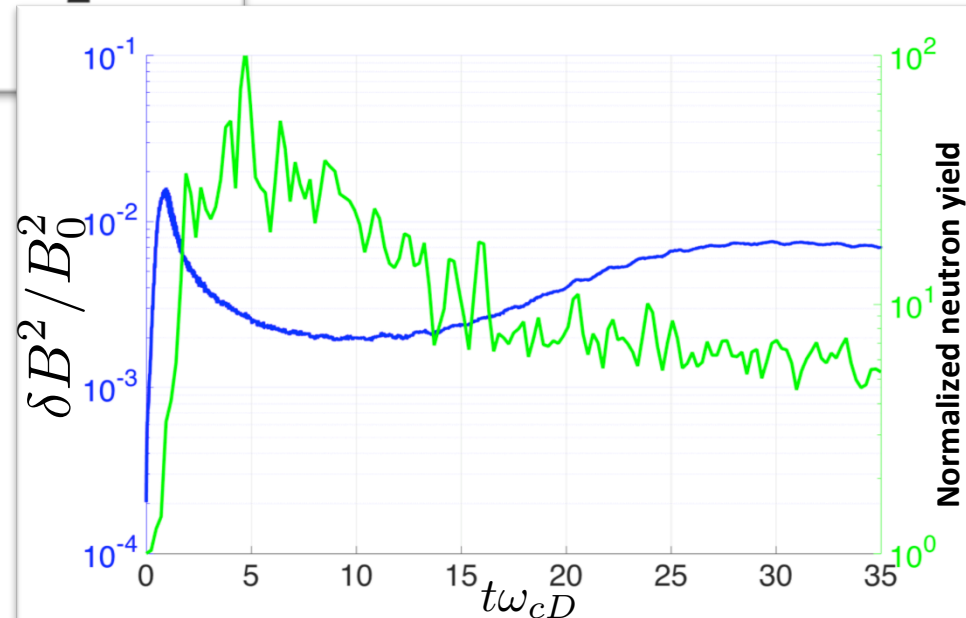
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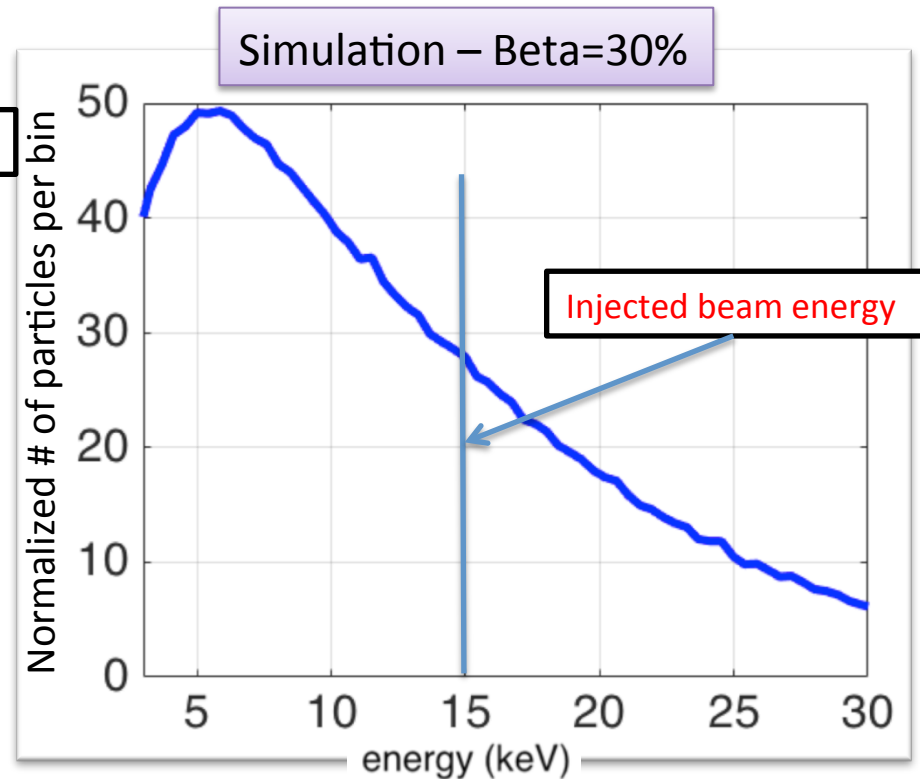
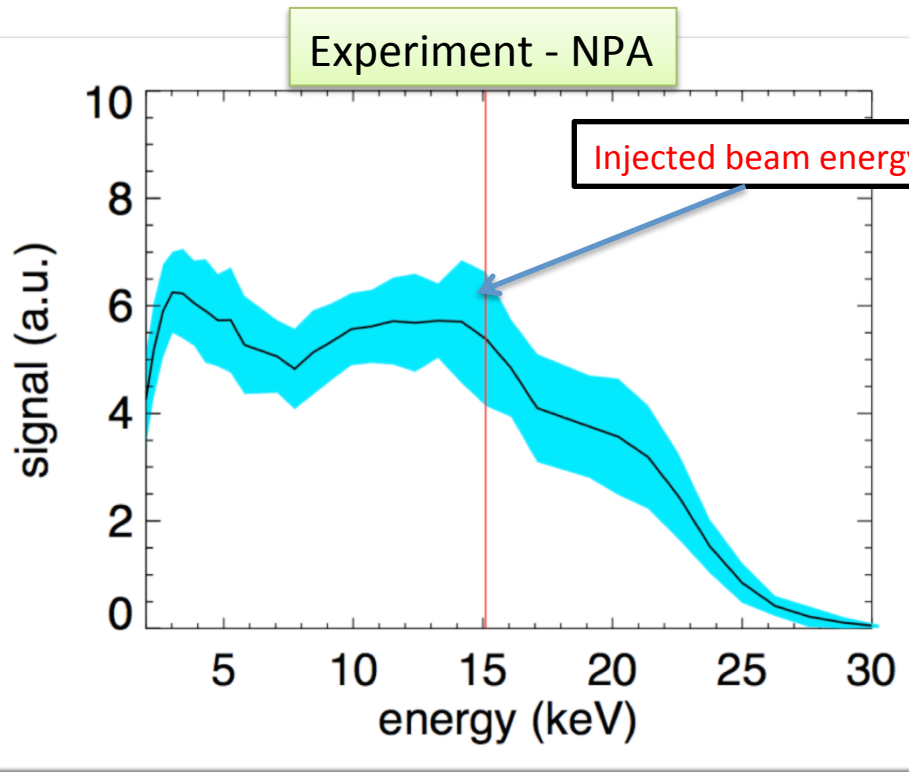
LEGENDS

-  Alfvén velocity (v_a)
-  Shear Alfvén velocity ($v_a \cdot \cos(\theta)$)
-  Buchsbaum-whistler (MS)
-  Alfvén Ion Cyclotron (AIC)

- Allow $k_{||} < k_{\text{perp}}$
- Broadband
- Ion Bernstein harmonics
- AIC mode
- **Enhanced neutron signal**
- Fast growth $\sim 0.2 \omega_{cD}$
- **Proton beam and deuteron thermal**



Proton (Beam) Energy Spectrum

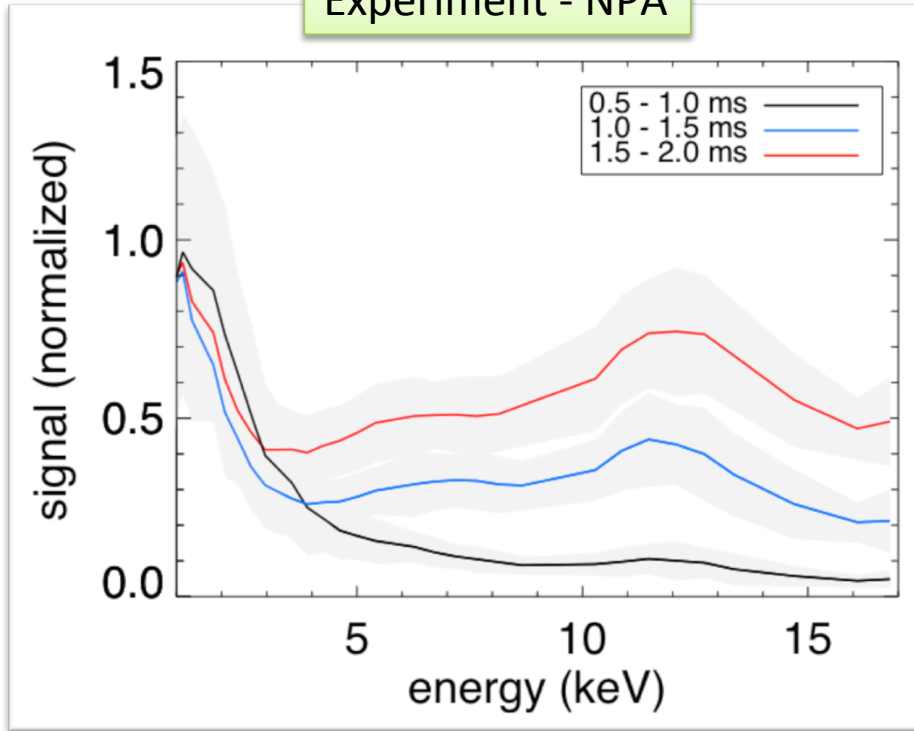


Courtesy: Ryan Clary

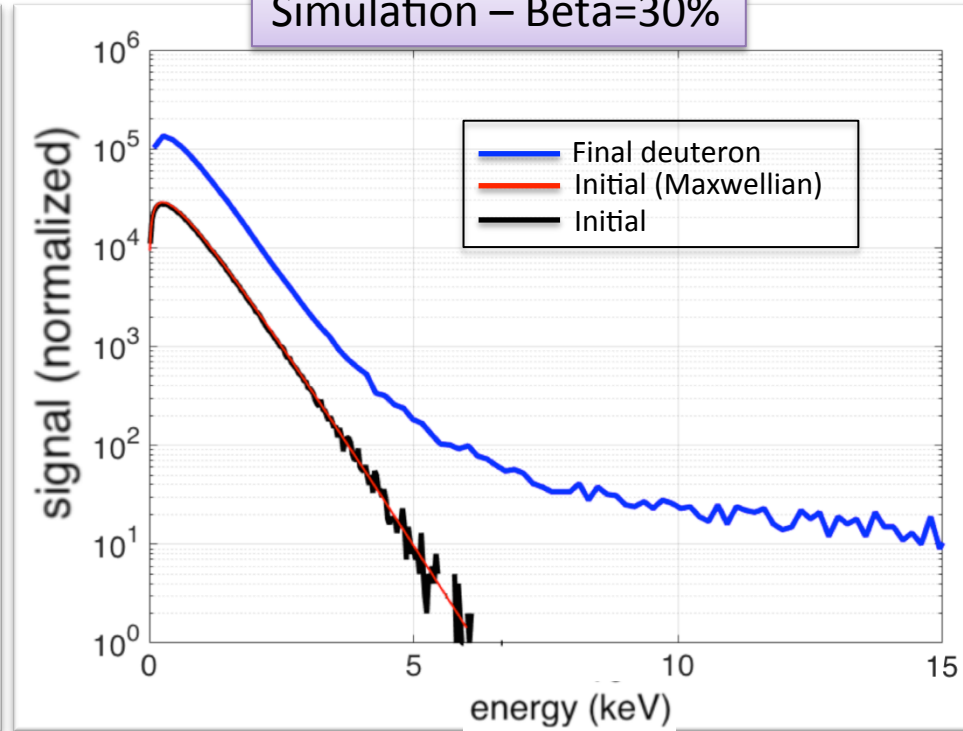
Experimental observation and simulation show scattering of beam proton to higher than injected energy

Deuteron (Thermal) Energy Spectrum

Experiment - NPA



Simulation - Beta=30%



Courtesy: Ryan Clary

Experimental observation and simulation show energetic tail in the deuteron distribution

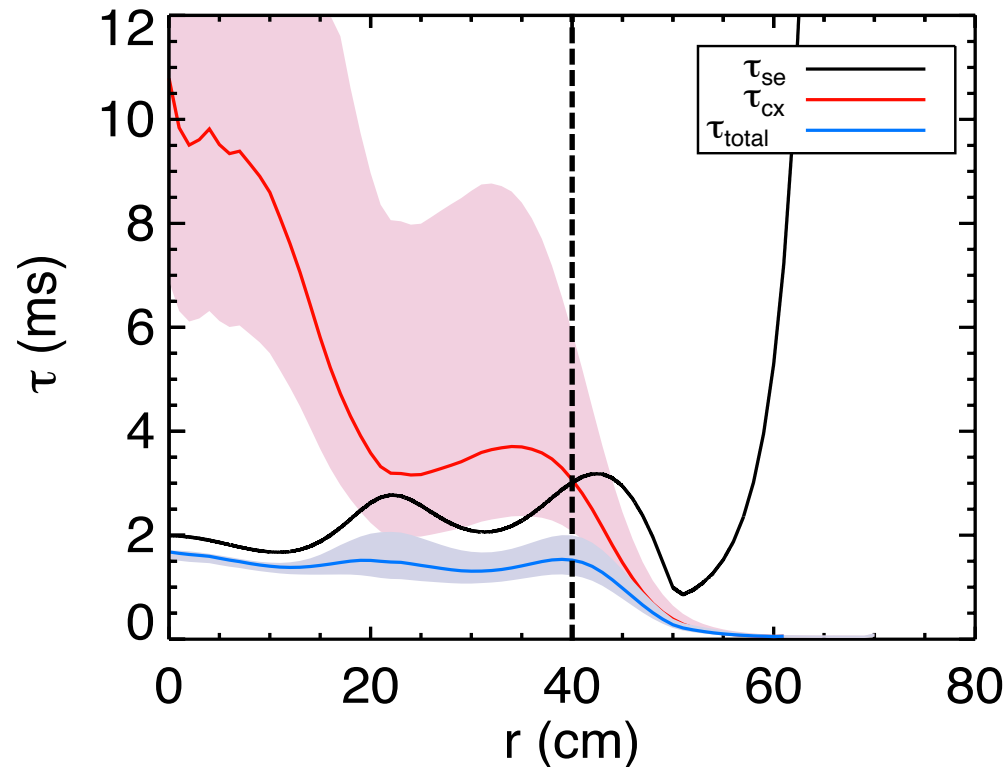
Summary

- Introduced hypothesis for enhanced neutron observed experimentally
- ES used for an analytical benchmark
- Studied three different regimes => each generating enhanced neutrons
- Concentrate on a regime near the impact parameter of the NPA
- ~10x increase in neutron yield compares well to experiment
- Despite the presence of possible robust μ -instabilities, FRC plasma remains robust and undestroyed.

Future Work

- ◆ Study beam anisotropy
- ◆ Field and density inhomogeneity
- ◆ DC steady state beam injection and modulation
- ◆ 1D3V longer runs (Perform at NERSC)
- ◆ 2D3V simulation to relax angle of propagation (Perform at NERSC)

Dominant fast ion collisional process differs inside and outside separatrix



- In FRC, $\tau_{se} < \tau_{cx} \rightarrow$ slowing down distribution
- In SOL, $\tau_{se} > \tau_{cx} \rightarrow df/dv > 0 \rightarrow$ beam-driven modes
- Fast ions sample both regions of plasma